CANONICAL COORDINATES AT A POINT FOR TWO SKEW-SYMMETRIC TENSORS

G. H. F. GARDNER

1. Introduction. The problem to be considered concerns an algebraic property of two skew-symmetric tensors at a point. Let P_{rs} and Q_{rs} be two skew-symmetric tensors. Under a transformation $C_{s'}^r = \partial x^r/\partial x^{s'}$ they become, in the new coordinate system, $P_{r's'}$ and $Q_{r's'}$, where

$$(1.1) P_{r's'} = C_{r'}^{m} C_{s'}^{n} P_{mn}, Q_{r's'} = C_{r'}^{m} C_{s'}^{n} P_{mn}.$$

We wish to find conditions on P_{rs} and Q_{rs} under which there will exist a nonsingular transformation $C_{s'}^r$ which does not reorient the axes, such that

$$(1.2) \quad \begin{array}{ll} P_{1'2'} = Q_{3'4'}, & Q_{1'2'} = -P_{3'4'}, & P_{2'3'} = Q_{1'4'}, \\ Q_{2'3'} = -P_{1'4'}, & P_{3'1'} = Q_{2'4'}, & Q_{3'1'} = -P_{2'4'} \end{array}$$

at a point A. In view of the restrictions on the transformation it is necessary that

$$\det C_{s'}^{r} > 0.$$

2. Let η^{rs} denote the diagonal array (1, 1, 1, -1) and ϵ^{rstu} the permutation symbol on 1234. We may write equations (1.2) as

$$\eta^{rm}\eta^{sn}P_{r's'}=\epsilon^{mntu}Q_{t'u'}.$$

This is equivalent to the equation

$$\eta^{rm}\eta^{sn}Q_{r's'} = -\epsilon^{mntu}P_{t'u'}$$

using latin indices ranging from 1 to 4 with the summation convention throughout.

A coordinate system in which (2.1) and (2.2) hold, we shall call a P-canonical coordinate system. If we interchange $P_{r's'}$ and $Q_{r's'}$ in (2.1) or (2.2) we get a Q-canonical coordinate system. We shall prove later that a P-system can be obtained from a Q-system by a transformation restricted as in (1.3) and vice-versa. For convenience we may define

$$(\epsilon PQ) = 1/2\epsilon^{mntu}P_{mn}Q_{tu}.$$

In this notation it is useful to note that

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$$(2.3) det P_{rs} = (\epsilon PP)^2$$

and also

(2.4)
$$\epsilon^{mntu} P_{mn} O_{us} = 1/2 \delta_s^t (\epsilon PQ).$$

Here we may dispose of some trivial cases. If at the point we are considering one of the tensors should happen to be identically zero, then the other must also be identically zero if canonical coordinates are to exist. This is clear from equation (2.1). If the two tensors should happen to be linearly dependent at the point, then both tensors must be identically zero, if canonical coordinates are to exist. For in such a system we would have

$$(2.5) aP_{r's'} + bQ_{r's'} = 0;$$

multiply this by $b\eta^{rm}\eta^{sn}$,

$$b^{2}\eta^{r}{}^{m}\eta^{s}{}^{n}Q_{r's'} = -ab\eta^{r}{}^{m}\eta^{s}{}^{n}P_{r's'}$$

$$= -ab\epsilon^{mn}{}^{tu}Q_{t'u'} \qquad \text{(by (2.1))}$$

$$= +a^{2}\epsilon^{mn}{}^{tu}P_{t'u'} \qquad \text{(by (2.5))}$$

$$= -a^{2}\eta^{r}{}^{m}\eta^{s}{}^{n}Q_{r's'} \qquad \text{(by (2.2))}.$$

Therefore we would have $Q_{r's'} \equiv 0$ and similarly $P_{r's'} \equiv 0$.

3. The following three theorems suffice to solve the problem.

THEOREM I. If p_{rs} and q_{rs} are two linearly independent, skew-symmetric tensors with matrices of rank 2 at a point A, then there exists a transformation of positive Jacobian to coordinates such that at the point A all six components of each tensor vanish except:

- (i) one in p_{rs} and one in q_{rs} , if $(\epsilon pq) \neq 0$,
- (ii) two in p_{rs} and two in q_{rs} , if $(\epsilon pq) = 0$.

THEOREM II. If p_{rs} and q_{rs} are two linearly independent, skew-symmetric tensors, with matrices of rank 2 at a point A, then both p-canonical and q-canonical coordinates can always be found for them at the point A by means of a transformation with a positive determinant.

THEOREM III. If P_{rs} and Q_{rs} are two linearly independent, skew-symmetric tensors, then a necessary and sufficient condition that both P-canonical and Q-canonical coordinates may be found at a point A by means of a transformation with a positive determinant is that, at A,

$$(3.1) \qquad (\epsilon PP) + (\epsilon QQ) = 0.$$

PROOF OF THEOREM I. The assumption that the two tensors are of

rank 2 is equivalent to assuming that the determinant of each vanishes and that neither tensor is identically zero, is since the tensors are of rank two and equations

$$(3.2) p_{rs}\lambda_a^s = 0 (a = 1, 2), q_{rs}\lambda_b^s = 0 (b = 3, 4)$$

can be satisfied by real values of $\lambda_{(a)}^s$ and $\lambda_{(b)}^s$. The vectors $\lambda_{(1)}^s$ and $\lambda_{(2)}^s$ are linearly independent and arbitrary within a linear homogeneous combination. For example $\lambda_{(1)}^r = B_{1m} \epsilon^{mrtu} p_{tu}$ and $\lambda_{(2)}^r = B_{2m} \epsilon^{mrtu} p_{tu}$ are solutions where B_{1m} and B_{2m} are any linearly independent sets of constants. This may be seen by using (2.4). $\lambda_{(a)}^r$ thus defines a 2-flat at the point A which we shall call the p-plane. Similarly $\lambda_{(b)}^r$ defines a 2-flat at the point A which we shall call the q-plane. These two planes will intersect in a common line if and only if (3.2) can be satisfied by a common vector λ^r . For this it is necessary and sufficient that det $(p_{rs} - kq_{rs}) = (\epsilon pp - 2k\epsilon pq + k^2\epsilon qq)^2 = 0$ for all k, that is,

$$(3.3) (\epsilon pq) = 0,$$

since $(\epsilon pp) = (\epsilon qq) = 0$, and k is arbitrary. Since p_{rs} and q_{rs} are linearly independent, the planes never coincide.

The argument is divided into two cases, (i) $(\epsilon pq) \neq 0$, (ii) $(\epsilon pq) = 0$.

(i) $(\epsilon pq) \neq 0$. Suppose we want, say, $p_{3'4'}$ and $q_{1'2'}$ to be the only surviving elements in the tensors. Then consider the transformation

$$(3.4) C_{s'}^r = \lambda_{(s)}^r$$

where $\lambda'_{(a)'}$ and $\lambda'_{(b)'}$ satisfy (3.2). The four vectors $\lambda'_{(s)'}$ are linearly independent since $(\epsilon pq) \neq 0$, by hypothesis. Hence det $C'_{s'} \neq 0$. Since each vector is arbitrary to within a constant factor, we can easily arrange to have det $C'_{s'} > 0$.

Substitution of (3.4) in (1.1) and use of (3.2) shows us immediately that the only nonzero elements in the two tensors are $p_{3'4'}$ and $q_{1'2'}$. For example,

$$p_{r'a'} = \lambda_{(r')}^m \lambda_{(a')}^n p_{mn} = 0$$
, and $q_{r'b'} = 0$.

This completes case (i). We may interpret transformation (3.4) in this manner: we choose the $x^{1'}$, $x^{2'}$ axes to lie in the p-plane and the $x^{3'}$, $x^{4'}$ axes to lie in the q-plane.

(ii) $(\epsilon pq) = 0$. The above transformation becomes singular since the planes now intersect in a line. Instead let us choose our x^r axes such that the x^1 axis lies in the p-plane but does not coincide with the line of intersection, the x^3 axis lies in the q-plane but does not coincide

¹ C. C. MacDuffee, Theory of matrices, Chelsea, p. 12.

with the line of intersection, and the x^2 , x^4 coordinate plane contains the line of intersection, but this line does not coincide with either the x^2 axis or the x^4 axis. We may write this type of transformation from x^r coordinates to $x^{r'}$ coordinates in the following form,

$$(3.5) C_{s'}^r = \mu_{(s')}^r,$$

where the vectors $\mu_{(s')}^r$ satisfy the equations

$$p_{rs}\mu_{1}^{r} = 0, q_{rs}\mu_{3}^{r} = 0,$$

$$p_{rs}(\mu_{2}^{r} + \mu_{4}^{r}) = 0, q_{rs}(\mu_{2}^{r} + \mu_{4}^{r}) = 0,$$

$$p_{rs}\mu_{a}^{r} \neq 0 (a = 2, 3, 4), q_{rs}\mu_{b}^{r} \neq 0 (b = 1, 2, 4),$$

It is clear that the vectors $\mu'_{(s')}$, restricted only by equations (3.6), can be chosen so that they are linearly independent. Then, by multiplying say $\mu'_{(1)}$ by a factor with a suitable sign, we shall have det $C'_{s'} > 0$.

Substitution of (3.5) in (1.1) and use of (3.6) shows immediately that the only surviving elements in the two tensors are $p_{3'4'}$, $p_{2'3'}$, $q_{1'2'}$, $q_{1'4'}$. For example,

$$p_{r'1'} = \mu_{(r')}^{m} \mu_{(1')}^{n} p_{mn} = 0,$$

$$p_{2'4'} = \mu_{(2')}^{m} \mu_{(4')}^{n} p_{mn} = -\mu_{(2')}^{m} \mu_{(2')}^{n} p_{mn} = 0,$$

$$p_{3'4'} = \mu_{(3')}^{m} \mu_{(4')}^{n} p_{mn} = -\mu_{(3')}^{m} \mu_{(2')}^{n} p_{mn} = p_{2'3'}.$$

In particular we notice that $p_{3'4'} = p_{2'3'}$ and $q_{1'2'} = q_{4'1'}$.

PROOF OF THEOREM II. There are two cases. Assume firstly $(\epsilon pq) \neq 0$. Suppose, as above, that the only nonzero elements in the two tensors are $p_{3'4'}$ and $q_{1'2'}$. Consider the diagonal transformation

$$(3.7) C_{s''}^{r'} = k_{(r)} \delta_s^{r} (r \text{ not summed}).$$

From this we obtain $p_{3''4''} = k_3k_4p_{3'4'}$ and $q_{1''2''} = k_1k_2q_{1'2'}$. We are demanding that det $C_{\delta''}^{r'} \equiv k_1k_2k_3k_4 > 0$. However $(\epsilon pq)'$ is a tensor density and therefore transforms thus:

$$(\epsilon pq)^{\prime\prime} = (\epsilon pq)^{\prime} \det C_{s^{\prime\prime}}^{r^{\prime}}.$$

That is, by (1.3), the sign of (ϵpq) is invariant under transformations with positive determinant. Now $(\epsilon pq)' = q_{1'2'}p_{3'4'}$, and so if $(\epsilon pq) > 0$ we can choose k_1 , k_2 , k_3 , k_4 such that $p_{3''4''} = q_{1''2''}$, and this is a q-canonical coordinate system. If $(\epsilon pq) < 0$ then k_1 , k_2 , k_3 , k_4 can be chosen so that $p_{3''4''} = -q_{1''2''}$ which is a p-canonical system. How-

ever these two systems are equivalent to one another since we can get from a p-system to a q-system by a nonsingular transformation which does not reorient the axes. To see this consider the transformation

$$C_{\bullet^{\prime\prime\prime}}^{\prime\prime\prime} = \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix}$$

applied to a q-canonical coordinate system in which $p_{3''4''} = q_{1''2''}$ are the only nonzero elements. We obtain in the new system $p_{1'''2'''} = q_{3'''4'''}$ and all other components are zero; but this is a p-canonical coordinate system, and det $C_{s'''}^{r'}=1$. Similarly transformation (3.8) turns a p-system in which p_{12} and q_{34} are the only nonzero elements into a q-system in which p_{34} and q_{12} are the only nonzero elements. Thus when $(\epsilon pq) \neq 0$ we can find a transformation $C_{s'}^{r}$, det $C_{s'}^{r} > 0$, which will give either p-canonical or q-canonical coordinates.

Now assume $(\epsilon pq) = 0$. Let us take a coordinate system in which each tensor has only two nonvanishing elements, say $p_{3'4'} = p_{2'3'}$ and $q_{1'2'} = q_{4'1'}$, and apply to it the diagonal transformation

$$(3.9) C_{s''}^{r'} = l_{(r)} \delta_{s''}^{r'} (r \text{ not summed})$$

where $l_{(r)} = (p_{3'4'}, l_2q_{1'2'}, l_4)$. In the $x^{r''}$ system we then have

$$p_{3''4''} = l_4 q_{1'2'} p_{3'4'}, \qquad q_{1''2''} = l_2 p_{3'4'} q_{1'2'},$$

$$p_{2''3''} = l_2 q_{1'2'} p_{3'4'}, \qquad q_{1''4''} = -l_4 p_{3'4'} q_{1'2'}.$$

If $p_{3'4'}q_{1'2'}>0$ we can take $l_2=l_4$, thus making det $C_{s''}^{'}>0$, and have a q-canonical system; while if $p_{3'4'}q_{1'2'}<0$ we can take $l_2=-l_4$ and have a p-canonical system. As above we can transform from a p-system to a q-system very easily. In this case apply the transformation

(3.10)
$$C_{a'''}^{r''} = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix}.$$

This completes the proof of Theorem II. It is now evident that for skew-symmetric tensors with vanishing determinants we can pass from any p-canonical coordinate system to a q-canonical coordinate system and vice-versa by means of a transformation of positive determinant.

PROOF OF THEOREM III. That equation (3.1) is a necessary condition may be seen by multiplying (2.1) by $Q_{m'n'}$ and (2.2) by $-P_{m'n'}$ and adding. Now if (3.1) holds and if the determinant of one tensor vanishes so must the determinant of the other. The case in which both determinants are zero has already been dealt with by Theorem II. We can therefore now assume that neither (ϵPP) nor (ϵQQ) vanishes. We shall show the sufficiency of (3.1) by reducing the problem to the case already considered in Theorem II.

Let us assume (3.1) is satisfied at the point A. Consider the equation

$$\det (P_{rs} - \theta Q_{rs}) \equiv (\epsilon PP - 2\theta \epsilon PO + \theta^2 \epsilon QO)^2 = 0.$$

Since $(\epsilon QQ) \neq 0$, $(\epsilon PP) \neq 0$ we may write this, using (3.1), as

$$(3.11) \theta^2 - 2\theta \frac{\epsilon PQ}{\epsilon QQ} - 1 = 0.$$

This quadratic equation always has two real distinct roots. Let them be θ_1 and $-1/\theta_1$. These roots are invariant under all transformations. If we now define

(3.12)
$$p_{rs} = P_{rs} - \theta_1 Q_{rs},$$

$$q_{rs} = \theta_1 P_{rs} + Q_{rs},$$

we have det $p_{rs} = \det q_{rs} = 0$ as in Theorem II. Equations (3.12) have a tensor character and therefore remain the same in all coordinate systems. Furthermore,

$$(\epsilon pq) = (\epsilon (P - \theta_1 Q)(\theta_1 P + Q))$$

$$\cdot = (1 - \theta_1^2)(\epsilon PQ) + \theta_1((\epsilon PP) - (\epsilon QQ));$$

using (3.1) and (3.11) we obtain

(3.13)
$$(\epsilon pq) = \frac{2\theta_1}{(\epsilon PP)} \left\{ (\epsilon PP)^2 + (\epsilon PQ)^2 \right\}.$$

Since we are assuming that $(\epsilon PP) \neq 0$, it follows that $(\epsilon pq) \neq 0$; moreover since we can take θ_1 as either the positive or negative root of (3.11), we may have either $(\epsilon pq) > 0$ or $(\epsilon pq) < 0$ as we please. By Theorem II we can find both p-canonical and q-canonical coordinates for p_{rs} and q_{rs} . We shall now establish the penultimate result that:

If the coordinates are p-canonical or q-canonical for p_{rs} and q_{rs} they are also respectively P-canonical or Q-canonical for P_{rs} and Q_{rs} , as related by (3.12).

Assuming the coordinates to be p-canonical for p_{rs} and q_{rs} and using (3.12) it is clear that

$$(3.14) \eta^{rm}\eta^{sn}(P_{rs}-\theta_1Q_{rs})=\epsilon^{mntu}(\theta_1P_{tu}+Q_{tu}),$$

$$(3.15) \eta^{rm}\eta^{sn}(\theta_1P_{rs}+Q_{rs})=-\epsilon^{mntu}(P_{tu}-\theta_1Q_{tu}).$$

Multiplying (3.15) by θ_1 and adding that to (3.14) we see that

$$(1 + \theta_1^2) \eta^{rm} \eta^{sn} P_{rs} = (1 + \theta_1^2) \epsilon^{mntu} Q_{tu}.$$

However $1+\theta_1^2\neq 0$, θ_1 being real, and therefore the coordinates are P-canonical for P_{rs} and Q_{rs} . Similarly if the coordinates are q-canonical for p_{rs} and q_{rs} , they are Q-canonical for P_{rs} and Q_{rs} . We have already seen that we may pass from a p-canonical coordinate system to a q-canonical coordinate system by a transformation of positive Jacobian. Hence we deduce at once that Theorem III is true.

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CARNEGIE INSTITUTE OF TECHNOLOGY