The Stability of Difference Approximations to a Self-Adjoint Parabolic Equation, Under Derivative Boundary Conditions

By C. M. Campbell and P. Keast

1. Introduction. A self-adjoint parabolic equation in one space variable is considered, under boundary conditions which involve the function and its space derivative. A type of numerical instability can arise, which is traceable to the boundary conditions, and which is caused by the existence of unbounded solutions of the original differential equation.

2. The Differential Equation. The equation to be examined is

$$(1) \quad \frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left(g(x) \frac{\partial u}{\partial x} \right)$$

in the region

$$R \equiv [a \leq x \leq b] \times [t \geq 0]$$

subject to the initial condition

$$(1a) \quad u(x, 0) = y(x), \quad a \leq x \leq b,$$

and the boundary conditions

$$(1b) \quad \begin{aligned} \partial u / \partial x - pu &= \phi_0(t), & x &= a, & t &\geq 0, \\ \partial u / \partial x + qu &= \phi_1(t), & x &= b, \end{aligned}$$

where p, q are constants, $\phi_0(t)$ and $\phi_1(t)$ are bounded as $t \rightarrow \infty$, and there are no discontinuities in the initial or boundary conditions, or at the corners of R . In addition, we assume $g(x) > 0$, for $a < x < b$.

Several authors (e.g. [1]–[5]) have considered the equation (1) subject to (1a) and (1b) for the particular case when $p \geq 0$ and $q \geq 0$. These conditions are not imposed in this paper. In [6] Keast and Mitchell have considered the problem of the present paper with $g(x) \equiv 1$.

3. The Difference Approximations. The region R is covered by a rectangular mesh, the nodal points of which are given by

$$x_i = a + ih; \quad i = 0, 1, \dots, N; \quad Nh = b - a$$

and

$$t_n = nk, \quad n \geq 0.$$

The constants h and k are the space and time increments respectively, and the ratio k/h^2 is denoted by r .

Received May 11, 1967. Revised October 23, 1967.

We consider an approximation to Eq. (1) which can be written in the form

$$(2) \quad w_i^{n+1} - w_i^n = \frac{r}{2} \Delta(B_i \nabla w_i^n) + \frac{r}{2} \Delta(B_i \nabla w_i^{n+1}), \quad i = 0, 1, \dots, N; \quad n \geq 0,$$

where w_i^m ($m = n, n+1$) is the numerical approximation to $u(x_i, t_m)$; the difference operators Δ and ∇ are defined by

$$(3) \quad \Delta f_i = f_{i+1} - f_i \quad \text{and} \quad \nabla f_i = f_i - f_{i-1},$$

and the coefficients B_i ($i = 0, 1, \dots, N+1$) are defined, using a method developed by Tikhonov and Samarskii [7], by

$$(4) \quad B_i = \left[\int_{x_{i-1}}^{x_i} \frac{dx}{g(x)} \right]^{-1}, \quad i = 0, 1, \dots, (N+1).$$

In Eq. (2), when $n = 0$, the values of w_i^0 ($i = 0, 1, \dots, N$) are obtained from the initial conditions (1a). When $i = 0$ or N , the values of w outside R , viz. w_{-1}^m and w_{N+1}^m ($m = n, n+1$), which occur in (2) are eliminated using the boundary conditions (1b), in which the derivatives are approximated by the equation

$$(5) \quad \partial u_j^m / \partial x = (w_{j+1}^m - w_{j-1}^m) / 2h, \quad j = 0, N \quad \text{and} \quad m = n, n+1.$$

If we denote the vector of values of w at each time level t_m by $\mathbf{w}^m$, i.e.

$$\mathbf{w}^m = \{w_0^m, w_1^m, \dots, w_N^m\}^T$$

and use the boundary conditions (1b) in (2) when $i = 0$ and $i = N$, we may write the totality of difference equations in the form

$$(6) \quad (I + rU/2)\mathbf{w}^{n+1} = (I - rU/2)\mathbf{w}^n + \mathbf{l}^n$$

where $\mathbf{l}^n$ is a vector involving the boundary conditions (with $l_1^n = l_2^n = \dots = l_{N-1}^n = 0$, for all n) and I and U are matrices of order $(N+1)$. The matrix I is the identity matrix, and U is given by

$$(7) \quad U = \begin{bmatrix} B_1 + B_0(1 + 2p/N) & -(B_1 + B_0) & & & \\ -B_1 & & B_1 + B_2 & & -B_2 \\ & \cdot & & \cdot & \\ & & \cdot & & \cdot \\ & & & \cdot & \\ & & & & \cdot \\ & & -B_{N-1} & (B_{N-1} + B_N) & -B_N \\ & & & -(B_N + B_{N+1}) & B_N + B_{N+1}(1 + 2q/N) \end{bmatrix}.$$

Let the numerical error of the difference method (6) be denoted by

$$\boldsymbol{\varepsilon}^n = \mathbf{w}^n - \hat{\mathbf{w}}^n, \quad \forall n \geq 0,$$

where $\mathbf{w}^n$ is the exact solution of Eq. (6) and $\hat{\mathbf{w}}^n$ is the numerical solution. Then we obtain the error equation

$$(I + rU/2)\boldsymbol{\varepsilon}^{n+1} = (I - rU/2)\boldsymbol{\varepsilon}^n.$$

We shall assume $(I + rU/2)$ to be nonsingular, so that

$$(8) \quad \mathbf{e}^{n+1} = (I + rU/2)^{-1}(I - rU/2)\mathbf{e}^n.$$

The difference method (6) will be stable, therefore, if and only if the norms $\|\mathbf{e}^n\|$ of the vectors given by (8) are bounded. Since

$$\mathbf{e}^{n+1} = G^{n+1}\mathbf{e}^0$$

where

$$G = (I + rU/2)^{-1}(I - rU/2)$$

then the norms $\|\mathbf{e}^n\|$ are bounded if and only if there is a constant K such that

$$\|G^{n+1}\| \leq K, \quad \forall n \geq 0.$$

The matrix U is not symmetric, but it is similar to a symmetric matrix $\tilde{U}$, where

$$(9) \quad \tilde{U} = D^{-1}UD$$

with D given by

$$D = \begin{bmatrix} ((B_1 + B_0)/B_1)^{1/2} & & & & \\ & 1 & & & \\ & & \ddots & & \\ & & & \ddots & \\ & & & & 1 & \\ & & & & & ((B_N + B_{N+1})/B_N)^{1/2} \end{bmatrix}.$$

Hence the eigenvalues of U are all real.

Thus G is similar to a symmetric matrix $\tilde{G}$, where

$$\tilde{G} = D^{-1}GD$$

and if $\rho(G)$ is the spectral radius of G (the maximum modulus eigenvalue) then

$$\|\tilde{G}^{n+1}\| = \rho^{n+1}(G).$$

Clearly $\|G^{n+1}\|$ is bounded if and only if $\|\tilde{G}^{n+1}\|$ is bounded. The stability condition is therefore: there exists $K > 0$, for all $n > 0$, $\rho^{n+1}(G) \leq K$.

This condition clearly implies

$$(10) \quad \rho(G) \leq 1.$$

It should be noted that condition (10) is the condition for asymptotic stability ($n \rightarrow \infty$) of Eq. (6). If the errors are to be bounded in a closed region $0 \leq t \leq T$, for some finite T , then (10) may be weakened to the condition

$$\rho(G) \leq 1 + O(k).$$

This point has been discussed by Gary [8]. In this paper we take (10) as our condition for stability.

If we denote the $(N + 1)$ real eigenvalues of the matrix U by $\lambda_0, \lambda_1, \dots, \lambda_N$, then condition (10) is satisfied if and only if

$$|(1 - r\lambda_i/2)/(1 + r\lambda_i/2)| \leq 1, \quad i = 0, 1, \dots, N.$$

We shall assume that $1 + r\lambda_i/2 > 0$ (an assumption which is justified later) and hence the necessary and sufficient condition for stability becomes

$$\lambda_i \geq 0, \quad i = 0, 1, \dots, N.$$

4. The Eigenvalues of U . If we apply the theorem of Gerschgorin to the matrix U in (7), we find that the spectrum of U lies within the union of the intervals

$$\begin{aligned} 0 \leq \lambda \leq 2(B_i + B_{i+1}), \quad i = 1, 2, \dots, (N-1), \\ (2B_0/N)p \leq \lambda \leq 2[B_0(1 + p/N) + B_1], \\ (2B_N/N)q \leq \lambda \leq 2[B_{N+1}(1 + q/N) + B_N]. \end{aligned}$$

If $p \geq 0$ and $q \geq 0$ then U is positive definite (or positive semidefinite if $p = q = 0$), and Eq. (6) is then stable. This is the case which has been considered by the authors [1]–[5]. If any eigenvalues of U are negative, then they are $O(1/N)$, and so for large N we may assume that $1 + r\lambda_i > 0$, $i = 0, 1, \dots, N$. This justifies the assumption made in Section 3.

We now consider for what values of p and q the matrix U is singular, i.e. $\det(U) = 0$. Expanding this determinantal equation we obtain

$$\begin{aligned} (4pq/N^2)B_0B_{N+1}T_{N-1}^{(1)} + (2p/N)B_0(B_N + B_{N+1})(T_{N-1}^{(1)} - B_NT_{N-2}^{(1)}) \\ (11) \quad + 2(q/N)B_{N+1}(B_0 + B_1)(T_{N-1}^{(1)} - B_1T_{N-2}^{(2)}) \\ + (B_0 + B_1)(B_N + B_{N+1})(T_{N-1}^{(1)} - B_NT_{N-2}^{(1)} - B_1T_{N-2}^{(2)} + B_1B_NT_{N-3}^{(2)}) = 0, \end{aligned}$$

where

$$T_r^{(s)} = \det \begin{bmatrix} (B_s + B_{s+1}) & & & & -B_{s+1} \\ -B_{s+1} & & & & \\ & \ddots & & & \\ & & \ddots & & \\ & & & \ddots & \\ & & & & -B_{s+r-1} \\ & & & & (B_{s+r-1} + B_{s+r}) \end{bmatrix}$$

is the determinant of an $(r \times r)$ symmetric matrix. The constant term (i.e. the term independent of p and q) in (11) is the expansion of $\det U$ when $p = q = 0$. But if $p = q = 0$, U is singular (the sum along each row is zero), and so this constant is zero.

It is easily shown that

$$T_r^{(s)} = \sum_{j=0}^r \left\{ \prod_{i=s; i \neq j+s}^{r+s} B_i \right\}$$

and hence that

$$T_{N-1}^{(1)} - B_NT_{N-2}^{(1)} = \prod_{i=1}^{N-1} B_i$$

and

$$T_{N-1}^{(1)} - B_1T_{N-2}^{(2)} = \prod_{i=2}^N B_i.$$

Equation (11) then becomes

$$(12) \quad \frac{4pq}{N^2} B_0 B_{N+1} \sum_{j=0}^{N-1} \prod_{i=1; i \neq j+1}^N B_i + \frac{2p}{N} \frac{(B_N + B_{N+1})}{B_N B_{N+1}} \prod_{i=0}^{N+1} B_i + \frac{2q}{N} \frac{B_0 + B_1}{B_0 B_1} \prod_{i=0}^{N+1} B_i = 0.$$

This may be rearranged to give

$$(12)' \quad \frac{4pq}{N^2} \sum_{i=1}^N \frac{1}{B_i} + \frac{2p}{N} \frac{B_N + B_{N+1}}{B_N B_{N+1}} + \frac{2q}{N} \frac{B_0 + B_1}{B_0 B_1} = 0.$$

For the method of Tikhonov and Samarskii this equation takes the form

$$(12a) \quad \frac{4pq}{N^2} \int_a^b \frac{dx}{g(x)} + \frac{2p}{N} \int_{x_{N-1}}^{x_{N+1}} \frac{dx}{g(x)} + \frac{2q}{N} \int_{x_{-1}}^{x_1} \frac{dx}{g(x)} = 0.$$

It should be noted that this is an approximation, to $O(1/N^2)$, of the equation

$$(13) \quad pq \int_a^b \frac{dx}{g(x)} + \frac{p}{g(b)} + \frac{q}{g(a)} = 0.$$

The matrix U is singular when p and q lie on a hyperbola, which we denote by $\Sigma = 0$, where

$$\Sigma \equiv pq + Lp + Mq = 0.$$

The coefficients L, M are both positive. We partition the (p, q) plane into sections A, B, D where

- (i) A is $pq + Lp + Mq \geq 0, p + q \geq 0$,
- (ii) B is $pq + Lp + Mq < 0$,
- (iii) D is $pq + Lp + Mq \geq 0, p + q < 0$,

(see Fig. 1).

Let $z(\lambda)$ be the characteristic polynomial of the matrix U , i.e.

$$z(\lambda) \equiv \det (U - \lambda I) .$$

Since U is similar to a symmetric matrix, the polynomial $z(\lambda)$ has $(N + 1)$ real zeros; in addition $z(\lambda)$ is clearly positive for large negative λ . But in the region B , $z(0)$ is negative, and so there must be a zero of $z(\lambda)$ in $\lambda < 0$. In fact $z(\lambda)$ must have an odd number of negative roots. At the point $p = -N/2$, $q = -N/2$ in D the matrix U has at least one negative root, since

$$U = \begin{bmatrix} B_1 & -(B_0 + B_1) & & & & \\ -B_1 & (B_1 + B_2) & & -B_2 & & \\ & \cdot & \cdot & \cdot & \cdot & \\ & & \cdot & \cdot & \cdot & \\ & & & -B_{N-1} & (B_{N-1} + B_N) & -B_N \\ & & & & -(B_N + B_{N+1}) & B_N \end{bmatrix}$$

which is not positive definite, since

$$U \begin{bmatrix} 1 \\ 1 \\ \vdots \\ 1 \end{bmatrix} = \begin{bmatrix} -B_0 \\ 0 \\ 0 \\ \vdots \\ 0 \\ -B_{N+1} \end{bmatrix} < 0 .$$

Thus, since $z(-\infty) > 0$ and $z(0) > 0$ in D , there is an even number of negative roots at the point $(-N, -N)$.

We now exhibit the dependence of $z(\lambda)$ on (p, q) by writing the characteristic polynomial as $z(p, q, \lambda)$. The function $z = 0$ is a surface in (p, q, λ) space, cutting each vertical line $p = \text{constant}$, $q = \text{constant}$, in $(N + 1)$ points. Let $\Gamma(p, q) = 0$ be any curve in the (p, q) plane. Let $\Gamma_{\lambda_0} = 0$ be the curve in (p, q, λ) space, on the surface $z = 0$, which consists of the points $(p, q, \lambda_0(p, q))$, where λ_0 is the smallest root of $z = 0$, for each point (p, q) on $\Gamma(p, q) = 0$. Thus $\Gamma = 0$ is the projection of $\Gamma_{\lambda_0} = 0$ on the (p, q) plane. Clearly, as (p, q) moves, the curve $\Gamma_{\lambda_0} = 0$ is continuous. Suppose that (p, q) moves from a point in the region A where $p > 0$ and $q > 0$, along $\Gamma = 0$. Then initially $\lambda_0 > 0$, and as (p, q) moves, λ_0 cannot change sign unless $\Gamma_{\lambda_0} = 0$ crosses the (p, q) plane; i.e. unless λ_0 passes through a zero. This cannot happen unless $\Gamma_{\lambda_0} = 0$ and $\Gamma = 0$ pass through a point on $\Sigma = 0$ in the (p, q) plane. Thus at every point in A , $\lambda_0 > 0$ except on the boundary $\Sigma = 0$, when $\lambda_0 = 0$, and the other eigenvalues are positive. If $\Gamma = 0$ moves into the region B in which U has an odd number of eigenvalues, λ_0 stays negative and no other eigenvalue changes sign. (A similar argument with $\Gamma_{\lambda_1} = 0$ will prove that λ_1 stays positive as $\Gamma = 0$ moves into the region B .)

At one point in D there are at least two negative eigenvalues λ_0 and λ_1 . Consid-

eration of the curves $\Gamma_{\lambda_0} = 0$ and $\Gamma_{\lambda_1} = 0$, as (p, q) crosses the boundary between B and D , shows that on the boundary $\lambda_0 < 0$ and $\lambda_1 = 0$, and in D there are exactly two negative roots.

Thus we may summarize as follows:

- (i) In the region A all eigenvalues are positive.
- (ii) On the boundary between A and B one eigenvalue is zero and the rest are positive.
- (iii) In the region B one eigenvalue is negative and the rest are positive.
- (iv) On the boundary between B and D one eigenvalue is negative, one eigenvalue is zero, and the rest are positive.
- (v) In D there are two negative eigenvalues and the rest are positive.

The difference scheme (6) will be stable therefore if and only if the values of p and q come from the region A or the boundary between the regions A and B .

5. Numerical Results. In order to demonstrate the results of the preceding sections, we considered the problem

$$\frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left(\frac{1}{x^L} \frac{\partial u}{\partial x} \right)$$

in the region

$$[2 \leq x \leq 3] \times [t \geq 0]$$

subject to the initial conditions

$$u(x, 0) = x^{L+1} - 4, \quad 2 \leq x \leq 3,$$

and the boundary conditions

$$\begin{aligned} \partial u / \partial x - pu &= (L+1)2^L - p(2^{L+1} - 4), & x = 2, \\ \partial u / \partial x + qu &= (L+1)3^L + q(3^{L+1} - 4), & x = 3, \end{aligned}$$

where L was taken equal to 1 and 3.

The method of Tikhonov and Samarskii was used to solve this problem with $N = 20$ and $r = 1$ and with the following values of p and q for $L = 1$ and $L = 3$:

Problem	(1)	(2)	(3)	(4)	(5)
$L = 1$					
$p =$	2.00	-0.20	-0.36	-2.00	-3.00
$q =$	1.00	1.00	1.00	1.00	-3.00
$L = 3$					
$p =$	2.00	-0.10	-0.18	-2.00	-3.00
$q =$	1.00	1.00	1.00	1.00	-3.00

Table 1 shows the maximum error for values of $K = 200(200)2000$ where K is the number of steps in time. Problems (1), (2), (3) are taken with (p, q) in the region A and are stable; problems (4) and (5) have (p, q) in the regions B and D respectively and are unstable.

TABLE 1

$L = 1$ K	(1)	(2)	(3)	(4)	(5)
200	$.3170 \times 10^{-4}$	$.3440 \times 10^{-4}$	$.3350 \times 10^{-4}$	$.6599 \times 10^{-4}$	$.1618 \times 10^{-3}$
400	$.4260 \times 10^{-4}$	$.6150 \times 10^{-4}$	$.6120 \times 10^{-4}$	$.3151 \times 10^{-3}$	$.2397 \times 10^{-2}$
600	$.4330 \times 10^{-4}$	$.8719 \times 10^{-4}$	$.9477 \times 10^{-4}$	$.1089 \times 10^{-2}$	$.2964 \times 10^{-1}$
800	$.4330 \times 10^{-4}$	$.1147 \times 10^{-4}$	$.1291 \times 10^{-3}$	$.3470 \times 10^{-2}$	.3547
1000	$.4330 \times 10^{-4}$	$.1407 \times 10^{-3}$	$.1638 \times 10^{-3}$	$.1079 \times 10^{-1}$	4.2105
1200	$.4330 \times 10^{-4}$	$.1645 \times 10^{-3}$	$.1986 \times 10^{-3}$	$.3327 \times 10^{-1}$	49.883
1400	$.4330 \times 10^{-4}$	$.1867 \times 10^{-3}$	$.2335 \times 10^{-3}$	.1024	590.64
1600	$.4330 \times 10^{-4}$	$.2070 \times 10^{-3}$	$.2685 \times 10^{-3}$	.3148	6992.5
1800	$.4330 \times 10^{-4}$	$.2250 \times 10^{-3}$	$.3032 \times 10^{-3}$	.9676	82779
2000	$.4330 \times 10^{-4}$	$.2415 \times 10^{-3}$	$.3382 \times 10^{-3}$	2.9740	979943

$L = 3$ K	(1)	(2)	(3)	(4)	(5)
200	.0038	.0058	.0060	.0100	.0142
400	.0046	.0084	.0087	.0195	.0366
600	.0052	.0103	.0107	.0319	.0798
800	.0057	.0119	.0125	.0487	.1645
1000	.0060	.0133	.0140	.0715	.3312
1200	.0062	.0145	.0154	1.0287	.6604
1400	.0062	.0156	.0167	1.4620	1.3113
1600	.0062	.0166	.0179	2.0613	2.5995
1800	.0062	.0175	.0190	2.8917	5.1499
2000	.0062	.0185	.0201	4.0430	10.2007

6. The Differential Equation. Consider Eq. (1) subject to the initial condition (1a) and the boundary conditions

$$\begin{aligned} \partial u / \partial x - pu &= 0, & x &= a, \\ \partial u / \partial x + qu &= 0, & x &= b, \end{aligned} \quad t \geq 0.$$

Let $u(x, t) = X(x)T(t)$ so that

$$(15) \quad dT/dt = -\lambda T$$

and

$$(16) \quad (d/dx)(g(x)dX(x)/dx) = -\lambda X,$$

where λ is a constant and

$$(17) \quad X'(a) - pX(a) = 0, \quad X'(b) + qX(b) = 0.$$

Eq. (16) subject to the boundary conditions (17) constitutes a Sturm-Liouville problem for the eigenfunctions $X(x)$ and eigenvalues λ . The eigenvalues are real, distinct and ordered: $\lambda_0 < \lambda_1 < \dots$ and the corresponding eigenfunctions $X_i(x)$ ($i = 0, 1, \dots$) form an orthonormal set in $a \leq x \leq b$ [9]. Then the solution of Eq. (1) subject to the initial and boundary conditions is

$$(18) \quad u(x, t) = \sum_{i=0}^{\infty} A_i X_i(x) e^{-\lambda_i t}$$

where

$$A_i = \int_a^b X_i(x) y(x) dx.$$

The solution (18) is uniformly bounded as $t \rightarrow \infty$ if and only if all the eigenvalues λ_i are nonnegative; i.e. if and only if $\lambda_0 \geq 0$.

It has been shown [10] that if $p \geq 0$ and $q \geq 0$ then $\lambda_0 \geq 0$. Thus the solutions of the differential equation are bounded in part of the region A in Fig. 1. We wish to prove that the differential equation has bounded solutions if and only if (p, q) is a point in A or on the boundary between A and B ; that is if and only if the difference scheme is stable.

The Eq. (16) has a zero eigenvalue if and only if

$$(d/dx)[g(x)(dX(x)/dx)] = 0$$

i.e. when

$$X(x) = C \int_a^b \frac{1}{g(x)} dx + X(a)$$

where $X(a)$ and C are constants which are obtained from the boundary conditions. The conditions (17) then give the relation

$$(19) \quad pq \int_a^b \frac{dx}{g(x)} + \frac{p}{g(b)} + \frac{q}{g(a)} = 0$$

which is equation (13). Thus there is a zero eigenvalue of the Sturm-Liouville problem only on the curve to which Σ tends as $N \rightarrow \infty$. By means of an argument similar to the one used on $z(\lambda)$ for the discrete case, it may be shown that $\lambda_0 \geq 0$ everywhere in the region A or on the upper branch of the curve (19), where $\lambda_0 = 0$. In addition it may be shown that there is one negative eigenvalue between the branches of the curve (19); one negative and one zero eigenvalue on the lower branch; and two negative eigenvalues in the region inside the lower branch of the curve (19), corresponding to the region D .

7. Conclusion. The numerical instability observed in this paper has therefore been traced to the existence of unbounded solutions of the differential equation as $t \rightarrow \infty$. This type of instability will be apparent only if the difference methods are run for large values of $t = nk$. In some problems it is necessary to run the difference methods beyond the stage where initial transients die out (see e.g. [8]) and in such problems asymptotic instability will clearly be of importance. Another important instance of this kind of instability will occur in the iterative solution of systems of equations arising from numerical approximations to Laplace's equation in two or more space variables. It is clear that derivative boundary conditions of the type discussed in this paper will have an effect on the eigenvalues of the matrices occurring in the system and may prevent convergence. This problem, however, will be discussed in a later paper.

The calculations were carried out on the IBM 1620 computer of the University of St. Andrews.

Department of Mathematics
University of St. Andrews
St. Andrews, Fife, Scotland

1. R. S. VARGA, *Matrix Iterative Analysis*, Prentice-Hall, Englewood Cliffs, N. J., 1962. MR 28 #1725.
2. M. E. ROSE, "On the integration of non-linear parabolic equations by implicit difference methods," *Quart. J. Appl. Math.*, v. 14, 1956, pp. 237-248. MR 19, 65.
3. J. F. GUNN, "The solution of elliptic difference equations by semi-explicit iterative techniques," *J. SIAM Ser. B Numer. Anal.*, v. 2, 1965, pp. 24-45. MR 31 #4199.
4. A. A. SAMARSKII, "Homogeneous difference schemes on non-linear equations of parabolic type," *Ž. Vyčisl. Mat. i. Mat. Fiz.*, v. 2, 1962, pp. 25-26. (Russian) MR 33 #5144.
5. I. V. FRYAZINOV, "On a difference approximation of the boundary conditions for the third boundary-value problem," *Ž. Vyčisl. Mat. i. Mat. Fiz.*, v. 4, 1964, pp. 1106-1112. (Russian) MR 31 #895.
6. P. KEAST & A. R. MITCHELL, "On the instability of the Crank Nicholson formula under derivative boundary conditions," *Comput. J.*, v. 9, 1966, pp. 110-114. MR 33 #5137.
7. A. N. TIKHONOV & A. A. SAMARSKII, "Homogeneous difference schemes," *Ž. Vyčisl. Mat. i. Mat. Fiz.*, v. 1, 1961, pp. 5-63. (Russian) MR 29 #5391.
8. J. GARY, "A generalization of the Lax-Richtmyer Theorem on finite difference schemes," *SIAM J. Numer. Anal.*, v. 3, 1966, pp. 467-473. MR 34 #5311.
9. R. COURANT & D. HILBERT, *Methods of Mathematical Physics*, Vol. 1, Interscience, New York, 1953. MR 16, 426.
10. N. S. KOSHYLYAKOV, M. M. SMIRNOV & E. B. GLINER, *Differential Equations of Mathematical Physics*, Amsterdam, North-Holland; Interscience, New York, 1964. MR 31 #1443.