A Method for the Approximation of Functions Defined by Formal Series Expansions in Orthogonal Polynomials*

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Abstract. An algorithm is described for numerically evaluating functions defined by formal (and possibly divergent) series as well as convergent series of orthogonal functions which are, apart from a factor, orthogonal polynomials. When the orthogonal functions are polynomials, the approximations are rational functions. The algorithm is similar in some respects to the method of Padé approximants. A rational approximation involving Tchebychev polynomials due to H. Machley and described by E. Kogbetliantz [1] is a special case of the algorithm. ■

1. Introduction. The solutions to many physical problems are obtained as expansions in (infinite) series of orthogonal functions. When the series are convergent they can in principle be approximated to any accuracy by truncating the series at the proper point. When the series are weakly convergent or divergent, the procedures for their numerical evaluation become rather ad hoc and a more general approach would be useful. In the following sections we describe an approximation to functions defined as infinite series of orthogonal functions which are (apart from a factor) orthogonal polynomials. The approximation takes the form of the ratio of two functions, the denominator function being a polynomial. When the orthogonal functions are polynomials, the approximating function is rational.

The derivation of the algorithm will be formal and no real proofs are given. Indication that the algorithm is at least sometimes valid is provided by the numerical examples in Section 6. While the derivations could probably be made rigorous in the case of absolutely convergent series, interesting cases occur with divergent series such as the examples of Section 6. Numerous other practical applications of the algorithm have been made in the past year at Los Alamos Scientific Laboratory. The success of these examples would justify an effort at finding a class of functions representable by divergent series for which the algorithm is applicable.

2. The Approximation. The approximations we shall discuss fall roughly into two problems. The first of these is the approximation problem, that is, given a function (or equivalently its expansion), find an easily calculated approximation to the function. The second problem is the summation of infinite series, that is, given an infinite series (which may not be convergent in the ordinary sense), assign a sum to the series which gives the same value as the ordinary summation method when the series is convergent. Clearly the two problems cannot be sharply separated. In this paper however we will lean more toward the second problem.

For simplicity we assume the functions we will encounter are real on the real

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interval $[a, b]$. The generalization to complex functions on the real interval $[a, b]$ is trivial and the final results are identical to those which will be given.

The quantity to be evaluated is given in the form,

$$(1) \quad f(x) = \sum_{i=0}^{\infty} c_i g_i(x).$$

The basis functions $\{g_i\}$ are assumed to be orthogonal on $[a, b]$ with weight function $\omega(x)$ and normalization μ_i and, apart from a factor, to be orthogonal polynomials,

$$(2) \quad g_n(x) = g_0(x) q_n(x),$$

where $\{q_n\}$ are orthogonal polynomials of exact degree n with weight function $g_0^2 \omega$.

We introduce a set of complementary polynomials $\{p_i\}$ orthogonal on $[a, b]$ with weight function ω . The approximation to f is then given by

$$(3) \quad f(x) = \sum_{m=0}^M a_m g_m(x) \bigg/ \sum_{n=0}^N b_n p_n(x) = g_0(x) \sum_{m=0}^M a_m q_m(x) \bigg/ \sum_{n=0}^N b_n p_n(x),$$

where the sets of coefficients $\{a_i\}$ and $\{b_i\}$ must satisfy

$$(4) \quad f(x) \sum_{n=0}^N b_n p_n(x) - \sum_{m=0}^M a_m g_m(x) = \sum_{m=M+N+1}^{\infty} \rho_m g_m(x).$$

We take $b_0 = 1$ and call this the standard normalization. Other choices of normalization of the coefficients may at times be more convenient. If we were to use the algorithm to compute the inverse of f as a series of polynomials $\{p_i\}$, then $a_0 = 1$ might be the most appropriate normalization.

To solve the set of Eq. (4) it is convenient to introduce a matrix H^f associated with the function f whose element H_{mn}^f is the coefficient of $g_m(x)$ in the expansion of $f(x)p_n(x)$ in terms of the set $\{g_i\}$. An integral representation of the elements H^f is given by

$$(5) \quad H_{mn}^f = \mu_m^{-1} \int_a^b f(x) g_m(x) p_n(x) \omega(x) dx.$$

The first column of H^f is proportional to the expansion coefficients of $f(x)$,

$$(6) \quad H_{m0}^f = p_0 c_m,$$

and subsequent columns can be generated from the first by recursion. To see this we note the following properties of $\{g_i\}$.

(i) *Orthogonality.*

$$(7) \quad \int_a^b g_m(x) g_n(x) \omega(x) dx = \mu_m \delta_{mn},$$

where $\delta_{mn} = 1$ if $m = n$ and $\delta_{mn} = 0$ otherwise.

(ii) *Closure.*

$$(8) \quad f(y) = \sum_{n=0}^{\infty} \mu_n^{-1} g_n(y) \int_a^b f(x) g_n(x) \omega(x) dx.$$

(iii) *Linear Recursion.*

$$(9) \quad g_{n+1}(x) = (\alpha_n x + \beta_n)g_n(x) - \gamma_n g_{n-1}(x),$$

which remains valid for $n = 0$ when $g_{-1}(x)$ is replaced by 0. The complementary polynomials $\{p_i\}$ satisfy (7) and (8) with normalization h_i and

$$(10) \quad p_{n+1}(x) = (A_n x + B_n)p_n(x) - C_n p_{n-1}(x),$$

again valid for $n = 0$ when $p_{-1}(x)$ is replaced by zero.

Note that if we write

$$(11) \quad p_n(x) = k_n x^n + k'_n x^{n-1} + \dots, \quad p_0(x) = k_0,$$

then

$$(12) \quad \begin{aligned} A_n &= k_{n+1}/k_n, & B_n &= A_n(k'_{n+1}/k_{n+1} - k'_n/k_n), \\ C_n &= A_n h_n / (A_{n-1} h_{n-1}) = k_{n+1} k_{n-1} h_n / (k_n^2 h_{n-1}). \end{aligned}$$

Similar relations hold for the $\{g_m\}$. From the recursion relations (9) and (10) for the $\{g_m\}$ and the $\{p_n\}$ we see the elements of H^f satisfy the recursion relation

$$(13) \quad \left[\frac{B_n}{A_n} - \frac{\beta_m}{\alpha_m} \right] H_{mn}^f + \frac{\gamma_{m+1}}{\alpha_{m+1}} H_{m+1,n}^f + \frac{1}{\alpha_{m-1}} H_{m-1,n}^f = \frac{1}{A_n} H_{m,n+1}^f + \frac{C_n}{A_n} H_{m,n-1}^f.$$

Since this relation remains valid for m or n equal to zero provided we set $H_{m,-1}^f = H_{-1,n}^f = 0$ in Eq. (13), any segment of the matrix can be generated by recursion from the first column or from the expansion coefficients of f .

Finally for later use we note that the elements of the inverse of H^f are given by

$$(14) \quad (H^f)_{mn}^{-1} = \mu_n h_m^{-1} H_{nm}^{(1/f)} = h_m^{-1} \int_a^b f^{-1}(x) g_n(x) p_m(x) \omega(x) dx,$$

as can be seen by multiplication by (5) and use of the closure relations (8).

We now write the exact relation

$$(15) \quad B(x)f(x) - A(x) = r(x),$$

where $r(x)$ is a remainder and is related to the error in the approximation. Multiply Eq. (15) by $\mu_m^{-1} g_m(x) \omega(x) dx$ and integrating over $[a, b]$ yields the set of equations

$$(16) \quad \sum_{n=0}^N H_{mn}^f b_n - a_m = \rho_m,$$

with

$$(17) \quad r(x) = \sum_{m=0}^{\infty} \rho_m g_m(x).$$

We now require ρ_m to be zero for $m \leq M + N$. We note that such a requirement is the basis of the Padé approximation to a power series. The equations to be solved for the coefficients of A and B are

$$(18) \quad \sum_{n=1}^N H_{mn}^f b_n = -H_{m0}^f = -k_0 c_m, \quad M + 1 \leq m \leq M + N,$$

and

$$(19) \quad a_m = \sum_{n=0}^N H_{mn}^f b_n, \quad 0 \leq m \leq M.$$

The coefficients of $r(x)$ are given by

$$(20) \quad \rho_m = \sum_{n=0}^N H_{mn}^f b_n, \quad M + N + 1 \leq m.$$

We define a "Padé-like" table as a two-dimensional array of elements, the $[N, M]$ element being the approximation with numerator of order M and denominator of order (or degree) N . It can easily be shown (by expansion of the determinants by minors) that the $[N, M]$ approximation can be written as the ratio of determinants

$$(21) \quad [N, M](x) = \frac{\begin{vmatrix} \sum_{m=0}^M H_{m0}^f g_m(x) & \cdots & \sum_{m=0}^M H_{mn}^f g_m(x) & \cdots & \sum_{m=0}^M H_{m,N}^f g_m(x) \\ H_{M+1,0}^f & \cdots & H_{M+1,n}^f & \cdots & H_{M+1,N}^f \\ \vdots & & \vdots & & \vdots \\ \vdots & & \vdots & & \vdots \\ H_{M+N,0}^f & \cdots & H_{M+N,n}^f & \cdots & H_{M+N,N}^f \end{vmatrix}}{\begin{vmatrix} p_0(x) & \cdots & p_n(x) & \cdots & p_N(x) \\ H_{M+1,0}^f & \cdots & H_{M+1,n}^f & \cdots & H_{M+1,N}^f \\ \vdots & & \vdots & & \vdots \\ \vdots & & \vdots & & \vdots \\ H_{M+N,0}^f & \cdots & H_{M+N,n}^f & \cdots & H_{M+N,N}^f \end{vmatrix}}.$$

Note that this expression uses a nonstandard normalization. In a similar manner the remainder $r(x)/B(x)$ can formally be expressed as the ratio of the two determinants above except that the sums over the index m in the numerator are over all m such that $m > M + N$.

When $M \geq N$, the approximation can be written in two forms. The first form is the improper fraction,

$$(22) \quad \frac{A(x)}{B(x)} = \frac{\sum_{m=0}^{N+L} a_m g_m(x)}{\sum_{n=0}^N b_n p_n(x)},$$

where the coefficients are given in Eq. (18) and (19) with $M = N + L$. The second form is the polynomial plus proper fraction

$$(23) \quad \frac{A(x)}{B(x)} = D(x) + \frac{\bar{A}(x)}{B(x)} = \sum_{l=0}^L d_l g_l(x) + \frac{\sum_{m=0}^{N-1} \bar{a}_m g_m(x)}{\sum_{n=0}^N b_n p_n(x)}.$$

The coefficients b_n are obtained in both cases by solving Eq. (18) with $M = N + L$.

If we expand the quantity

$$(24) \quad B(x)f(x) - B(x)D(x) - \bar{A}(x) = r(x),$$

we find the coefficients $\bar{a}_m$ are given by

$$(25) \quad \bar{a}_m = \sum_{n=0}^N H_{mn}^f b_n - \sum_{l=0}^L H_m^B d_l, \quad 0 \leq m \leq N - 1,$$

and the coefficients d_l are given by solving

$$(26) \quad \sum_{l=0}^L H_m^B d_l = \sum_{n=0}^N H_{mn}^f b_n, \quad N \leq m \leq N + L.$$

As the notation indicates, H^B is the matrix associated with the polynomial $B(x)$. The Eqs. (26) are easily solved due to the simple structure of H^B . (The submatrix involved in Eq. (26) is triangular.)

The coefficients for the two forms (22) and (23) are related by

$$(27) \quad d_L = a_{N+L}/H_{N+L,L}^B,$$

$$d_{L-l} = \left[a_{N+L-l} - \sum_{m=0}^{l-1} d_{L-m} H_{N+L-l,L-m}^B \right] / H_{N+L-l,L-l}^B, \quad 1 \leq l \leq L,$$

and

$$(28) \quad \bar{a}_m = a_m - \sum_{l=0}^L H_{m,L-l}^B d_{L-l}, \quad 0 \leq m \leq N-1.$$

These latter equations contain no direct reference to f , and indeed they provide a way of calculating the proper fraction plus finite sum from any improper fraction. If the fraction is rational, they provide a method for applying the Euclidian algorithm to find common factors without first transforming the (finite) series of orthogonal polynomials to a power series.

3. Relation to a Least Squares Approximation. One measure of the weighted error of the rational approximation to the function $f(x)$ on the interval $[a, b]$ is the quantity

$$(29) \quad e = \int_a^b [B^*(x)f(x) - A^*(x)]^2 \omega(x) dx.$$

(If the functions were complex, the squared quantity would be replaced by the squared magnitude.) This error is a function of the coefficients $\{a_i^*\}$ and $\{b_i^*\}$. For the moment consider $B^*(x)$ to be determined and $A^*(x)$ to be a polynomial of degree M . Then e is a quadratic function of the coefficients $\{a_i^*\}$. This error is minimized when for each $k \leq M$,

$$\frac{\partial e}{\partial a_k^*} = -2 \int_a^b [B^*(x)f(x) - A^*(x)] g_k(x) \omega(x) dx = 0,$$

or

$$(30) \quad a_k^* = \sum_{n=0}^N H_{kn}^f b_n^*, \quad 0 \leq k \leq M.$$

Thus the set $\{a_k^*\}$ which minimizes the error e is just that set given by the algorithm once the set $\{b_i^*\}$ is determined.

Now consider the minimization e with respect to the set of coefficients $\{b_i^*\}$. For each $k \leq N$,

$$\frac{\partial e}{\partial b_k^*} = 2 \int_a^b [B^*(x)f(x) - A^*(x)] f(x) p_k(x) \omega(x) dx = 0,$$

or

$$(31) \quad \sum_{l=0}^{\infty} \left[\sum_{n=0}^N H'_{ln} b_n^* - a_l^* \right] \mu_l H'_{lk} = 0, \quad 0 \leq k \leq N,$$

where $a_l^* \equiv 0$ for $l > M$.

This set of equations is to be solved for the b_k^* 's consistently with the set of equations (30) and a subsidiary condition involving the normalization for the a_k^* 's to give the least squares approximation to $f(x)$. Although in principle the coefficients can be obtained from H' , in practice this is not always possible so the coefficients in Eq. (31) cannot be regarded as known.

To obtain the equations of the algorithm, we hold M fixed and let $N \rightarrow \infty$. The set of Eqs. (31) becomes

$$(32) \quad \sum_{l=0}^{\infty} \left[\sum_{n=0}^{\infty} H'_{ln} b_n^* \right] \mu_l H'_{lk} = \sum_{l=0}^M a_l^* \mu_l H'_{lk}, \quad 0 \leq k < \infty.$$

Now we can multiply by the inverse of H' (14) and obtain,

$$(33) \quad \begin{aligned} \sum_{n=0}^{\infty} H'_{ln} b_n^* &= a_l^*, & 0 \leq l \leq M, \\ \sum_{n=0}^{\infty} H'_{ln} b_n^* &= 0, & M+1 \leq l. \end{aligned}$$

If we now truncate the set of Eqs. (33) we arrive at the equations of the algorithm. It is then clear that the algorithm does not yield a least squares approximation except possibly in the limit $N \rightarrow \infty$.

4. Some Remarks on Convergence. In this section we will examine the convergence of the first two rows of the Padé-like table under suitable restrictions on f . We note that the limit

$$(34) \quad \lim_{m \rightarrow \infty} \left[\frac{1}{\alpha_{m-1}} + \frac{\gamma_{m+1}}{\alpha_{m+1}} - \frac{\beta_m}{\alpha_m} \right]$$

is finite. This condition is satisfied by all of the classical orthogonal polynomials. We restrict f , defined by

$$(35) \quad f(x) = \sum_{m=0}^{\infty} c_m g_m(x),$$

to be such that

$$(36) \quad c_m \sim N m^{\nu}.$$

We recall the following transformation for improving the convergence of the series (35). By formally rearranging the series and using the recursion relation for $\{g_i\}$ we can write

$$(37) \quad f = (x - a)^{-1} \sum_{m=0}^{\infty} c'_m g_m(x),$$

where

$$(38) \quad c'_m = \frac{c_{m-1}}{\alpha_{m-1}} + \frac{\gamma_{m+1} c_{m+1}}{\alpha_{m+1}} - \left[\frac{\beta_m}{\alpha_m} + a \right] c_m.$$

If a is chosen such that

$$(39) \quad \lim_{m \rightarrow \infty} \left[\frac{1}{\alpha_{m-1}} + \frac{\gamma_{m+1}}{\alpha_{m+1}} - \frac{\beta_m}{\alpha_m} \right] = a,$$

then

$$(40) \quad c_m' \sim N' m^{r-1},$$

and the series (37) is (asymptotically) more rapidly convergent. We shall return to this result shortly.

We note that the elements $[0, M]$ of the first row of the Padé-like table are simply the truncated sums

$$(41) \quad [0, M](x) = \sum_{m=0}^M a_m g_m(x) = \sum_{m=0}^M c_m g_m(x).$$

If the sum (35) is convergent to f , then the sequence of elements of the first row of the table are convergent to f .

Now consider the elements of the second row of the table. Notice that

$$(42) \quad \begin{aligned} H_{M+1,1}/H_{M+1,0} &= A_0 \left[\frac{B_0}{A_0} - \frac{\beta_{M+1}}{\alpha_{M+1}} + \frac{\gamma_{M+2}}{\alpha_{M+2}} \frac{c_{M+2}}{c_{M+1}} + \frac{1}{\alpha_M} \frac{c_M}{c_{M+1}} \right] \\ &\sim A_0 \left[\frac{B_0}{A_0} + \frac{1}{\alpha_M} + \frac{\gamma_{M+2}}{\alpha_{M+2}} - \frac{\beta_{M+1}}{\alpha_{M+1}} \right] + \mathcal{O}(M^{-1}). \end{aligned}$$

For large M , we use (12) and (13) to write $B(x)$ in a form which displays the position of the pole in the approximation,

$$(43) \quad B_1(x) \sim -K \left[\frac{1}{\alpha_{M-1}} + \frac{\gamma_{M+1}}{\alpha_{M+1}} - \frac{\beta_M}{\alpha_M} - x \right],$$

where K depends on the normalization. As $M \rightarrow \infty$, the pole moves to a position identical to that in Eqs. (37) and (39).

The numerator $A_M(x)$ is given by

$$(44) \quad A_M(x) = \sum_{m=0}^M a_m g_m(x),$$

where for large M ,

$$(45) \quad a_m = K \left\{ \frac{c_{m-1}}{\alpha_{m-1}} + \frac{\gamma_{m+1} c_{m+1}}{\alpha_{m+1}} - \left[\frac{\beta_m}{\alpha_m} - \frac{1}{\alpha_M} - \frac{\gamma_{M+2}}{\alpha_{M+2}} + \frac{\beta_{M+2}}{\alpha_{M+1}} \right] c_m \right\}.$$

In the limit $M \rightarrow \infty$, (with $K = 1$) this is identical to (38) and (39) and in this limit the approximation $[1, M]$ is identical to the transformation mentioned at the beginning of this section. It is tempting to conjecture that as one moves down the table, the row-wise convergence improves. Numerical evidence indicates that this is the case, but proof remains to be found.

We conclude this section with a formula for the remainder $R_N^M(x)$,

$$(46) \quad R_N^M(x) = f(x) - A_M(x)/B_N(x),$$

when $\{g_i\} = \{P_n^{(\alpha, \beta)}(x)\}$, the set of Jacobi polynomials. Referring to Szegő [3, Eqs. (4.62.19) and (9.2.1)], we find

$$\begin{aligned}
 & \sum_{n=N+1}^{\infty} [h_n^{(\alpha, \beta)}]^{-1} P_n^{(\alpha, \beta)}(x) Q_n^{(\alpha, \beta)}(y) \\
 (47) \quad &= \frac{2^{-\alpha-\beta} \Gamma(N+2) \Gamma(N+\alpha+\beta+2)}{(2N+\alpha+\beta+2) \Gamma(N+\alpha+1) \Gamma(N+\beta+1)} \\
 & \cdot \left[\frac{P_{N+1}^{(\alpha, \beta)}(x) Q_N^{(\alpha, \beta)}(y) - P_N^{(\alpha, \beta)}(x) Q_{N+1}^{(\alpha, \beta)}(y)}{x-y} \right].
 \end{aligned}$$

The expansion coefficients of $B_N R_N^M$ are zero when their index is less than $M+N$ and equal to those of $B_N f$ otherwise. It then follows that

$$\begin{aligned}
 (48) \quad R_N^M(x) &= \frac{\Gamma(N+M+\alpha) \Gamma(N+M+\alpha+\beta+2) 2^{1-\alpha-\beta} e^{i\pi\alpha}}{(2N+2M+\alpha+\beta+2) \Gamma(N+\alpha+1) \Gamma(N+\beta+2) B_N(x)} \\
 & \times \frac{1}{2\pi i} \oint_C \left[\frac{P_{N+M+1}^{(\alpha, \beta)}(x) Q_{N+M}^{(\alpha, \beta)}(y) - P_{N+M}^{(\alpha, \beta)}(x) Q_{N+M+1}^{(\alpha, \beta)}(y)}{x-y} \right] B f \omega dy,
 \end{aligned}$$

where $e^{i\pi\alpha}(y) = (y-1)^\alpha(y+1)^\beta$ and C is an ellipse with foci at ± 1 and enclosing no singularity of f . With special values for α and β , (48) gives representations for the remainder in approximations involving Gegenbauer, Legendre, and Tchebychev polynomials.

5. Numerical Examples. We now consider the application of the algorithm to a particular case, though some of the remarks made later are based on experience with other functions. Consider the function

$$(49) \quad -i\eta \left(\frac{1-x}{2} \right)^{-1-i\eta} = \sum_{l=0}^{\infty} (2l+1) \frac{(1+i\eta)_l}{(1-i\eta)_l} P_l(x),$$

where $\{g_l\} = \{P_l\}$, the set of Legendre polynomials. The symbol $(a)_l$ in Eq. (49) is defined by

$$(50) \quad (a)_0 = 1, \quad (a)_l = a(a+1) \cdots (a+l-1).$$

In particular we consider Eq. (49) with $\eta = 1$.

The function in Eq. (49) when multiplied by the proper factor is the Coulomb scattering amplitude which one obtains in the nonrelativistic scattering of charged particles. The coefficients of the Legendre polynomials behave asymptotically as $l^{1+2i\eta}$ as $l \rightarrow \infty$. The series is divergent, but nevertheless represents the function on the left [4].

Referring to the Padé-like table, we will consider the three-diagonal approximations [4, 4], [8, 8], and [12, 12]. These three approximations were calculated on a LASL-Control Data 6600 computer. Single-precision floating point numbers are represented on this machine by a 48 bit mantissa (about 14 decimal digits). In Fig. 1 we have plotted the magnitude of the relative error. From Fig. 1 (and other calculations) several things are immediately apparent. As we move down the diagonal of the Padé-like table (or more generally as $M+N$ increases), the relative error is reduced over the interval $[-1, 1]$. In these numerical calculations, the minimum error is of course bounded by the accuracy of the representation of the number (in this case $\sim 10^{-14}$). Since this function has a branch point at $x = 1$, and sufficiently close to the branch point no finite configuration of poles and zeroes can "mock up"

the branch point, the approximation fails close to $x = 1$.

FIGURE 1

If we calculate the positions of the poles and zeroes in the complex plane for various finite approximations we find that they tend to lie along the branch cut from $x = 1$ to infinity and tend to accumulate at $x = 1$. We show this behavior in Fig. 2, where we have plotted their positions for the $[12, 12]$ approximation. If in Eq. (49) we allow η to vary, we find that as $\eta \rightarrow 0$, the poles and zeroes near the cut move closer to the real axis. Off the real axis the error in the approximation is somewhat larger than that indicated in Fig. 1. To the left of $x = 1$ and along the lines with the imaginary part of x equal to $\pm 1/2$, the magnitude of the relative error is about 10^{-8} and 10^{-9} for the $[8, 8]$ and $[12, 12]$ approximations and about 10^{-4} for the $[4, 4]$ approximation. Thus the rational approximation is a useful analytic continuation into the complex plane.

In practice it is found that in addition to poles along the branch cut of the function of Eq. (49), poles appear elsewhere in the complex plane in positions not related to the analytic structure of the approximated function. The positions of these poles vary rapidly from an approximation of one order to the next, and (if allowed by the order of M) they are usually cancelled by zeroes of the numerator (the numerator and denominator possess common factors). This cancelling by the

zeroes might be expected from the least-squares property of the numerator. We will call these spurious poles "nuisance poles." A similar phenomenon has been observed in Padé approximants [5, 6]. Although the form of the approximation is such that it is possible but not convenient to use the Euclidian algorithm for removing these common factors, the nuisance poles need cause no real problem. If a nuisance pole of the $[N, M]$ approximation should occur too close to a region of interest, then this will usually not be so for the $[N \pm 1, M]$ or $[N, M \pm 1]$ or one of the other adjacent approximations.

FIGURE 2

It is interesting to compare the rational approximation obtained from series of orthogonal polynomials with the Padé approximants obtained from a power series. The power series expansion from which the Padé approximant to a function is constructed depends on the value (and derivatives) of the function at a single point and converges most rapidly near this point [1]. The polynomial expansion depends on weighted averages of the function over the interval of orthogonality and (when convergent) converges about equally rapidly over the entire interval. These properties are reflected in the rational approximations constructed from the polynomial expansions. To illustrate these points we consider the function $-\log((1-x)/2)$, whose power series expansion about $x = 0$ is given by

$$(51) \quad -\log((1-x)/2) = \log(2) + \sum_{n=1}^{\infty} x^n/n, \quad |x| < 1,$$

and whose Legendre polynomial expansion in x is given by

$$(52) \quad -\log((1-x)/2) = 1 + \sum_{n=1}^{\infty} (2n-1)/(n(n+1))P_n(x), \quad -1 \leq x < 1.$$

The expansion coefficients in both cases have a similar asymptotic behavior.

We have calculated the $[4, 4]$ Padé approximant from (51) and the $[4, 4]$ rational

polynomial approximant from (52). In Fig. 3 we compare the magnitudes of the relative errors. As previously mentioned, we see that the $[4, 4]$ Padé approximant has a very small error near $x = 0$ with the error rising slowly away from $x = 0$. By contrast, the $[4, 4]$ rational polynomial approximation fits the function with about the same (average) accuracy over the interval $(-1, 1)$ with eight zeroes of the (relative) error function lying in the interval $(-1, 1)$.

FIGURE 3

One more point should be made. If the expansion (52) is truncated at $n = 8$, rearranged in powers of x , and the $[4, 4]$ Padé approximant constructed from this truncated series, the result must be identical to the $[4, 4]$ rational polynomial approximation from the uniqueness of the rational form.** The coefficients from which the Padé approximants are constructed are a function of the truncation point and it is not often that the infinite series (52) can be rearranged into a power series (51).

We will now consider in detail two other closely related examples. These examples are limiting cases of the expansion of $(1 - x)^\alpha$ as $\alpha \rightarrow -1$ [4]. In the previous example the indications (admittedly based on scanty evidence) were that the diagonal sequence of approximations was converging. It may happen that no diagonal sequence of approximations exists.

Consider first the rational approximations to the function expanded in a (divergent) series of Legendre polynomials $\{P_l\}$,

$$(53) \quad (x - 1)^{-1} = \sum_{l=0}^{\infty} (2l + 1)\psi(2l + 1)P_l(x),$$

where ψ is the logarithmic derivative of the gamma function. For our purposes it is sufficient to know that ψ satisfies the recursion relation

** I wish to thank G. A. Baker, Jr. for a discussion on this point.

$$(54) \quad \psi(l+1) = \psi(l) + 1/l.$$

It can be shown by induction using the recursion relations (14) and by the symmetry of H^f that

$$(55) \quad \begin{aligned} H_{mn}^f &= (2m+1)\psi(m+1), & m \geq n, \\ &= (2m+1)\psi(n+1), & m \leq n. \end{aligned}$$

Referring then to the determinantal form (21) of the rational approximation it is easily seen that

$$(56) \quad \begin{aligned} \text{(i)} \quad [0, M](x) &= \sum_{m=0}^M (2m+1)\psi(m+1)P_m(x), \\ \text{(ii)} \quad [1, M](x) &= 1/(x-1), & 0 \leq M, \\ \text{(iii)} \quad [N, 0](x) &= 1/(x-1), & 1 \leq N, \\ \text{(iv)} \quad [N, M](x) &\text{ is indeterminant for } M > 0 \text{ and } N > 1. \end{aligned}$$

Since the original series was divergent, the sequence of elements in the first row is not convergent. The elements of the second row and those of the first column are (obviously) convergent to the function $(x-1)^{-1}$, and no other row or column nor any diagonal sequence even exists.

A similar result obtains for the generalized function

$$(57) \quad \delta(1-x) = \sum_{l=0}^{\infty} (l+1/2)P_l(x),$$

where the $\{P_i\}$ are again Legendre polynomials and δ is the Dirac delta function. The elements of the related matrix H^f are given by

$$(58) \quad H_{mn}^f = (m+1/2).$$

The elements of the Padé-like table are then given by

$$(59) \quad \begin{aligned} [0, M](x) &= \sum_{m=0}^M (m+\tfrac{1}{2})P_m(x), & 0 \leq M, \\ [1, M](x) &= 0/(x-1), & 0 \leq M, \\ [N, 0](x) &= 0/(x-1), & 1 \leq N, \\ [N, M](x) &\text{ is indeterminant for } M > 0 \text{ and } N > 1. \end{aligned}$$

The remarks after the previous example are applicable also in this case.

6. Conclusion. Many additional properties of the approximation can be derived for particular choices of the set $\{g_i\}$. These properties may depend on special properties of the set under consideration and so are not generalizable to other sets of orthogonal functions. These special properties may be very important in any given problem.

As an example, suppose the interval of orthogonality is even, $[-a, a]$ and the functions $g_n(x)$ and $p_n(x)$ are even or odd according as the index n is even or odd. Then if a function is defined by an even (odd) series and an odd-odd (even-even)

approximation is attempted, the Eqs. (18) have no solution. Also the even-odd or odd-even approximations are equal to adjacent nonvanishing elements in the table. These remarks may be verified by referring to the determinantal form of the approximation (21) and noting that alternate elements of H^f vanish. In these even or odd cases, some computing economy can be effected by rewriting the recursion relation (12) for H^f so that it relates the elements $H_{m,n}^f$, $H_{m\pm 2,n}^f$, $H_{m,n\pm 2}^f$, and reformulating the algorithm so that only the even (odd) elements of $\{g_i\}$ are involved (we take $B(x)$ to have only even terms in either case).

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1. E. G. KOGBETLIANTZ, "Generation of elementary functions," in *Mathematical Methods for Digital Computers*, Wiley, New York, 1960. MR 22 #8681.
2. A. ERDÉLYI, et al., *Higher Transcendental Functions*, Vol. II, McGraw-Hill, New York, 1953. MR 15, 419.
3. G. SZEGÖ, *Orthogonal Polynomials*, Amer. Math. Soc. Colloq. Publ., Vol. 23, Amer. Math. Soc., Providence, R. I., 1959. MR 1, 14.
4. J. HOLDEMAN, "Legendre polynomial expansions of hypergeometric functions with applications." (To appear.)
5. G. A. BAKER, JR., J. L. GAMMEL & J. G. WILLS, "An investigation of the applicability of the Padé approximant method," *J. Math. Anal. Appl.*, v. 2, 1961, pp. 405-418. MR 23 #B3125.
6. G. A. BAKER, JR., "The theory and application of the Padé approximant method," in *Advances in Theoretical Physics*, Vol. I, Academic Press, New York, 1965, pp. 1-58. MR 32 #5253.