BEHAVIOR OF SOLUTIONS OF SECOND ORDER SELF-ADJOINT DIFFERENTIAL EQUATIONS

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Introduction. In this paper the modified polar coordinate transformation,

(I) $$y(x) = \rho(x) \sin \theta(x), \qquad y'(x) = \frac{w(x)}{r(x)} \rho(x) \cos \theta(x),$$

will be applied to the self-adjoint equation,

(1) $$(ry')' + qy = 0.$$

For $x \geq a$, let $q(x)$ be a function of class C, $r(x)$ be a positive function of class C, $w(x)$ be a positive function of class C', and $y(x)$ be a nontrivial solution of equation (1).[1] The reader can show that there exist functions $\rho(x)$ and $\theta(x)$ of class C' which satisfy (I) and $\rho(x) > 0$. Furthermore,

(II$_1$) $$\rho' = \rho \left[\left(\frac{w}{r} - \frac{q}{w} \right) \frac{\sin 2\theta}{2} - \frac{w'}{w} \cos^2 \theta \right]$$

and

(II$_2$) $$\theta' = \frac{1}{2} \left(\frac{w}{r} + \frac{q}{w} \right) + \frac{1}{2} \left(\frac{w}{r} - \frac{q}{w} \right) \cos 2\theta + \frac{w'}{2w} \sin 2\theta.$$

The transformation (I) is an extension of the polar transformation

(I') $$y(x) = \rho(x) \sin \theta(x), \qquad y'(x) = \rho(x) \cos \theta(x)$$

of the normal form of the ordinary wave differential equation

(1') $$y'' + q(x)y = 0$$

which was introduced by Prüfer [6]. For well-known applications of (I') to the self-adjoint equation (1) the reader is referred to [2, pp. 161–167] and [3, pp. 274–281]. W. M. Whyburn [8] has used this transformation in studying solutions of a system of two first order nonlinear equations. More recently F. V. Atkinson [1] has employed

[1] That is, on $x \geq a$, $y(x)$ satisfies (1), y and ry' are of class C', and y is not identically zero.

a special case of (I) in transforming (1'), namely,

$$(\mathrm{I''}) \qquad y(x) = \rho(x)\cos\theta(x), \qquad \frac{y'(x)}{(q(x))^{1/2}} = \rho(x)\sin\theta(x).$$

The first section consists of consequences of the first equation (II_1), the ρ-equation. Certain theorems on functional bounds for solutions of equations (1) are given which are extensions of results of Levinson [5] and Leighton [4]. In the second section the equation (II_2), the θ-equation, will be used to establish sufficient conditions for oscillation of solutions of (1) and these are compared with necessary conditions established by Leighton [4].

Finally, an asymptotic form of solutions of (1) is obtained.

1. Boundedness. The first theorem is an extension of a theorem of Levinson.

THEOREM 1. *If $w'(x) \leqq 0$, $Q(x) = \int_a^x \left| w/r - q/w \right| dt$, $x \geqq a$, then*

$$(\mathrm{III}) \qquad \rho(a)\exp\left[-Q(x)/2\right] \leqq \rho(x) \leqq \frac{\rho(a)\,w(a)}{w(x)}\exp\left[Q(x)/2\right],$$

and for each solution $y(x)$ of (1):

$$(2) \qquad\qquad |y(x)| \leqq \frac{\rho(a)\,w(a)}{w(x)}\exp\left[Q(x)/2\right],$$

$$(3) \qquad\qquad |y'(x)| \leqq \frac{\rho(a)\,w(a)}{r(x)}\exp\left[Q(x)/2\right].$$

PROOF. Equation (II_1) yields the inequalities:

$$-\frac{1}{2}\left|\frac{w}{r} - \frac{q}{w}\right| \leqq \frac{\rho'}{\rho} \leqq \frac{1}{2}\left|\frac{w}{r} - \frac{q}{w}\right| - \frac{w'}{w}$$

from which (III) is obtained by integration. Inequalities (2) and (3) follow from (III).

COROLLARY 1.1. *If, in addition to the hypotheses of Theorem 1, $Q(x)$ is bounded for $x \geqq a$, then for each solution $y(x)$ of (1):*

$$(4) \qquad\qquad y(x) = O\!\left(\frac{1}{w(x)}\right) \quad and \quad y'(x) = O\!\left(\frac{1}{r(x)}\right), \qquad as\ x \to \infty.$$

Let $w = (qr)^{1/2}$, then $Q(x) = 0$ and a result of Leighton is established:

COROLLARY 1.2. *If $q(x) > 0$, $x \geqq a$, $q(x) \cdot r(x)$ is of class C' and $(qr)' \leqq 0$, then*

(5') $$y(x) = O\left(\frac{1}{(rq)^{1/2}}\right) \quad \text{and} \quad y'(x) = O\left(\frac{1}{r}\right) \qquad \text{as } x \to \infty.$$

COROLLARY 1.3. *If, in addition to the hypotheses of Corollary 1.2, $w(x)$ (or $r(x)$) is bounded away from zero, then $y(x)$ (or $y'(x)$) is bounded.*

2. **Oscillation.** Note that $y(x)$ has a zero only when $\theta(x)$ is a multiple of π. If $\theta(x)$ is equal to an integral multiple of π at $x = x_1$, then $\theta'(x_1) = w(x_1)/r(x_1) > 0$ and, therefore, if $\theta(x)$ takes on integral multiples of π for infinitely many values of x, then $\lim_{x \to \infty} \theta(x) = \infty$. Therefore, in order that $y(x)$ be oscillatory it is necessary and sufficient that

$$\lim_{x \to \infty} \theta(x) = \infty.$$

The θ-equation (II$_2$) yields the inequality:

(IV) $$\theta' \geqq \frac{1}{2}\left(\frac{w}{r} + \frac{q}{w}\right) - \frac{1}{2}\left|\frac{w}{r} - \frac{q}{w}\right| - \frac{1}{2}\frac{|w'|}{w}$$

from which the following theorem is readily obtained.

THEOREM 2. *If there exists a positive function $w(x)$ of class C' on $x \geqq a$ such that*

$$\lim_{x \to \infty} \int_a^x \left[\left(\frac{w}{r} + \frac{q}{w}\right) - \left|\frac{w}{r} - \frac{q}{w}\right| - \frac{|w'|}{w}\right] dt = \infty,$$

then every nontrivial solution of (1) is oscillatory.

As in Corollary 1.2, the special choice $w = (qr)^{1/2}$ gives a simplified form of (IV) from which the following *sufficient* condition for oscillation is derived:

COROLLARY 2.1. *If $q(x) > 0$, $(qr)' \leqq 0$ on $x \geqq a$, and*

$$\lim_{x \to \infty} \left[\int_a^x \left(\frac{q}{r}\right)^{1/2} dt + \frac{1}{4}\ln(q(x)r(x))\right] = \infty$$

then every nontrivial solution of (1) is oscillatory.

It is of interest to compare this sufficient condition with Leighton's corresponding necessary condition

$$\lim_{x \to \infty} \int_a^x \left(\frac{q}{r}\right)^{1/2} dt = \infty.$$

This corollary establishes that Leighton's condition is also sufficient for cases where the function qr is bounded away from zero. However, in many interesting cases,

$$\lim_{x\to\infty} r(x)q(x) = 0,$$

for example, the Euler equation where $q(x) = k/x^2$ and $r(x) = 1$. For this example Corollary 2.1 shows that oscillation occurs for $k > 1/4$.

3. Asymptotic behavior.

THEOREM 3. *If k is a positive number such that $\int_a^\infty |k/r - q/k| \, dt < \infty$, $\beta(x) = (1/2)\int_a^x (k/r + q/k)dt$, and $y(x)$ is any nontrivial solution of (1), then there exists a positive number A and a number α such that*

$$\lim_{x\to\infty} \left[y(x) - A \sin (\beta(x) + \alpha) \right] = 0.$$

PROOF. Let $w = k$ in equations (II). From a well-known theorem, it follows that $(1/2)\int_a^\infty (k/r - q/k) \sin 2\theta dt$ exists (i.e. is finite). Call this value A_1. Then since $\rho(x)$ satisfies (II$_1$):

$$\lim_{x\to\infty} \rho(x) = \rho(a) \cdot e^{A_1} = A > 0.$$

Furthermore, from (II$_2$):

$$\theta(x) = \theta(a) + \beta(x) + \frac{1}{2} \int_a^x (k/r - q/k) \cos 2\theta dt.$$

Hence, if $\alpha = \theta(a) + (1/2)\int_a^\infty (k/r - q/k) \cos 2\theta dt$, then

$$\lim_{x\to\infty} \left[\theta(x) - \beta(x) - \alpha \right] = 0.$$

Finally, it follows that

$$\lim_{x\to\infty} \left[y(x) - A \sin (\beta(x) + \alpha) \right] = 0.$$

Under the hypothesis of Theorem 3, the oscillation or nonoscillation of $y(x)$ depends on whether or not $\beta(x) \to \infty$ as $x \to \infty$. If

$$\lim_{x\to\infty} \int_a^x dt/r(t) = \infty,$$

then $\lim_{x\to\infty} \int_a^x q(t)dt = \infty$, since $\left| \int_a^x (k/r)dt - \int_a^x (q/k)dt \right| \leq \int_a^x |k/r - q/k| \, dt$. Hence $\beta(x) \to \infty$ as $x \to \infty$, and there exists a sequence $\{x_n\}$

such that $\beta(x_n)+\alpha = (4n-3)\pi/2$, $x_n \to \infty$, and $\lim_{n\to\infty} y(x_n) = A > 0$. This result is summarized in

COROLLARY 3.1. *If, in addition to the hypotheses of Theorem 3,*

$$\lim_{x\to\infty} \int_a^x \frac{dt}{r(t)} = \infty,$$

then $\lim \sup_{x\to\infty} y(x) > 0$, *and* $y(x)$ *oscillates as* $x \to \infty$.

On the other hand, if $\int_a^\infty dt/r(t) < \infty$ and since $|q/k| \leqq k/r + |q/k - k/r|$, then $\int_a^\infty |q(t)| dt < \infty$ and $\lim_{x\to\infty} \beta(x)$ exists. Therefore, the next result follows easily.

COROLLARY 3.2. *If, in addition to the hypotheses of Theorem 3,*

$$\int_a^\infty \frac{dt}{r(t)} < \infty,$$

then $\lim_{x\to\infty} y(x)$ *and* $\lim_{x\to\infty} r(x)y'(x)$ *exist.*

For a different proof of an equivalent result by Wintner see [7, p. 58].

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