# THE BEHAVIOR OF THE ANALYTICALLY CONTINUED RESOLVENT OPERATOR NEAR $\kappa=0$ AND AN APPLICATION TO ENERGY DECAY 

KAZUHIRO YAMAMOTO

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#### Abstract

We shall study the behavior of the analytically continued resolvent operator $R^{+}(\kappa)$ for perturbations of $-\Delta$ in a neighborhood of $\kappa=0$. As an application, making use of Vainberg's argument, we shall show the local energy decay of solutions to generalized wave equations whose stationary problems are not positive definite.


## 1. Introduction and results

In this paper we shall study the behavior of the analytically continued resolvent operator $R^{+}(\kappa)$ near $\kappa=0$ of the following problem:

$$
\left\{\begin{array}{l}
\left(L-\kappa^{2}\right) u=-\alpha(x)\left[\partial_{j} a_{j k}(x) \partial_{k}-q(x)\right] u-\kappa^{2} u=f \quad \text { in } \Omega,  \tag{1.1}\\
B u=0 \text { on } \partial \Omega
\end{array}\right.
$$

where $\Omega$ is $R^{n}$ or an exterior domain of $R^{n}(n \geq 3)$ with the $C^{2}$-class boundary $\partial \Omega$ and $B$ is either the Dirichlet boundary condition or the third boundary condition of the form $B u=\nu_{j}(x) a_{j k} \partial_{k} u+\sigma(x) u$ with the unit outer normal vector $\nu(x)=\left(\nu_{1}, \ldots, \nu_{n}\right)$ at $x \in \partial \Omega$. Assumptions on the coefficients of $L$ and $B$ are as follows:
(1.2) The function $\alpha(x)$ is bounded, measurable, and real valued, and is uniformly positive in $R^{n}$. The real symmetric matrix $\left(a_{j k}(x)\right)$ is uniformly positive in $R^{n}$ and its components are in $C^{1}\left(R^{n}\right)$. Moreover there exist positive constants $C$ and $b$ such that

$$
|\alpha(x)-1|+\left|a_{j k}(x)-\delta_{j k}\right|+\left|\nabla a_{j k}(x)\right| \leq C e^{-2 b|x|} \quad \text { in } R^{n} .
$$

(1.3) The function $q(x)$ is real and in $L_{\mathrm{loc}}^{p}\left(R^{n}\right)$, where $p=n / 2$ for $n \geq 5$, $p>2$ for $n=4$ and $p=2$ for $n=3$. The real valued function $\sigma(x)$ is in $C^{2}(\partial \Omega)$. Moreover there exist positive numbers $C$ and $b$ such that for sufficiently large $|x| \quad|q(x)| \leq C e^{-2 b|x|}$.

Let $H_{\rho_{1}}^{m}(\Omega)$ be a set $\left\{f \in H_{\mathrm{loc}}^{m}\{\bar{\Omega}) ; \rho_{1} \partial_{x}^{\alpha} f \in L^{2}(\Omega)\right.$ for $\left.|\alpha| \leq m\right\}$ and $L_{\rho_{1}}^{2}(\Omega)=H_{\rho_{1}}^{0}(\Omega)$. We denote by $\Gamma_{a}$ a set $\{\kappa \in C ; \operatorname{Im} \kappa>-a\}$ if $n$ is odd,

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and a set $\{\kappa \in C \backslash 0 ;|\operatorname{Im} a|<a,-\infty<\arg \kappa<\infty\} \cup\{\kappa \in C ; 0<\arg \kappa<\pi\}$ if $n$ is even. Then in Theorem 4.6 of our previous paper [9] under the above assumptions we show that for any positive $a<b$ there exists $R^{+}(\kappa)$ which is a finitely meromorphic function in $\Gamma_{a}$ with values in $\mathscr{B}\left(L_{\rho}^{2}(\Omega), H_{\rho^{-1}}^{2}(\Omega)\right)$ such that $u(x)=R^{+}(\kappa) f$ is the solution of (1.1) if $\kappa$ is an analytic point of $R^{+}(\kappa)$, where $\rho(x)=e^{-a|x|}$.

The purpose of this paper is to study the behavior of $R^{+}(\kappa)$ in a neighborhood of $\kappa=0$. It is well known that $L$ is a selfadjoint operator on $\mathscr{H}$ with the domain $D(L)=\left\{f \in H^{2}(\Omega) ; B u=0\right.$ on $\left.\partial \Omega\right\}$, where $\mathscr{H}=\left\{f(x) ;\|f\|_{\mathscr{K}}^{2}=\right.$ $\left.\left(f, f \alpha^{-1}\right)_{L^{2}(\Omega)}<\infty\right\}$ (see Lemma 2.1 of [8]). We denote by $H_{D}(\Omega)$ the completion of $C_{0}^{\infty}(\bar{\Omega})\left(C_{0}^{\infty}(\Omega)\right)$ by the Dirichlet norm $\|\nabla f\|_{L^{2}(\Omega)}$ if $B$ is the third (Dirichlet) boundary condition. Let $H$ be $H_{D}(\Omega) \times \mathscr{H}$ and $A=\left(\begin{array}{cc}0 & 1 \\ -L & 0\end{array}\right)$ be an operator with the domain $\left\{f={ }^{t}\left(f_{1}, f_{2}\right) \in H ; f_{2} \in L^{2}(\Omega) \cap H_{D}(\Omega)\right.$, $\partial_{x}^{2} f_{1} \in L^{2}(\Omega)$ for $|\alpha|=2, B f_{1}=0$ (on $\left.\partial \Omega\right\}$.) Put $E_{n}=\left\{\kappa ; \kappa^{2}\right.$ is a nonpositive eigenvalue of $L\}$ if $n \geq 5$ and put $E_{n}=\{\kappa ; \kappa=i \mu, \mu \geq 0, \pm \mu$ is an eigenvalue of $A\}$ if $n=3,4$. We note that $E_{n} \backslash\{0\}=\left\{\kappa ; \kappa^{2}\right.$ is a negative eigenvalue of $L\} \quad(n=3,4)$ and that from Lemma 2.1 of [8] the number of elements of $E_{n}$ is finite and the dimensions of each corresponding eigenspace are finite. We shall prove the following:
Theorem 1.1. We have the following two statements:
(i) If $n$ is odd and $0 \notin E_{n}$, then $R^{+}(\kappa)$ is analytic at $\kappa=0$.
(ii) If $n$ is even and $0 \notin E_{n}$, then there exists $\delta$ such that

$$
\begin{equation*}
R^{+}(\kappa)=\sum_{p, q=0}^{\infty} R_{p q}\left(\kappa^{n-2} \ln \kappa\right)^{p} \kappa^{q} \tag{1.4}
\end{equation*}
$$

where $\kappa \in\left\{\kappa \in \Gamma_{a} ;|\kappa|<\delta,-\pi / 2<\arg \kappa<3 \pi / 2\right\}, R_{p q} \in \mathscr{B}\left(L_{\rho}^{2}(\Omega)\right.$, $\left.H_{\rho^{-1}}^{2}(\Omega)\right)$ and the double series is absolutely convergent in the uniform operator norm.

Under the assumptions $\alpha(x)=1, q(x)=0$, and $\sigma=0$ the above theorem is proved in Lemma 1 of [4] and Theorem 1.2 of [2]. In this case $0 \notin E_{n}$ is automatically satisfied.

As an application of Theorem 1.1 we shall consider the local energy decay of the solution of the following wave equation:

$$
\left\{\begin{array}{l}
\partial_{t}^{2} u+L u=0 \quad \text { in } R \times \Omega  \tag{1.5}\\
B u=0 \quad \text { on } R \times \partial \Omega, \quad \partial_{t}^{j-1} u=f_{j} \text { on } t=0(j=1,2),
\end{array}\right.
$$

where we assume that all coefficients of $L$ and $B$ and $\partial \Omega$ are in the $C^{\infty}$ class. It is known that there exists a group of linear operators $U(t)$ on $H$ such that $A$ is the infinitesimal generator of $U(t)$. (See Theorem 2.7 in [8].) We remark that for any $f={ }^{f}\left(f_{1}, f_{2}\right) \in H$ the first component of $U(t) f$ satisfies (1.5) in the distribution sense. (See the proof of Lemma 3.2.) In order to show a decay of the local energy $E\left(f ; \Omega_{a}\right)=\frac{1}{2} \int_{\Omega_{a}}\left\{a_{j k} \partial_{k} f_{1} \partial_{j} \bar{f}_{1}+q\left|f_{1}\right|^{2}+\alpha^{-1}\left|f_{2}\right|^{2}\right\}+$ $\frac{1}{2} \int_{\partial \Omega} \sigma\left|f_{1}\right|^{2} d S$, where $\Omega_{a}=\Omega \cap\{|x|<a\}, f={ }^{t}\left(f_{1}, f_{2}\right) \in H$ and $\sigma=0$ if $B$ is the Dirichlet condition, we suppose a nontrapping condition of (1.5), which is stated in Definition 3.1. We note that if $\alpha(x)=1, a_{j k}(x)=\delta_{j k}$, and the complement set of $\Omega$ is convex or star-shaped, a nontrapping condition of (1.5)
holds. Making use of Vainberg's argument in [6], we can show the following local energy decay:
Theorem 1.2. We assume that $\alpha(x)-1, a_{j k}(x)-\delta_{j k}, q(x)$ belongs to $C_{0}^{\infty}(\bar{\Omega})$. If the nontrapping condition of Definition 3.1 holds and $0 \notin E_{n}$, then for any $f \in H \cap \mathscr{E}^{\prime}\left(\Omega_{a}\right)$,

$$
\begin{equation*}
E\left(U(t)(1-P) f ; \Omega_{a}\right) \leq a(t)\|f\|_{H}^{2} \quad \text { for all } t>0 \tag{1.6}
\end{equation*}
$$

where $P$ is the projection to the eigenspace associated to negative eigenvalues of $L$, and if $n$ is odd, then $a(t)=C_{1} \exp \left(-C_{2} t\right)$ with positive constants $C_{j}$ $(j=1,2)$ and if $n$ is even, then $a(t)=C_{3}(1+t)^{2(1-n)} \quad\left(C_{3}>0\right)$. Moreover if $E_{n}=\varnothing$ and $\operatorname{Ker} A=\{0\}$, then by putting $E(f)=E\left(f ; \Omega_{\infty}\right)$ we see that for $f \in H \cap \mathscr{E}^{\prime}\left(\mathbf{\Omega}_{a}\right)$

$$
\begin{equation*}
E\left(U(t) f ; \Omega_{a}\right) \leq a(t) E(f) \tag{1.7}
\end{equation*}
$$

Under the assumption $\alpha(x)=1, q(x)=0$, and $\sigma=0$, (1.7) is proved in [4] and Theorem 4.3 of [2].

## 2. The proof of Theorem 1.1

In the odd case Theorem 1.1 is proved in Theorem 4.6 of [9]. So in this section we always assume that $n$ is even and $n \geq 4$. First we shall consider the fundamental solution $R_{0}^{+}(\kappa)$ of $-\Delta+\kappa^{2}$ defined by $\left[R_{0}^{+}(\kappa) f\right](x)=$ $\int F_{\kappa}^{+}(x-y) f(y) d y$. Here by making use of the Hankel function $H_{p}^{(1)}(z)$ of the first kind with $p=(n-2) / 2, F_{\kappa}^{+}(x)$ is defined by $i(\kappa / 2 \pi|x|)^{p} H_{p}^{(1)}(\kappa|x|) / 4$. The behavior of $R_{0}^{+}(\kappa)$ near $\kappa=0$ is as follows:
Lemma 2.1. There exists $A(\kappa)$ and $B(\kappa)$ which are analytic in $\{\kappa \in C ;|\kappa|<$ a\} with values in $\mathscr{B}\left(L_{\rho}^{2}\left(R^{n}\right), H_{\rho^{-1}}^{2}\left(R^{n}\right)\right)$, where $a$ is an arbitrary positive number and $\rho(x)=e^{-a|x|}$, such that

$$
\begin{equation*}
R_{0}^{+}(\kappa)=A(\kappa) \kappa^{n-2} \ln \kappa+B(\kappa) \tag{2.1}
\end{equation*}
$$

where $[B(0) f](x)=c_{0} \int|x-y|^{2-n} f(y) d y$ with some constant $c_{0}$.
Proof. From [7] (see (5), p. 74; (1), p. 61; (2), p. 62) it follows that $F_{\kappa}^{+}(x)=$ $A_{1}(\kappa|x|) \kappa^{n-2}(\ln \kappa|x| / 2+1)+|x|^{-p} \sum_{m=0}^{p-1} c_{m}(\kappa|x|)^{2 m}+\kappa^{n-2} B_{1}(\kappa|x|)$, where $A_{1}(z)$ and $B_{1}(z)$ are entire functions such that for any $j\left|\partial_{z}^{j} A_{1}(z)\right|+\left|\partial_{z}^{j} B_{1}(z)\right|$ $\leq C_{j} e^{|z|}$. Let $A(\kappa) f=\int A_{1}(\kappa|x-y|) f(y) d y$. Then by the argument of proving Propositions 2.4 and 2.5 in [9] we see that $A(\kappa)$ is analytic in $\{\kappa \in$ $C ;|\kappa|<a\}$ with values in $\mathscr{B}\left(L_{\rho}^{2}\left(R^{n}\right), H_{p^{-1}}^{2}\left(R^{n}\right)\right)$. Similarly $B_{2}(\kappa) f=$ $\int B_{2}(\kappa|x-y|) f(y) d y$, where

$$
B_{2}(\kappa|x|)=A_{1}(\kappa|x|) \kappa^{n-2}(\ln |x| / 2+1)+|x|^{-p} \sum_{m=1}^{p-1} c_{m}(\kappa|x|)^{2 m}+\kappa^{n-2} B_{1}(\kappa|x|)
$$

has the same property. The remainder term $\left(B_{3} f\right)(x)=c_{0} \int|x-y|^{2-n} f(y) d y$ belongs to $H_{\rho^{-1}}^{1}\left(R^{n}\right)$ for $f \in L_{\rho}^{2}\left(R^{n}\right)$ and satisfies the relation $-\Delta\left(B_{3} f\right)=$ $c_{0}^{\prime} B_{3} f$ with some constant $c_{0}^{\prime}$. Thus by the argument of proving Proposition 2.4 of [9] it follows that $B_{3} \in \mathscr{B}\left(L_{\rho}^{2}\left(R^{n}\right), H_{\rho^{-1}}^{2}\left(R^{n}\right)\right)$. The proof is completed.

Next according to the argument in $\S 3$ of [9] we shall check the behavior of the resolvent operator of $-\partial_{j} a_{j k}(x) \partial_{k}$ in $R^{n}$. Let $\varphi_{m}(x)$ be a $C_{0}^{\infty}\left(R^{n}\right)$ function such that $\varphi_{m}(x)=1$ for $|x|<m, \varphi_{m}(x)=0$ for $|x|>m+1$, $0 \leq \varphi_{m}(x) \leq 1$, and for any multi-index $\beta\left|\partial_{x}^{\beta} \varphi_{m}(x)\right| \leq C_{\beta}$, where $C_{\beta}$ does not depend on $m$. Put $A=\partial_{j}\left(\delta_{j k}-a_{j k}(x)\right)\left(1-\varphi_{m}(x)\right) \partial_{k}$. Then $A$ belongs to $\mathscr{B}\left(H_{\rho^{-1}}^{2}\left(R^{n}\right), L_{\rho}^{2}\left(R^{n}\right)\right)$ and if $m$ is sufficiently large, then $\|A\|$ is sufficiently small. Put $T_{1}(\kappa)=A R_{0}^{+}(\kappa)$. Then for any $g \in L_{\rho}^{2}\left(R^{n}\right),-\left(\partial_{j} \tilde{a}_{j k}(x) \partial_{k}+\kappa^{2}\right) \times$ $R_{0}^{+}(\kappa) g=\left(1+T_{1}(\kappa)\right) g$, where $\tilde{a}_{j k}(x)=\delta_{j k}(x)+\left(a_{j k}(x)-\delta_{j k}\right)\left(I-\varphi_{m}(x)\right)$. From Lemma 2.1 it follows that $R_{1}^{+}(\kappa)=R_{0}^{+}(\kappa)\left(I+T_{1}(\kappa)\right)^{-1}$ is denoted by

$$
F\left(\kappa^{n-2} \ln \kappa, \kappa\right)=\sum_{p, q=0}^{\infty} A_{p q}\left(\kappa^{n-2} \ln \kappa\right)^{p} \kappa^{q}
$$

where $A_{p q} \in \mathscr{B}\left(L_{\rho}^{2}\left(R^{n}\right), H_{\rho^{-1}}^{2}\left(R^{n}\right)\right)$ and $\sum_{p, q=0}^{\infty}\left\|A_{p q}\right\||\lambda|^{p}|\kappa|^{q}<\infty$ if $|\lambda|$ and $|\kappa|$ are sufficiently small.

We shall consider the following problem:

$$
\begin{cases}\left(\partial_{j} a_{j k}(x) \partial_{k}+\lambda_{0}\right) v_{1}=\partial_{j}\left(\tilde{a}_{j k}-a_{j k}\right)(x) \partial_{\kappa} A_{p q} g & \text { in }|x|<N,  \tag{2.2}\\ v_{1}(x)=0 \quad \text { on }|x|=N\end{cases}
$$

where $\operatorname{Im} \lambda_{0} \neq 0, g \in L_{\rho}^{2}\left(R^{n}\right)$, and $N$ is sufficiently large. Put $V_{1}(\kappa)=$ $\sum_{p, q=0}^{\infty} V_{p q}\left(\kappa^{n-2} \ln \kappa\right)^{p} \kappa^{q}$, where $V_{p q} g=v_{1}(x)$. Then from $V_{p q} \in \mathscr{B}\left(L_{\rho}^{2}\left(R^{n}\right)\right.$, $H^{2}(\{|x|<N\})$ ) and $\left\|V_{p q}\right\| \leq C\left\|A_{p q}\right\|$, where $C$ does not depend on $p$ and $q$, we see that $\sum_{p, q=0}^{\infty}\left\|V_{p q}\right\||\lambda|^{p}|\kappa|^{q}<\infty$ if $|\lambda|$ and $|\kappa|$ are sufficiently small, and that $V_{1} \in \mathscr{B}\left(L_{\rho}^{2}\left(R^{n}\right), H^{2}(\{|x|>N\})\right)$. Now from (2.2) it follows that

$$
-\left(\partial_{j} a_{j k}(x) \partial_{k}-\kappa^{2}\right)\left(R_{1}^{+}(\kappa)+\varphi_{1}^{2} V_{1}(\kappa)\right) g=\left(I+T_{2}(\kappa)\right) g,
$$

where $\varphi_{1}(x) \in C_{0}^{\infty}(\{|x|<N\})$ such that $\varphi_{1}(x)=1$ on $\operatorname{supp}\left(a_{j k}-\tilde{a}_{j k}\right)$ and $T_{2}(\kappa)=\varphi_{1}^{2}\left(\lambda_{0}-\kappa^{2}\right) V_{1}(\kappa)-\left[\partial_{j} a_{j k} \partial_{k}, \varphi_{1}^{2}\right] V_{1}(\kappa)$, where $[A, B]=A B-B A$. We have the following:
Lemma 2. $I+T_{2}(0)$ is an invertible operator in $\mathscr{B}\left(L_{\rho}^{2}\left(R^{n}\right), L_{\rho}^{2}\left(R^{n}\right)\right)$.
Proof. Since $T_{2}(0)$ is a compact operator, we may show that $\operatorname{Ker}\left(I+T_{2}(0)\right)=$ $\{0\}$. We suppose $\left(I+T_{2}(0)\right) g=0$. Then $-\partial_{j} a_{j k}(x) \partial_{k}\left(R_{1}^{+}(\kappa)+\varphi_{1}^{2} V_{1}(0)\right) g=0$. If we put $h(x)=\left(I+T_{1}(0)\right)^{-1} g$, then $(1+|x|)^{N} h(x) \in L^{2}\left(R^{n}\right)$ for any $N$ and $R_{1}^{+}(0) g=R_{0}^{+}(0) h$. From Lemma 2.1, $R_{0}^{+}(0) h(x)=c_{0} \int|x-y|^{2-n} h(y) d y$. By $|x|^{2-n} \in L^{1}(\{|x|<1\})$ it follows that $\left|\int_{D_{1}}\right| x-\left.y\right|^{2-n} h(y) d y \mid \leq(1+|x|)^{-N} h_{1}(x)$, where $D_{1}=\{y ;|x-y|<1\}$ and $h_{1}(x) \in L^{2}\left(R^{n}\right)$. On the other hand $\left|\int_{D_{2}}\right| x-\left.\left.y\right|^{2-n} h(y) d y\right|^{2} \leq C \int_{D_{2}}|x-y|^{2 \varepsilon-n}\left(1+|y|^{-2 N} d y\right.$, where $\varepsilon$ is sufficiently small and $D_{2}=\{y ;|x-y|>1\}$. From the well-known inequality (see (2.2) in [9]) we see that $\left|R_{0}^{+}(0) h(x)\right| \leq C(1+|x|)^{\varepsilon-n / 2}+(1+|x|)^{-N} h_{1}(x)$. Similarly $\nabla R_{0}^{+}(0) h(x)$ is dominated by a function of the same type. So we can use the argument in the proof of Proposition 3.1 in [9] and can show that $g=0$. The proof is completed.

We put $R_{2}^{+}(\kappa)=\left(R_{1}^{+}(\kappa)+\varphi_{1}^{2} V_{1}(\kappa)\right)\left(I+T_{2}(\kappa)\right)^{-1}$. Then $R_{2}^{+}(\kappa)$ is the resolvent operator of $-\partial_{j} a_{j k}(x) \partial_{\kappa}-\kappa^{2}$ and is denoted by $F_{1}\left(\kappa^{n-2} \ln \kappa, \kappa\right)=$ $\sum_{p, q=0}^{\infty} B_{p q}\left(\kappa^{n-2} \ln \kappa\right)^{p} \kappa^{q}, \quad$ where $\quad B_{p q} \in \mathscr{B}\left(L_{\rho}^{2}\left(R^{n}\right), \quad H_{\rho^{-1}}^{2}\left(R^{n}\right)\right)$ and
$\sum_{p, q=0}^{\infty}\left\|B_{p q}\right\||\lambda|^{p}|\kappa|^{q}<\infty$ if $|\lambda|$ and $|\kappa|$ are sufficiently small. We shall consider the following problem:

$$
\left\{\begin{array}{l}
\left(\partial_{j} a_{j k} \partial_{k}-\lambda_{1}\right) v_{2}=\left(\partial_{j} a_{j k} \partial_{k}-\lambda_{1}\right)(1-\psi) B_{p q} E g \quad \text { in } \Omega_{1},  \tag{2.3}\\
B v_{2}=0 \quad \text { on } \partial \Omega, v_{2}=0 \text { on }|x|=N_{1}
\end{array}\right.
$$

where $\operatorname{Im} \lambda_{1} \neq 0, g \in L_{\rho}^{2}(\Omega), E g=g$ in $\Omega$ and $E g=0$ in $\Omega^{c}, \Omega_{1}=$ $\Omega \cap\left\{|x|<N_{1}\right\}$ and $\psi(x) \in C^{\infty}\left(R^{n}\right)$ such that $\psi(x)=1$ if $|x|>N_{1}-1$ and $\psi(x)=0$ near $\partial \Omega$. Put $V_{2}(\kappa)=\sum_{p, q=0}^{\infty} W_{p q}\left(\kappa^{n-2} \ln \kappa\right)^{p} \kappa^{q}+\psi R_{2}^{+}(\kappa) E$, where $W_{p q} \in \mathscr{B}\left(L_{\rho}^{2}(\Omega), H^{2}\left(\Omega_{1}\right)\right)$ is defined by $W_{p q} g=v_{2}$. We note that $\sum_{p, q=0}^{\infty}\left\|W_{p q}\right\||\lambda|^{p}|\kappa|^{q}<\infty$ if $|\lambda|$ and $|\kappa|$ are sufficiently small. From (2.3) we see that $-\left(\partial_{j} a_{j k} \partial_{k}-q+\alpha^{-1} \kappa^{2}\right)\left\{R_{2}^{+}(\kappa) E g-\varphi_{2}\left(R_{2}^{+}(\kappa) E-V_{2}(\kappa)\right) g\right\}=$ $\left(I+T_{3}(\kappa)\right) g$, where $\varphi_{2} \in C_{0}^{\infty}\left(\left\{x ;|x|<N_{1}\right\}\right)$ such that $\varphi_{2}(x)=1$ near $\Omega^{c}$ and $T_{3}(\kappa) \in \mathscr{B}\left(L_{\rho}^{2}(\Omega), L_{\rho}^{2}(\Omega)\right)$ is defined by
$\left\{\left(\alpha^{-1}-1\right) \kappa^{2}-q\right\} R_{2}^{+} E g+\left\{\left[\partial_{j} a_{j k} \partial_{k}, \varphi_{2}\right]+\varphi_{2}\left(\alpha^{-1} \kappa^{2}-\lambda_{1}-q\right)\right\}\left(R_{2}^{+}(\kappa) E-V_{2}(\kappa)\right) g$.
Lemma 2.3. If $0 \notin E_{n}$, then $I+T_{3}(0)$ is invertible.
Proof. From Proposition 4.2 in [9], $T_{3}(0)$ is a compact operator. So we may prove $\operatorname{Ker}\left(I+T_{3}(0)\right)=\{0\}$. We shall show that if $\left(I+T_{3}(0)\right) g=0$, then $R_{2}^{+}(0) E g-\varphi_{2}\left(R_{2}^{+}(0) E-V_{2}(0)\right) g$ belongs to $L^{2}(\Omega)$, if $n \geq 6$, and it belongs to $H_{D}(\Omega)$, if $n=4$. From the definition $R_{2}^{+}(0) E g-R_{0}^{+}(0) h$ has a compact support, where $h=\left(I+T_{1}(0)\right)^{-1}\left(I+T_{2}(0)\right)^{-1} E g$ and $(1+|x|)^{N} h(x) \in L^{2}\left(R^{n}\right)$ for all $N$. Since $|x|^{2-n} \in L^{2}(\{x ;|x| \geq 1\}) \cap L^{1}(\{x ;|x| \leq 1\})$ if $n \geq 6$, we see that $R_{2}^{+}(0) E g$ belongs to $L^{2}\left(R^{n}\right)$, if $n \geq 6$. By similar argument and one of Lemma 2.2, $\nabla R_{2}^{+}(0) E g \in L^{2}\left(R^{n}\right)$ and $\left|R_{2}^{+}(0) E g(x)\right| \leq C(1+|x|)^{\varepsilon-n / 2}+$ $(1+|x|)^{-N} h_{2}(x)$, where $h_{2}(x) \in L^{2}\left(R^{n}\right)$. So by the argument in the proof of Proposition 4.4 in [9] when $\kappa=0$ and $n=3 R_{2}^{+}(0) E g-\varphi_{2}\left(R^{+}(0) E-V_{2}(0)\right) g \in$ $H_{D}$ if $n=4$. It follows that if $0 \notin E_{n}$, then $R_{2}^{+}(0) E g-\varphi_{2}\left(R_{2}^{+}(0) E-V_{2}(0)\right) g=$ 0 . The argument of deriving $g=0$ from this condition is similar to one of the proof of Proposition 4.4 in [9]. The proof is completed.

Finally we put

$$
R^{+}(\kappa)=\left\{R_{2}^{+}(\kappa) E-\varphi_{2}\left(R_{2}^{+}(\kappa)-V_{2}(\kappa)\right)\right\}\left(I+T_{3}(\kappa)\right)^{-1} \alpha^{-1}
$$

where $\left(\alpha^{-1} f\right)(x)=\alpha^{-1}(x) f(x)$. Then $R^{+}(\kappa)$ has all the properties stated in Theorem 1.1.

## 3. The local energy decay

In this section we shall show the local energy decay of solutions to (1.5). Here we assume that all coefficients of $L$ and $B$ and $\partial \Omega$ are in the $C^{\infty}$ class, and the supports of $\alpha(x)-1, a_{j k}(x)-\delta_{j k}, q(x)$ are compact. First, we state the definition of a nontrapping condition. The generalized bicharacteristics of problem (1.5) are defined in Definition 3.1 in [3]. The projection into $\bar{\Omega}$ of generalized bicharacteristics is called generalized geodesics. We note that these geodesics are parameterized by time $t$.
Definition 3.1. Problem (1.5) is said to be nontrapping if for any sufficiently large $a$ there exists $T_{a}$ such that there are no generalized geodesics $\gamma(t)$ which satisfy the condition $\left\{\gamma(t) \in \bar{\Omega} ; t \in\left[0, T_{a}\right]\right\} \subset\{x ;|x|<a\}$.

It is known that if $\alpha(x)=1, a_{j k}(x)=\delta_{j k}$, and $\Omega^{c}$ is convex or starshaped, then problem (1.5) is nontrapping. In this section we always assume this nontrapping condition.

Lemma 3.2. Assume that the support of $f \in H_{D}$ is contained in $\Omega_{a}=\Omega \cap\{x$; $|x|<a\}$. If $a$ is sufficiently large, then $[U(t) f](x) \in C^{\infty}\left(D_{a}\right)$, where $D_{a}=$ $\left\{(t, x) \in R \times \bar{\Omega} ; t>T_{a},|x|<t-T_{a}+a\right\}$.
Proof. Let $u(t, x)$ be the first component of $[U(t) f](x)$. Since $D(A)$ is a dense set in $H$, we see that $\partial_{t}^{2} u+L u=0$ as $\mathscr{D}^{\prime}(R \times \Omega)$. From this fact, Theorem 4.3.1, and Theorem 2.5.6 of [1] it follows that the trace of $B u$ on $R \times \partial \Omega$ exists as an element of $\mathscr{D}^{\prime}(R \times \partial \Omega)$. From $B u=0$ on $R \times \partial \Omega$ for $f \in D(A)$ by approximating elements of $D(A)$ to $u$, we see that $B u=0$ on $R \times \partial \Omega$. Thus we can use theorems on a propagation of singularities for $u(t, x)$. From the finite propagation property and the fact that $D(A)$ is a dense subset of $H, u(t, x)=0$ in $\{(t, x) ;|x|>|t|+a\}$. Let $\gamma(t)$ be a generalized geodesic such that $\left|\gamma\left(t_{0}\right)\right|<t_{0}-T_{a}+a$ and $t_{0}>T_{a}$. Then from $L=-\Delta$ in $\{x ;|x|>a\}$ and the nontrapping condition it follows that $|\gamma(0)|>a$. Thus $u(t, x)$ is $C^{\infty}$ near $\{(t, \gamma(t)) ;|t|<\delta\}$, where $\delta$ is sufficiently small. By theorems on a propagation of singularities (see for example Theorem 5.10 in [3]) it follows that $u(t, x)$ is $C^{\infty}$ near $\left(t_{0}, \gamma\left(t_{0}\right)\right)$. We note that the condition $(2.2)_{ \pm}$in [3] is valid for the third boundary condition. The proof is completed.

Next we shall state the behavior of $R^{+}(\kappa)$ near $|\operatorname{Re} \kappa|=\infty$.
Theorem 3.3. Assume that problem (1.5) is nontrapping. Then there exist positive constants $\alpha$ and $\beta$ such that $R^{+}(\kappa)$ is analytic in $U_{\alpha, \beta}=\{\kappa ;-\pi / 2<\arg \kappa<$ $3 \pi / 2,|\operatorname{Im} \kappa|<\alpha \ln |\operatorname{Re} \kappa|-\beta\}$ with values in $\mathscr{B}\left(L_{a}^{2}(\Omega), H^{2}\left(\Omega_{a}\right)\right)$, where $L_{a}^{2}(\Omega)=L^{2}(\Omega) \cap \mathscr{E}^{\prime}\left(\Omega_{a}\right)$ with $\Omega_{a}=\Omega \cap\{|x|<a\}$. Moreover there exist positive constants $C$ and $T$ such that for $j=0,1,2, \kappa \in U_{\alpha, \beta}$,

$$
\begin{equation*}
\left\|R^{+}(\kappa) f\right\|_{H^{2-j}\left(\Omega_{a}\right)} \leq C|\kappa|^{1-j} e^{T|\operatorname{Im} \kappa|}\|f\|_{L^{2}\left(\Omega_{a}\right)} \tag{3.1}
\end{equation*}
$$

The statement of the above theorem is the same as Theorem 7 in [6]. However, in order to prove it the author of [6] assumes an existence of the Green function of problem (1.5) which is defined in condition $D^{\prime}$ in [6, p. 11]. Since there is no guarantee on the existence of it, we shall give a proof without assuming the existence of the Green function. But the outline is almost the same as the proof of Theorem 7 in [6].

The following is a sketch of the proof of Theorem 3.3 without assuming the existence of the Green function. Let $h(t, x)$ be a $C^{\infty}\left(R^{n+1}\right)$ function such that $h(t, x)$ depends only on $t$ in a neighborhood of $\partial \Omega, h(t, x)=1$ in the complement set of $\left\{(t, x) ; t>T_{a}+1,|x|<t-\left(T_{a}+1\right)+a\right\}$ and $h(t, x)=0$ in $\left\{(t, x) ; t>T_{a}+2,|x|<t-\left(T_{a}+2\right)+a\right\}$. Put $(E \varphi)(t, x)=h(t, x) u(t, x)$, where $u(t, x)$ is the first component of $U(t)^{t}(0, \varphi)$ and define $(\widetilde{E}(\kappa) \varphi)=$ $\int_{0}^{\infty} e^{i \kappa t}(E \varphi)(x, t) d t$. Then we have the following:

Lemma 3.4. The operator $\tilde{E}(\kappa)$ is an entire function of $\kappa$ with values in $\mathscr{B}\left(H_{a}^{s}(\Omega), H^{s+2}\left(\Omega_{a}\right)\right)(s=0,1)$ and there exist positive constants $C$ and $T$ such that for $s=0,1$ and $j=0,1,2$,

$$
\begin{equation*}
\|\widetilde{E}(\kappa) \varphi\|_{H^{s+2-j}\left(\Omega_{a}\right)} \leq C|\kappa|^{1-j} e^{T|\operatorname{Im} \kappa|}\|\varphi\|_{H^{s}\left(\Omega_{a}\right)} \tag{3.2}
\end{equation*}
$$

Proof. By Lemma 3.2 and the closed graph theorem the mapping from $L_{a}^{2}(\Omega)$ to $u(t, x) \in C^{\infty}(D)$, where $D=\left\{(t, x) \in R \times \bar{\Omega}: t>T_{a},|x|<t-T_{a}+a\right\}$, is continuous. So for any $j, \beta$, and $D_{1} \Subset D$ there exists a constant $C=$ $C\left(j, \beta, D_{1}\right)$ so that

$$
\begin{equation*}
\sup _{D_{1}}\left|\left(\partial_{x}^{\beta} \partial_{t}^{j} u\right)(t, x)\right| \leq C\|\varphi\|_{L^{2}\left(\Omega_{a}\right)} \tag{3.3}
\end{equation*}
$$

By $\partial_{t}(E \varphi)(0, x)=\varphi$ and integration by parts it follows that

$$
\begin{equation*}
\left(L-\kappa^{2}\right) \widetilde{E} \varphi=-\varphi+\int_{0}^{\infty} f(t, x) e^{i \kappa t} d t \tag{3.4}
\end{equation*}
$$

where $f=\left[\partial_{t}^{2}+L, h\right] u$ belongs to $C^{\infty}(R \times \bar{\Omega})$ from (3.3). Making use of the properties of $U(t)$, the continuous inclusion $H_{a}^{1}(\Omega) \subset H^{1}\left(\Omega_{a}\right)$, and the elliptic estimate for $L(u(t, \cdot))=-\left(\partial_{t}^{2} u\right)(t, x)$, we can prove that $\partial^{j} E / \partial t^{j} \in$ $C\left(R ; \mathscr{B}\left(H_{a}^{s}(\Omega), H^{s+1-j}\left(\Omega_{a}\right)\right)\right)$, where $s=0,1$ and $0 \leq j \leq s+1$. Now the estimate (3.2) is derived from the above property of $E$, (3.3), the elliptic estimate for (3.4), and $\widetilde{E} \varphi=i \kappa^{-1} \int_{0}^{\infty} e^{i \kappa t} \partial_{t}(E \varphi) d t$. The proof of Lemma 3.4 is completed.

From Lemma 3.2 we can extend $u(t, x)$ near $D_{2}=\left[T_{a}+1, T_{a}+2\right] \times\{x \in$ $\left.R^{n}:|x|<a\right\}$ to be an element of $C^{\infty}\left(D_{2}\right)$. Here we may assume that the extended function $\tilde{u}(t, x)$ also satisfies (3.3) for any compact subset $D_{1}$ of $\left.\{(t, x) \in R \times \bar{\Omega}): t>T_{a},|x|<t-T_{a}+a\right\} \cup D_{2}$. Let us consider the solution of the wave equation $\left(\partial_{t}^{2}-\Delta\right) w=-\tilde{f}$ in $R^{n}, \partial_{t}^{j} w=0$ on $t=0 \quad(j=0,1)$, where $\tilde{f}=\left[\partial_{t}^{2}+L, h\right] \tilde{u} \in C^{\infty}\left(R^{n+1}\right)$. Put $\widetilde{V}(\kappa) \varphi=\int_{0}^{\infty} e^{i \kappa t} w(t, x) d t$. Then we have the following:
Lemma 3.5. $\widetilde{V}(\kappa)$ is an analytic function of $K$ in $D=\{\kappa \in C \backslash\{0\}:=\pi / 2<$ $\arg \kappa<3 \pi / 2\}$ with values in $\mathscr{B}\left(L_{a}^{2}\left(R^{n}\right), H^{s}(\{x:|x|, a\})\right)$, where $L_{a}^{2}\left(R^{n}\right)=$ $L^{2}\left(R^{n}\right) \cap \mathscr{E}^{\prime}(\{|x|<a\})$ and $s$ is arbitrary integer. Moreover there exist constants $C_{s}$ and $T$ so that for any $\kappa \in D$ there exist constants $C_{s}$ and $T$ so that for any $\kappa \in D$

$$
\begin{equation*}
\|\tilde{V}(\kappa) \varphi\|_{H^{s}(\{|x|<a\})} \leq C_{s}|\kappa|^{-2} e^{T|\operatorname{Im} \kappa|}\|\varphi\|_{L^{2}\left(R^{n}\right)} \tag{3.5}
\end{equation*}
$$

Proof. Let $\psi(x) \in C^{\infty}\left(R^{n}\right)$ such that $\psi(x)=0$ if $|x|<a_{1} \quad\left(a_{1}>a\right), \psi(x)=$ 1 if $|x|>A_{1}$, and $L=-\Delta$ in supp $\psi$. Put $v(t, x)=\psi E \varphi+w$. Then $v$ satisfies the Cauchy problem $\left(\partial_{t}^{2}-\Delta\right) v=\psi f-\tilde{f}-[\psi, \Delta] E \varphi$ in $[0, \infty) \times R^{n}, \partial_{t}^{j} v=0$ on $t=0,(j=0,1)$. Here the support of $\psi f-\tilde{f}-[\psi, \Delta] E \varphi$ is contained in $\left\{(t, x):|x|<M_{0}, 0<t<T_{0}\right\}$ for some $M_{0}$ and $T_{0}$. From the finite progagation property of $v$ it follows that $\left(\partial_{t}^{2}-\Delta\right) v=0$ in $\left[T_{0}, \infty\right) \times R^{n}, \partial_{t}^{j} v=\varphi_{j}$ on $t=T_{0} \quad(j=0,1)$, where the supports of $\varphi_{j}$ are contained in $\{|x|<M\}$. Making use of the inequality $\sup _{t \in[0, T]}\left\|\partial_{t}^{j} w(t, \cdot)\right\|_{H^{s}\left(R^{n}\right)} \leq C_{T, j, s}\|\varphi\|_{L^{2}(\Omega)}$ which is derived from (3.3) and a strictly hyperbolic estimate for $w$, and the continuous inclusion $H_{a}^{1}(\Omega) \subset H^{1}\left(\Omega_{a}\right)$, we see that $\left\|\varphi_{j}\right\|_{L^{2}\left(R^{n}\right)} \leq C\|\varphi\|_{L^{2}(\Omega)} \quad(j=0,1)$. If $n$ is odd, by the existence of a lacuna for $v, w(t, x)=0$ for $t>T_{1}$ and $|x|<a$. If $n$ is even, $v=d\left(E_{x}(t) * \varphi_{0}\right) / d t+E_{x}(t) * \varphi_{1}$, where,

$$
E_{x}(t) * \varphi_{1}(2 \pi)^{-p}\left(\frac{1}{t} \frac{d}{d t}\right)^{p-1} \int_{|y| \leq t} \varphi_{1}(x-y) /\left(t^{2}-|y|^{2}\right)^{1 / 2} d y
$$

with $p=n / 2$. If $t>T_{1}=a+M$ and $|x|<a$, then

$$
\int_{|y| \leq t} \varphi_{1}(x-y) /\left(t^{2}-|y|^{2}\right)^{1 / 2} d y=\int_{|y| \leq M} \varphi_{1}(y) /\left(t^{2}-|x-y|^{2}\right)^{1 / 2} d y
$$

It follows that for $\operatorname{Re} t>T_{1},|x|<a, w(t, x)$ is infinitely differentiable jointly with respect to $t, x$ is analytic in $t$ and satisfies sup $\left|\partial_{t}^{j} \partial_{x}^{\alpha} w\right| \leq C_{j, \alpha} t^{-3}\|\varphi\|_{L^{2}(\Omega)}$, where the supremum is taken in $|x|<a$. Thus by the argument in the proof of Lemma 6 of [6] we have the desired properties on $\widetilde{V}(\kappa)$. The proof is completed.

Let $\psi_{1}(x) \in C^{\infty}\left(R^{n}\right)$ such that $\psi_{1}(x)=1$ for $|x|>a$ and $\psi_{1}(x)=0$ near $\partial \Omega$, and put $\widetilde{W}(\kappa) \varphi=\widetilde{E}(\kappa) \varphi+\psi_{1} \widetilde{V}(\kappa) \varphi$. Then $\left(L-\kappa^{2}\right) \widetilde{W} \varphi=-(I+T(\kappa)) \varphi$ and $B \widetilde{W}(\kappa) \varphi=0$, where $T(\kappa) \varphi=-(L+\Delta) \widetilde{V} \varphi-\left(L-\kappa^{2}\right)\left(\psi_{1}-1\right) \widetilde{V} \varphi \in$ $L_{a}^{2}(\Omega)$. From (3.2) and (3.5) there exist positive constants $\alpha, \beta$ such that $\|T(\kappa)\|_{\mathscr{B}\left(L_{\alpha}^{2}(\Omega), L_{\alpha}^{2}(\Omega)\right)} \leq 1 / 2$ for $\kappa \in U_{\alpha, \beta}$. Put $\widetilde{R}(\kappa)=-\widetilde{W}(\kappa)(I+T(\kappa))^{-1}$, which satisfies (3.1). The final problem to complete the proof of Theorem 3.3 is to show the equality $R^{+}(\kappa) \varphi=\widetilde{R}(\kappa) \varphi$ for $\varphi \in L_{a}^{2}(\Omega)$. In order to prove this we may show the following:
Lemma 3.6. There exists a positive constant $A$ such that for any $\varphi \in L_{a}^{2}(\Omega)$ and $\kappa$ with $\operatorname{Im} \kappa>A, \widetilde{R}(\kappa) \varphi \in H^{2}(\Omega)$.
Proof. Let $\psi_{1}(x)$ be the function that appeared in the definition of $\widetilde{W}(\kappa)$ and put $v_{1}(t, x)=\psi_{1}(E \varphi+w)$, which satisfies $\left(\partial_{t}^{2}+L\right) v_{1}=F(t, x)$, where $F=$ $\psi_{1}\left(\partial_{t}^{2}+L\right)(E \varphi+w)+\psi_{1}(\Delta-L) w$. From the properties of $U(t)$ and a strictly hyperbolic estimate for $w$ that appeared in the proof of Lemma 3.5 it follows that for all $t \geq 0\|F(t, \cdot)\|_{L^{2}\left(R^{n}\right)} \leq C_{\varphi}$. By the proof of Theorem 3.8 in [8], $u(t, x)=[U(t) \varphi]_{1}(x)=\sum_{j=1}^{m} \lambda_{j}^{-1}\left(\varphi, p_{j}\right)_{\mathscr{E}}\left(e^{\lambda_{j} t}-e^{-\lambda_{j} t}\right) p_{j} / 2+L_{+}^{-1 / 2} \sin t L_{+}^{1 / 2} \varphi_{1}$, where $\left\{-\lambda_{j}^{2}\right\}$ are negative eigenvalues of $L,\left\{p_{j}\right\}$ are linearly independent eigenvectors of $\left\{-\lambda_{j}^{2}\right\}, L_{+}=\int_{0}^{\infty} \lambda d E(\lambda)$ with the spectral resolution $\{E(\lambda)\}$ of $L$, and $\varphi_{1}=\varphi-\sum_{j=1}^{m}\left(\varphi, p_{j}\right)_{\mathscr{H}} p_{j}$. Thus we have $v_{1} \in C\left(R ; H^{1}\left(R^{n}\right)\right)$ and $\partial_{t} v_{1} \in C\left(R ; L^{2}\left(R^{n}\right)\right)$. Making use of an estimate for hyperbolic operator on $v_{1}$, we see that there exist $A>0$ and $\gamma>0$ so that $e^{-t A}\left\|v_{1}(t, \cdot)\right\|_{H^{1}\left(R^{n}\right)} \leq$ $C\left\{\left\|\psi_{1} \varphi\right\|_{L^{2}\left(R^{n}\right)}+\int_{0}^{t} e^{-s \gamma}\|F(s, \cdot)\|_{L^{2}\left(R^{n}\right)} d s\right\}$. Thus $\widetilde{R}(\kappa) \varphi \in H^{1}(\Omega)$ if $\operatorname{Im} \kappa>A$. Since $\widetilde{R}(\kappa) \in H_{\mathrm{loc}}^{1}(\bar{\Omega})$, from the Fourier transform and $\left(\Delta-\kappa^{2}\right) \widetilde{R}(\kappa) \varphi \in L^{2}(\Omega)$ it follows that $\widetilde{R}(\kappa) \varphi \in H^{2}(\Omega)$. The proof is completed.
Proof of Theorem 1.2. We use the argument in the proof of Theorem 8 in [6]. The solution $u(t, x)$ of (1.5) is equal to $\int_{i \gamma-\infty}^{i \gamma+\infty} R^{+}(\kappa) \varphi e^{-i \kappa t} d \kappa / 2 \pi$, where $\left(f_{1}, f_{2}\right)=(0, \varphi)$ with $\varphi \in D(L) \cap L_{a}^{2}(\Omega)$ and $\gamma$ is sufficiently large. We note that the integral converges from the equality $R^{+}(\kappa) \varphi=\kappa^{-2}\left(R^{+}(\kappa) L \varphi-\varphi\right)$. By the unique continuation property of solutions to $\left(L-\kappa^{2}\right) u=0$, where $\kappa \in R$, and Theorem 4.7(a) in [9] it follows that any point in $R \backslash\{0\}$ is an analytic point of $R^{+}(\kappa)$. By deforming the contour of the integration, we see that $u(t, x)=i \sum_{j=1}^{m} \operatorname{Re} s\left[R^{+}(\kappa) \varphi e^{-\kappa t} ; i \lambda_{j}\right]-\int_{\Gamma} R^{+}(\kappa) \varphi e^{-i \kappa t} d \kappa$, where $t$ is sufficiently large and $\Gamma$ is the curve shown in Figures 1 or 2 in [6]. From $R^{+}(\kappa)=\left(L-\kappa^{2}\right)^{-1}$ for $0<\arg \kappa<\pi$ it is not difficult to show that $i \operatorname{Res}\left[R^{+}(\kappa) \varphi e^{-i \kappa t} ; i \lambda_{j}\right]=e^{\lambda_{j} t} \sum_{l}\left(\varphi, p_{l}\right)_{\mathscr{H}} p_{l} / 2 \lambda_{j}$, where $\left\{p_{l}\right\}$ is a base of the
eigenvector space of the negative eigenvalues $-\lambda_{j}^{2}$ of $L$. From Theorems 1.1, 3.3, and Theorem 8 and Lemma 9 of [6] we have (1.6) for $f=(0, \varphi)$ with $\varphi \in D(L) \cap L_{a}^{2}(\Omega)$. We note that the solution $u(t, x)$ of (1.5) is equal to the first component of $\partial_{t} U(t)^{t}\left(0, f_{1}\right)+U(t)^{t}\left(0, f_{2}\right)$. Moreover if $E_{n}=\varnothing$ and $\operatorname{Ker} A=\{0\}$, then we can show that $\|f\|_{H_{D}}^{2}$ is equivalent to $E(f)$. We complete the proof of Theorem 1.2.

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Department of Mathematics, Nagoya Institute of Technology, Nagoya, Gokiso-cho, 466, JAPAN

