LARGE-TIME BEHAVIOR OF SOLUTIONS TO CERTAIN QUASILINEAR PARABOLIC EQUATIONS IN SEVERAL SPACE DIMENSIONS

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ABSTRACT. We consider the Cauchy problem, $u_t + \operatorname{div} f(u) = \Delta u$ for $x \in \mathbb{R}^n$, t > 0 with $u(x,0) = u_0(x)$. For n = 1, suppose f'' > 0 and $\int |u_0 - \phi| dx < \infty$ where ϕ is piecewise constant and $\phi(x) \to u^+$ (u^-) as $x \to +\infty$ ($-\infty$). A result of Il'in and Oleinik states that if $\phi(x - kt)$ is an entropy solution of $u_t + \operatorname{div} f(u) = 0$, then u(x,t) approaches a traveling wave solution, $\tilde{u}(x - kt)$, as $t \to \infty$, with $\tilde{u}(x) \to u^+$ (u^-) as $x \to +\infty$ ($-\infty$). We give two examples which show that this result does not hold for $n \ge 2$.

This work concerns the asymptotic behavior of the solution to the Cauchy problem,

(1)
$$u_t + \operatorname{div} f(u) = \Delta u \quad \text{for } x \in \mathbb{R}^n, \ t > 0,$$

(2)
$$u(x,0) = u_0(x) \quad \text{for } x \in \mathbb{R}^n,$$

where $f \in C^2(\mathbf{R}; \mathbf{R}^n)$, and $n \ge 2$. We assume throughout the paper that u_0 is a bounded function which approaches a piecewise constant state as $|x| \to \infty$ in the sense that $\int_{\mathbf{R}^n} |u_0 - \phi| dx < \infty$, where

(3)
$$\phi(x) = a \cdot \chi_{\Omega}(x) + b \cdot \chi_{\mathbf{R}' \setminus \Omega}(x)$$

with $a \ge b$, and Ω is a connected open subset of \mathbb{R}^n (with piecewise smooth boundary if $n \ge 2$). We also assume without loss of generality that f(a) = f(b) = 0, since the transformation $\bar{x} = x - kt$, $\bar{t} = t$, where k = [f(a) - f(b)]/(a - b), yields an equation of the form (1) which satisfies this condition.

For n = 1, a result of Il'in and Oleinik [3] states that if f is strictly convex and

$$\phi(x) = \phi_c(x) \equiv \begin{cases} a, & x < c, \\ b, & x > c, \end{cases}$$

for some $c \in \mathbb{R}$, then $\lim_{t \to \infty} u(x, t) \equiv \tilde{u}(x)$ exists and satisfies

(4)
$$\operatorname{div} f(\tilde{u}) = \Delta \tilde{u} \quad \text{in } \mathbf{R}^n,$$

(5)
$$\tilde{u}(x) - \phi(x) \to 0 \text{ as } |x| \to \infty, \text{ and}$$

$$\int_{\mathbf{R}^n} (\tilde{u} - \phi) dx = \int_{\mathbf{R}^n} (u_0 - \phi) dx.$$

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Note that any function ϕ of the form (3) satisfies

$$(6) u_t + \operatorname{div} f(u) = 0.$$

Also since f is convex, the functions $\{\phi_c\}$ are precisely those of the form (3) which satisfy the entropy condition (see (E) below). Thus in one space dimension, $\lim_{t\to\infty} u(x,t)$ exists and satisfies (4) and (5) whenever ϕ satisfies (E).

It is easy to see that the converse is also true. That is, if $\lim_{t\to\infty} u(x,t)$ exists and satisfies (4) and (5), there is a function ψ of the form (3) which satisfies the entropy condition and

$$\int_{\mathbf{R}} |\phi - \psi| \, dx < \infty.$$

The entropy condition for solutions of (6) when $n \ge 1$ was formulated by Kruzkov and Vol'pert in [4 and 6]. When applied to ϕ as in (3), it can be stated as follows:

(E)
$$\langle f(c), n(x) \rangle \leq 0$$
 for all $c \in (b, a)$

 H^{n-1} -almost everywhere on $\partial\Omega$, where n(x) is the outward-pointing normal to Ω at x.

In this paper we give two examples which demonstrate that Il'in and Oleinik's result on the large-time behavior of solutions to (1) and (2) fails in dimension $n \ge 2$. In both examples, f satisfies the strong convexity condition formulated by Conway in [1]. (In fact, we take f = (0, ..., 0, F) with F strictly convex.) We take $u_0 = \phi$ where ϕ has the form (3) and ϕ satisfies (E).

In our first example, $\tilde{u}(x) \equiv \lim_{t \to \infty} u(x, t)$ exists and satisfies the elliptic equation (4), but does not inherit the asymptotic values of ϕ as $|x| \to \infty$, as in (5). In our second example, \tilde{u} fails to exist.

EXAMPLE 1. Let f(u) = (0, 0, ..., F(u)) where F(u) is a smooth strictly convex function with F(0) = F(1) = 0. Set $\phi(x) = \chi_{\Omega}(x)$ where

$$\Omega = \{ x = (x', x_n) : x_n < -|x'| \}.$$

One readily checks that ϕ satisfies condition (E).

We take u(x, t) to be the solution of (1) and (2) with $u_0(x) = \phi(x)$. Such a solution will exist and be unique in the class

$$C([0,\infty);L^1_{loc}(\mathbf{R}^n))\cap L^\infty(\mathbf{R}^{n+1})\cap C^2(\mathbf{R}^n\times(0,\infty))$$

and satisfies $0 \le u \le 1$. This can be seen by considering a sequence of smooth, bounded functions $u_0^N(x) \to \phi(x)$ in $L^1_{loc}(\mathbb{R}^n)$. Taking the corresponding solutions of the Cauchy problem [5, V, Theorem 8.1] and arguing as in [6, §17.2], one obtains both existence and uniqueness in the indicated class. We will show that $\lim_{t\to\infty} u(x,t) \equiv 0$.

Consider first for the sake of comparison the function $g_d(x_n, t)$ satisfying

(7)
$$D_{x_{n}x_{n}}g_{d} = D_{t}g_{d} + D_{x_{n}}(F(g_{d})) \quad \text{for } -\infty < x_{n} < \infty, \ t > 0,$$

$$g_{d}(x_{n}, 0) = \begin{cases} 1, & x_{n} < d, \\ 0, & x_{n} > d. \end{cases}$$

From [3] the functions

$$\int_{-\infty}^{-y} \left(1 - g_d(s, t)\right) ds, \qquad \int_{y}^{\infty} g_d(s, t) ds$$

exist, are continuous for $y \in \mathbf{R}$ and $t \ge 0$, and together with

$$(1-g_d(-y,t)), \quad g_d(y,t)$$

tend to zero as $y \to +\infty$ uniformly in t.

Set $V_d(x_n, t) = \int_{x_n}^{\infty} g_d(s, t) ds$; this is continuous for $t \ge 0$ and satisfies

(8)
$$D_{x_n x_n} V_d = -F(-D_{x_n} V_d) + D_t V_d, \qquad V_d(x_n, 0) = (d - x_n)^+.$$

Moreover

$$\tilde{g}_d(x_n) \equiv \lim_{t \to \infty} g_d(x_n, t)$$
 and $\tilde{V}_d(x_n) \equiv \lim_{t \to \infty} V_d(x_n, t)$

both exist and are solutions to (7) and (8); they are uniquely determined respectively by

$$\int_{d}^{\infty} \tilde{g}_{d}(s) ds + \int_{-\infty}^{d} (\tilde{g}_{d}(s) - 1) ds = 0, \text{ and } \tilde{V}_{d}(x_{n}) = \int_{x}^{\infty} \tilde{g}_{d}(s) ds.$$

Finally the convergence to \tilde{g}_d is uniform in x_n .

From the maximum principle

$$0 \le u(x,t) \le g_0(x_n,t)$$
 for $x \in \mathbb{R}^n$, $t \ge 0$.

Thus $U(x, t) = \int_{x_{-}}^{\infty} u(x', s, t) ds$ is well defined and satisfies

$$\Delta U = -F(-D_{x_{-}}U) + D_{t}U, \qquad U(x,0) = (-|x'| - x_{n})^{+}.$$

For any d < 0, $w_d \equiv U - V_d$ satisfies

$$\Delta w_d = a(x, t, d) D_{x_n} w_d + D_t w_d$$
 and $w_d(x, 0) \le -d \cdot \chi_{\{|x'| \le -d\}}(x)$.

Again using the maximum principle, $w_d \le h(x', t)$ for $t \ge 0$ where

$$\Delta_{x'}h = D_t h$$
 for $x' \in \mathbb{R}^{n-1}$, $t > 0$

and

$$h(x',0) = -d \cdot \chi_{\{|x'| \le -d\}}(x'), \qquad x' \in \mathbb{R}^{n-1}.$$

Hence

$$\overline{\lim}_{t\to\infty} U(x,t) \leqslant \tilde{V}_d(x_n) + \lim_{t\to\infty} h(x',t) = \tilde{V}_d(x_n).$$

Since $g_d(x_n, t) = g_0(x_n - d, t)$, $\tilde{V}_d(x_n) \to 0$ as $d \to -\infty$. From parabolic estimates it follows that

$$u(x,t) = -D_{x,}U(x,t) \to 0$$
 as $t \to \infty$.

Example 2. Consider f(u) as above and $u_0(x) = \phi(x) = \chi_{\Omega}(x)$ where

$$\Omega = \{ (x', x_n) : x_n < \psi(x') \} \quad \text{with } \psi \in C^1(\mathbf{R}^{n-1}), -1 \le \psi \le 0,$$

and

$$\overline{\lim}_{r\to\infty} r^{1-n} \int_{|x'|\leqslant r} \psi \, dx' \neq \underline{\lim}_{r\to\infty} \int_{|x'|\leqslant r} \psi \, dx'.$$

Again one can check that (E) is satisfied. We show that $\lim_{t\to\infty} u(x,t)$ cannot exist pointwise almost everywhere on \mathbb{R}^n .

Using the previous remarks on g_d we see that

$$g_{-1}(x_n,t) \leqslant u(x,t) \leqslant g_0(x_n,t)$$

and that the function

$$U(x',t) = \int_{-\infty}^{0} (1 - u(x',s,t)) ds - \int_{0}^{\infty} u(x',s,t) ds$$

is well defined and satisfies

$$\Delta_{x'}U = D_{t}U \quad \text{for } t > 0, \ x' \in \mathbb{R}^{n-1},$$

$$U(x', 0) = -\psi(x') \quad \text{for } x' \in \mathbb{R}^{n-1}.$$

If $\lim_{t\to\infty} u(x,t)$ is well defined, then $\lim_{t\to\infty} U(x',t) \equiv \tilde{U}(x')$ exists. Hence \tilde{U} is a bounded harmonic function and thus a constant. By a result of [2] this is true iff

$$\lim_{r \to \infty} r^{1-n} \int_{|x'| \le r} \psi \, dx' \quad \text{exists.}$$

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