

## On Some Orthogonal Polynomial Integrals

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**Abstract.** The modified moments of the weight functions  $w(x) = x^\rho(1-x)^{\alpha}\ln(1/x)$ , on  $[0, 1]$ , with respect to the shifted Jacobi polynomials  $P_n^{(\alpha,\beta)}(x) = P_n^{(\alpha,\beta)}(2x-1)$ , and  $w_p(x) = x^\rho e^{-x}(\ln x)^p$ ,  $p = 1, 2$ , on  $[0, \infty)$ , with respect to the generalized Laguerre polynomials  $L_n^{(\alpha)}(x)$ , are explicitly evaluated.

**1. A Jacobi Polynomial Integral.** In a recent paper, Gautschi [3], generalizing a result of Blue [2], has considered and explicitly evaluated the modified moments of the weight function

$$w(x) = x^\mu \ln(1/x), \quad \mu > -1,$$

on  $[0, 1]$ , with respect to the shifted Legendre polynomials  $P_n^*(x) = P_n(2x-1)$ .

We further generalize these results by considering the weight function

$$(1.1) \quad w(x) = x^\rho(1-x)^\alpha \ln(1/x), \quad \alpha, \rho > -1,$$

and evaluating its modified moments on  $[0, 1]$  with respect to the shifted Jacobi polynomials  $P_n^{*(\alpha,\beta)}(x) = P_n^{(\alpha,\beta)}(2x-1)$ .

It is convenient from now on to replace  $\rho$  by  $\beta + \mu$ ; thus, the modified moments we have to examine assume the form

$$(1.2) \quad \nu_n^{(\alpha,\beta)}(\mu) = \int_0^1 x^{\beta+\mu}(1-x)^\alpha \ln(1/x) P_n^{*(\alpha,\beta)}(x) dx,$$

$$\alpha, \beta, \beta + \mu > -1, n = 0, 1, 2, \dots$$

We easily see that

$$\begin{aligned} \nu_n^{(\alpha,\beta)}(\mu) &= -2^{-(\alpha+\beta+\mu+1)} \int_{-1}^1 (1-t)^\alpha (1+t)^{\beta+\mu} \ln(\tfrac{1}{2}(1+t)) P_n^{(\alpha,\beta)}(t) dt \\ &= -2^{-(\alpha+\beta+\mu+1)} \left\{ \int_{-1}^1 (1-t)^\alpha (1+t)^{\beta+\mu} \ln(1+t) P_n^{(\alpha,\beta)}(t) dt \right. \\ &\quad \left. - \ln 2 \cdot \int_{-1}^1 (1-t)^\alpha (1+t)^{\beta+\mu} P_n^{(\alpha,\beta)}(t) dt \right\}, \end{aligned}$$

hence, by putting

$$(1.3) \quad I_n^{(\alpha,\beta)}(\mu) = \int_{-1}^1 (1-t)^\alpha (1+t)^{\beta+\mu} P_n^{(\alpha,\beta)}(t) dt,$$

$$\alpha, \beta, \beta + \mu > -1, n = 0, 1, 2, \dots,$$

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we obtain

$$(1.4) \quad \nu_n^{(\alpha, \beta)}(\mu) = 2^{-(\alpha + \beta + \mu + 1)} \left\{ I_n^{(\alpha, \beta)}(\mu) \ln 2 - \frac{d}{d\mu} I_n^{(\alpha, \beta)}(\mu) \right\}.$$

The following expression for (1.3),

$$(1.5) \quad I_n^{(\alpha, \beta)}(\mu) = 2^{\alpha + \beta + \mu + 1} \frac{\Gamma(\mu + 1)}{n! \Gamma(\mu - n + 1)} \frac{\Gamma(\beta + \mu + 1) \Gamma(n + \alpha + 1)}{\Gamma(n + \alpha + \beta + \mu + 2)},$$

is known ([1], [4, p. 256]). Indeed, (1.5) is easily obtained, multiplying on both sides of Rodrigues' formula,

$$(1 - t)^\alpha (1 + t)^\beta P_n^{(\alpha, \beta)}(t) = \frac{(-1)^n}{2^n n!} \frac{d^n}{dt^n} \{ (1 - t)^{n + \alpha} (1 + t)^{n + \beta} \},$$

by  $(1 + t)^\mu$ , integrating from  $-1$  to  $1$  and carrying out  $n$  partial integrations on the right-hand side.

Differentiating (1.5) with respect to  $\mu$  gives

$$(1.6) \quad \frac{d}{d\mu} I_n^{(\alpha, \beta)}(\mu) = I_n^{(\alpha, \beta)}(\mu) \{ \ln 2 + \psi(\mu + 1) + \psi(\beta + \mu + 1) \\ - \psi(\mu - n + 1) - \psi(n + \alpha + \beta + \mu + 2) \},$$

where  $\psi(x) = \Gamma'(x)/\Gamma(x)$  is the logarithmic derivative of the gamma function, and, if  $\mu$  coincides with an integer  $m < n$ ,  $m \geq 0$ , the right-hand member must be replaced by its limit as  $\mu \rightarrow m$ .

We first consider the case where  $\mu \neq 0, 1, 2, \dots, n - 1$ , whenever  $n \geq 1$ . By inserting (1.5) and (1.6) in (1.4), we obtain

$$(1.7) \quad \nu_n^{(\alpha, \beta)}(\mu) = \frac{\Gamma(\mu + 1) \Gamma(\beta + \mu + 1) \Gamma(n + \alpha + 1)}{n! \Gamma(\mu - n + 1) \Gamma(n + \alpha + \beta + \mu + 2)} \\ \cdot \{ \psi(\mu - n + 1) + \psi(n + \alpha + \beta + \mu + 2) - \psi(\mu + 1) - \psi(\beta + \mu + 1) \},$$

with  $\alpha, \beta, \beta + \mu > -1$ ,  $n = 0, 1, 2, \dots$  and  $\mu \neq 0, 1, 2, \dots, n - 1$  if  $n \geq 1$ .

Taking into account the recurrence relations  $\Gamma(x + 1) = x\Gamma(x)$  and  $\psi(x + 1) = \psi(x) + 1/x$ , we may derive a useful algorithm for the computation of the modified moments  $\nu_n^{(\alpha, \beta)}(\mu)$ . Indeed, it is easily seen that, if we put

$$a_0^{(\alpha, \beta)}(\mu) = \frac{\Gamma(\alpha + 1) \Gamma(\beta + \mu + 1)}{\Gamma(\alpha + \beta + \mu + 2)}, \\ b_0^{(\alpha, \beta)}(\mu) = \psi(\alpha + \beta + \mu + 2) - \psi(\beta + \mu + 1),$$

and we construct the two sequences  $\{a_n^{(\alpha, \beta)}(\mu)\}$  and  $\{b_n^{(\alpha, \beta)}(\mu)\}$ , defined by the recurrence relationships

$$a_n^{(\alpha, \beta)}(\mu) = a_{n-1}^{(\alpha, \beta)}(\mu) \frac{(\alpha + n)(\mu - n + 1)}{n(\alpha + \beta + \mu + n + 1)},$$

$$b_n^{(\alpha, \beta)}(\mu) = b_{n-1}^{(\alpha, \beta)}(\mu) + \frac{1}{\alpha + \beta + \mu + 1 + n} - \frac{1}{\mu + 1 - n},$$

we have

$$\nu_n^{(\alpha, \beta)}(\mu) = a_n^{(\alpha, \beta)}(\mu) b_n^{(\alpha, \beta)}(\mu).$$

Therefore, this last expression also shows that (1.7) can be written in the following rational form with respect to  $n$

$$(1.8) \quad \nu_n^{(\alpha, \beta)}(\mu) = \frac{\Gamma(\alpha + 1)\Gamma(\beta + \mu + 1)}{\Gamma(\alpha + \beta + \mu + 2)} \left\{ \psi(\alpha + \beta + \mu + 2) - \psi(\beta + \mu + 1) \right. \\ \left. + \sum_{k=1}^n \left( \frac{1}{\alpha + \beta + \mu + 1 + k} - \frac{1}{\mu + 1 - k} \right) \right\} \\ \cdot \prod_{k=1}^n \frac{(\alpha + k)(\mu + 1 - k)}{k(\alpha + \beta + \mu + 1 + k)},$$

where  $\alpha$ ,  $\beta$ , and  $\mu$  satisfy the above-mentioned conditions.

To examine the remaining case  $n \geq 1$  and  $\mu = m = 0, 1, \dots, n-1$ , we recall that for any integer  $r \geq 0$ ,

$$\lim_{\epsilon \rightarrow 0} \frac{\psi(-r + \epsilon)}{\Gamma(-r + \epsilon)} = (-1)^{r-1} r!.$$

Then, from (1.7), we obtain

$$\nu_n^{(\alpha, \beta)}(m) = \lim_{\mu \rightarrow m} \nu_n^{(\alpha, \beta)}(\mu) \\ = \frac{\Gamma(n + \alpha + 1)\Gamma(m + 1)\Gamma(\beta + m + 1)}{n!\Gamma(n + \alpha + \beta + m + 2)} \lim_{\epsilon \rightarrow 0} \frac{\psi(m + \epsilon - n + 1)}{\Gamma(m + \epsilon - n + 1)},$$

and finally

$$(1.9) \quad \nu_n^{(\alpha, \beta)}(m) = (-1)^{n-m} \frac{m!(n-m-1)!}{n!} \frac{\Gamma(n + \alpha + 1)\Gamma(\beta + m + 1)}{\Gamma(n + \alpha + \beta + m + 2)}, \\ \alpha, \beta > -1, m = 0, 1, 2, \dots, n-1, n \geq 1.$$

This completes the evaluation of the integrals (1.2). Integrals of the form

$$\int_0^1 x^{\beta + \mu} (1-x)^{\alpha} (\ln(1/x))^p P_n^{*(\alpha, \beta)}(x) dx,$$

may be similarly evaluated by repeatedly differentiating (1.7) with respect to  $\mu$ .

**2. Some Examples.** The results derived in the previous section show that if one has to evaluate modified moments of a given weight function of type (1.1) for given values of  $\rho$  and  $\alpha$ , then one may choose as polynomial basis the Jacobi polynomials

$P_n^{(\alpha, \beta)}(2x - 1)$ , with  $\beta$  being a free parameter. For instance, in the case of the weight function

$$w(x) = x^\rho \ln(1/x), \quad \rho > -1,$$

we can construct the modified moments associated with the basis  $\{P_n^{(0, \beta)}(2x - 1)\}$  instead of the particular one,  $\{P_n^{(0, 0)}(2x - 1)\}$  considered by Gautschi [3].

It may be of some interest to note that the choice  $\rho = \beta$  yields very simple expressions for the corresponding modified moments,

$$(2.1) \quad \nu_n^{(0, \beta)}(0) = \int_0^1 x^\beta \ln(1/x) P_n^{(0, \beta)}(2x - 1) dx = \begin{cases} 1/(\beta + 1)^2, & n = 0, \\ \frac{(-1)^n \Gamma(\beta + 1)(n - 1)!}{\Gamma(n + \beta + 2)}, & n \geq 1. \end{cases}$$

Also, in the case of the more general weight functions (1.1), the formulas we obtain are particularly simple when we let  $\rho = \beta$ ,

$$(2.2) \quad \int_0^1 x^\beta (1 - x)^\alpha \ln(1/x) P_n^{(\alpha, \beta)}(2x - 1) dx = \begin{cases} \frac{\Gamma(\alpha + 1)\Gamma(\beta + 1)}{\Gamma(\alpha + \beta + 2)} \{ \psi(\alpha + \beta + 2) - \psi(\beta + 1) \}, & n = 0, \\ (-1)^n \frac{\Gamma(n + \alpha + 1)\Gamma(\beta + 1)}{n\Gamma(n + \alpha + \beta + 2)}, & n \geq 1. \end{cases}$$

An example of (1.1), with  $\alpha \neq 0$ , could be the weight function

$$w(x) = x^\rho (1 - x)^{-1/2} \ln(1/x), \quad \rho > -1,$$

for which, recalling that [4, p. 60]

$$P_n^{*(-1/2, -1/2)}(x) = T_n^*(x) \prod_{k=1}^n \frac{2k - 1}{2k},$$

where  $T_n^*(x) = T_n(2x - 1)$  is the shifted Chebyshev polynomial of degree  $n$ . Setting

$$\tau_n(\rho) = \int_0^1 x^\rho (1 - x)^{-1/2} \ln(1/x) T_n^*(x) dx, \quad \rho > -1,$$

and applying (1.9) and (1.8), we have

$$(2.3) \quad \tau_n(\rho) = \begin{cases} (-1)^{n-m} \frac{m!(n-m-1)!}{(n+m)!} \pi \prod_{k=1}^m \frac{2k-1}{2}, & \rho + 1/2 = m < n, \quad m \geq 0 \text{ an integer,} \\ \frac{\sqrt{\pi}\Gamma(\rho+1)}{\Gamma(\rho+3/2)} \left\{ \psi(\rho+3/2) - \psi(\rho+1) \right. \\ \quad \left. + \sum_{k=1}^n \left( \frac{1}{\rho+1/2+k} - \frac{1}{\rho+3/2-k} \right) \right\} \\ \quad \cdot \prod_{k=1}^n \frac{\rho+3/2-k}{\rho+1/2+k}, & \text{otherwise.} \end{cases}$$

**3. Two Laguerre Polynomial Integrals.** In this section we consider the problem of evaluating the modified moments of the weight functions

$$(3.1) \quad w_p(x) = e^{-x} x^\rho (\ln x)^p, \quad \rho > -1, p = 1, 2, \dots,$$

on  $[0, \infty)$ , with respect to the generalized Laguerre polynomials  $L_n^{(\alpha)}(x)$ .

We first examine the case  $p = 1$ , which is relative to a weight function of mixed sign. By introducing, for notational convenience, a new parameter  $\mu$  such that  $\rho = \alpha + \mu$ , we refer to the integrals

$$(3.2) \quad N_{1,n}^{(\alpha)}(\mu) = \int_0^\infty e^{-x} x^{\alpha+\mu} \ln x L_n^{(\alpha)}(x) dx, \quad \alpha, \alpha + \mu > -1, n = 0, 1, 2, \dots,$$

for which, if we set

$$(3.3) \quad I_n^{(\alpha)}(\mu) = \int_0^\infty e^{-x} x^{\alpha+\mu} L_n^{(\alpha)}(x) dx,$$

we have

$$(3.4) \quad N_{1,n}^{(\alpha)}(\mu) = \frac{d}{d\mu} I_n^{(\alpha)}(\mu).$$

The evaluation of  $I_n^{(\alpha)}(\mu)$ , hence of  $N_{1,n}^{(\alpha)}(\mu)$ , can be carried out in a way much similar to that concerning the Jacobi polynomials described in Section 1. To this end, we need to recall the Rodrigues formula for Laguerre polynomials

$$e^{-x} x^\alpha L_n^{(\alpha)}(x) = \frac{1}{n!} \frac{d^n}{dx^n} (e^{-x} x^{n+\alpha}).$$

Indeed, inserting this last formula in (3.3) and integrating by parts  $n$  times, we obtain

$$I_n^{(\alpha)}(\mu) = \frac{(-1)^n}{n!} \frac{\Gamma(\mu+1)\Gamma(\mu+\alpha+1)}{\Gamma(\mu-n+1)}.$$

At this point it is not difficult to derive from (3.4) the following two expressions

$$(3.5) \quad N_{1,n}^{(\alpha)}(\mu) = \frac{(-1)^n}{n!} \frac{\Gamma(\mu+1)\Gamma(\mu+\alpha+1)}{\Gamma(\mu-n+1)} \cdot \{ \psi(\mu+1) + \psi(\mu+\alpha+1) - \psi(\mu-n+1) \},$$

with  $\alpha, \alpha + \mu > -1, n = 0, 1, 2, \dots$  and  $\mu \neq 0, 1, \dots, n-1$  if  $n \geq 1$ , and

$$(3.6) \quad N_{1,n}^{(\alpha)}(m) = (-1)^{m-1} \frac{m! (n-m-1)!}{n!} \Gamma(m+\alpha+1),$$

$$m = 0, 1, \dots, n-1, n \geq 1;$$

the second being obtained from the first by taking its limit as  $\mu \rightarrow m$ .

Finally, we remark that (3.5) may be put in the following form

$$(3.7) \quad N_{1,n}^{(\alpha)}(\mu) = \Gamma(\mu + \alpha + 1) \left\{ \psi(\mu + \alpha + 1) - \sum_{k=1}^n \frac{1}{k - \mu - 1} \right\} \prod_{k=1}^n \frac{k - \mu - 1}{k}.$$

The evaluation of the modified moments

$$N_{p,n}^{(\alpha)}(\mu) = \int_0^\infty e^{-x} x^{\alpha+\mu} (\ln x)^p L_n^{(\alpha)}(x) dx, \quad p \geq 2,$$

associated to the weight functions (3.1), can be obtained by repeatedly differentiating (3.5) with respect to  $\mu$ . We shall only examine, with some details, the case  $p = 2$ .

Differentiating (3.5) once, with respect to  $\mu$ , gives

$$(3.8) \quad N_{2,n}^{(\alpha)}(\mu) = \frac{(-1)^n}{n!} \frac{\Gamma(\mu + 1)\Gamma(\mu + \alpha + 1)}{\Gamma(\mu - n + 1)} \cdot \{(\psi(\mu + 1) + \psi(\mu + \alpha + 1) - \psi(\mu - n + 1))^2 + \psi'(\mu + 1) + \psi'(\mu + \alpha + 1) - \psi'(\mu - n + 1)\},$$

which holds for all  $n \geq 0$  with  $\mu \neq 0, 1, 2, \dots, n-1$ , when  $n \geq 1$ .

A more convenient form of (3.8), obtained by using the previously recalled properties of the functions  $\Gamma(x)$  and  $\psi(x)$ , together with the recurrence relation  $\psi'(x+1) = \psi'(x) - 1/x^2$ , is

$$(3.9) \quad N_{2,n}^{(\alpha)}(\mu) = \Gamma(\mu + \alpha + 1) \left\{ \left( \psi(\mu + \alpha + 1) - \sum_{k=1}^n \frac{1}{k - \mu - 1} \right)^2 + \psi'(\mu + \alpha + 1) - \sum_{k=1}^n \frac{1}{(k - \mu - 1)^2} \right\} \prod_{k=1}^n \frac{k - \mu - 1}{k}.$$

If  $\mu = m = 0, 1, 2, \dots, n-1, n \geq 1$ , from (3.8) we have

$$(3.10) \quad N_{2,n}^{(\alpha)}(m) = \lim_{\mu \rightarrow m} N_{2,n}^{(\alpha)}(\mu) = \frac{(-1)^n}{n!} \Gamma(m + 1)\Gamma(m + \alpha + 1) \{A_n(m) - 2B_n(m)\},$$

where

$$A_n(m) = \lim_{\epsilon \rightarrow 0} \frac{\psi^2(m + \epsilon - n + 1) - \psi'(m + \epsilon - n + 1)}{\Gamma(m + \epsilon - n + 1)},$$

$$B_n(m) = \lim_{\epsilon \rightarrow 0} \{ \psi(m + \epsilon + 1) + \psi(m + \epsilon + \alpha + 1) \} \frac{\psi(m + \epsilon - n + 1)}{\Gamma(m + \epsilon - n + 1)}.$$

By means of the two series expansions

$$\Gamma(x) = \frac{(-1)^r}{r!} \frac{1}{x + r} + \sum_{k=0}^{\infty} a_k (x + r)^k, \quad r = 0, 1, 2, \dots,$$

$$\psi(x) = \frac{-1}{x + r} + \psi(1 + r) + \sum_{k=0}^{\infty} b_k (x + r)^k,$$

which are valid for  $|x + r| < 1$ , it is easily seen that

$$\begin{aligned} A_n(m) &= (-1)^{n-m} 2(n-m-1)! \psi(n-m), \\ B_n(m) &= (-1)^{n-m} (n-m-1)! \{ \psi(m+1) + \psi(m+\alpha+1) \}. \end{aligned}$$

Hence, substituting these last two expressions into (3.10), we obtain the final result

$$\begin{aligned} (3.11) \quad N_{2,n}^{(\alpha)}(m) &= (-1)^m 2 \frac{m! (n-m-1)!}{n!} \Gamma(m+\alpha+1) \\ &\quad \cdot \{ \psi(n-m) - \psi(m+1) - \psi(m+\alpha+1) \}, \\ &\quad m = 0, 1, \dots, n-1, n \geq 1. \end{aligned}$$

**4. Some Particular Cases.** The results derived in Section 3 assume a very simple form when  $\mu = 0$ , that is in the cases where the weight functions

$$e^{-x} x^\alpha \ln x \quad \text{and} \quad e^{-x} x^\alpha (\ln x)^2, \quad \alpha > -1,$$

and the polynomials  $L_n^{(\alpha)}(x)$  have the same parameter  $\alpha$ .

For the first weight function, applying (3.7) and (3.6), we find

$$(4.1) \quad \int_0^\infty e^{-x} x^\alpha \ln x L_n^{(\alpha)}(x) dx = \begin{cases} \Gamma(\alpha+1) \psi(\alpha+1), & n = 0, \\ -\Gamma(\alpha+1)/n, & n \geq 1, \end{cases}$$

which may be regarded as a generalization of the well-known integral representation

$$\gamma = -\int_0^\infty e^{-x} \ln x dx,$$

of the Euler-Mascheroni constant  $\gamma = -\psi(1) = .57721\ 56649 \dots$

For the second weight function, by using (3.9) and (3.11), we obtain

$$(4.2) \quad \int_0^\infty e^{-x} x^\alpha (\ln x)^2 L_n^{(\alpha)}(x) dx = \begin{cases} \Gamma(\alpha+1) \{ \psi^2(\alpha+1) + \psi'(\alpha+1) \}, & n = 0, \\ \frac{2}{n} \Gamma(\alpha+1) \left\{ \sum_{k=1}^{n-1} \frac{1}{k} - \psi(\alpha+1) \right\}, & n \geq 1. \end{cases}$$

Two other cases of interest, which may be used, for instance, in constructing the modified moments of the weight functions

$$\exp(-x^2) \ln |x| \quad \text{and} \quad \exp(-x^2) (\ln |x|)^2,$$

on  $(-\infty, \infty)$ , with respect to the Hermite polynomials  $H_n(x)$ , are obtained by means of Szegő's relationships [4, p. 106] between Hermite and Laguerre polynomials,

$$(4.3) \quad H_{2n}(x) = (-1)^n 2^{2n} n! L_n^{(-1/2)}(x^2), \quad H_{2n+1}(x) = (-1)^n 2^{2n+1} n! x L_n^{(1/2)}(x^2).$$

Indeed, setting  $x = t^2$  in the integral (3.2), with  $\alpha = -1/2$  and  $\mu = 0$ , from (3.7) and (3.6) we obtain

$$4 \int_0^\infty \exp(-t^2) \ln t L_n^{(-1/2)}(t^2) dt = \begin{cases} \sqrt{\pi} \psi(1/2) = -\sqrt{\pi}(\gamma + 2 \ln 2), & n = 0, \\ -\sqrt{\pi}/n, & n \geq 1, \end{cases}$$

hence, by applying the first relation in (4.3),

$$(4.4) \quad \int_0^\infty \exp(-x^2) \ln x H_{2n}(x) dx = \begin{cases} \frac{-\sqrt{\pi}}{4} (\gamma + 2 \ln 2), & n = 0, \\ (-1)^{n-1} (n-1)! 2^{2(n-1)} \sqrt{\pi}, & n \geq 1. \end{cases}$$

Similarly, if we assume  $\alpha = \frac{1}{2}$  and  $\mu = -\frac{1}{2}$  in (3.2), then (3.7), together with the second relation in (4.3), gives

$$(4.5) \quad \begin{aligned} & \int_0^\infty \exp(-x^2) \ln x H_{2n+1}(x) dx \\ &= (-1)^{n-1} 2^{n-1} \left( \gamma + 2 \sum_{k=1}^n \frac{1}{2k-1} \right) \prod_{k=1}^n (2k-1), \end{aligned} \quad n = 0, 1, 2, \dots$$

The integrals involving the weight function  $\exp(-x^2)(\ln x)^2$  can be dealt with in the same way. Recalling that  $\psi'(\frac{1}{2}) = \pi^2/2$  and  $\psi'(1) = \pi^2/6$ , by use of (3.9) and (3.11), this leads to the following results,

$$(4.6) \quad \begin{aligned} & \int_0^\infty \exp(-x^2) (\ln x)^2 H_{2n}(x) dx \\ &= \begin{cases} \frac{\sqrt{\pi}}{8} \{ \psi^2(\frac{1}{2}) + \psi'(\frac{1}{2}) \} = 1.94752 \, 21803 \dots, & n = 0, \\ (-1)^n 2^{2(n-1)} (n-1)! \sqrt{\pi} \left( \gamma + 2 \ln 2 + \sum_{k=1}^{n-1} \frac{1}{k} \right), & n \geq 1, \end{cases} \end{aligned}$$

and

$$(4.7) \quad \begin{aligned} & \int_0^\infty \exp(-x^2) (\ln x)^2 H_{2n+1}(x) dx \\ &= (-1)^n 2^{n-2} \left\{ \left( \gamma + \sum_{k=1}^n \frac{2}{2k-1} \right)^2 + \frac{\pi^2}{6} - \sum_{k=1}^n \frac{4}{(2k-1)^2} \right\} \prod_{k=1}^n (2k-1), \end{aligned} \quad n = 0, 1, 2, \dots$$

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