STRONG MOMENT PROBLEMS FOR RAPIDLY DECREASING SMOOTH FUNCTIONS

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ABSTRACT. It is shown that the existence of rapidly decreasing smooth solutions of various moment problems follows from the theorem of Ritt on the existence of analytic functions with a prescribed asymptotic power series at the vertex of a given sector.

1. Introduction

The problem of moments, as well as its many generalizations, have been the object of research for over a century. In its basic form it consists of finding a positive function $\phi(x)$, defined on a given interval I, that satisfies

(1.1)
$$\int_{I} x^{n} \phi(x) dx = \mu_{n}, \qquad n \in \mathbb{N},$$

where $\{\mu_n\}_{n\in\mathbb{N}}$ is a given sequence of real or complex numbers. We refer to [1, 17] for an account of the classical aspects of this problem.

The moments play a rather important role in several areas of current interest, including the theory of asymptotic expansions of generalized functions [7, 8, 10], the theory of orthogonal polynomials [12, 14], and the theory of distributional solutions of differential and functional equations [13, 19].

Recently, Durán [6] proved that when $I = [0, \infty)$, for an arbitrary sequence $\{\mu_n\}_{n\in\mathbb{N}}$, problem (1.1) admits solutions of the class $\mathscr{S}(0, \infty) = \{\phi \in \mathscr{S}: \phi(x) = 0, x \leq 0\}$. Here $\mathscr{S} = \mathscr{S}(\mathbb{R})$ is the Schwartz space of rapidly decreasing smooth functions [16]; namely,

$$\mathcal{S}(\mathbb{R}) = \left\{ \phi \in C^{\infty}(\mathbb{R}) \colon \lim_{|x| \to \infty} x^n \phi^{(m)}(x) = 0 \,,\, n \,,\, m \in \mathbb{N} \right\}.$$

The arbitrariness of the sequence $\{\mu_n\}_{n\in\mathbb{N}}$ had already been obtained by Boas [4], who worked in the class of functions of bounded variation. This contrasts with the situation when I is a compact interval, since growth conditions on certain linear combinations of the μ_n 's are required for existence, even in the case when ϕ is a distribution [9].

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In the present article we show that this result as well as the existence of solutions of more general problems such as the strong moment problem,

(1.2)
$$\int_{-\infty}^{\infty} x^n \phi(x) \, dx = \mu_n, \qquad n \in \mathbb{Z}.$$

can be obtained from the following theorem of Ritt [18, p. 41]:

Theorem. Let $\{b_n\}_{n \in \mathbb{N}}$ be a sequence of complex numbers. Then there exists a bounded analytic function F(z) defined in the sector $S_a(\alpha, \beta) = \{re^{i\theta} : 0 < r < a, \alpha < \theta < \beta\}$ with the asymptotic power series

(1.3)
$$F(z) \sim \sum_{n=0}^{\infty} b_n z^n, \qquad z \to 0, \ z \in S_a(\alpha, \beta). \quad \Box$$

Actually, the function F(z) can be chosen analytic and bounded in the whole sector $S(\alpha, \beta) = S_{\infty}(\alpha, \beta) = \{re^{i\theta} : r > 0, \alpha < \theta < \beta\}$. This can be obtained from the result for $S_a(\alpha, \beta)$, $a < \infty$, as follows. Write F(z) = f(h(z)), where $h: S(\alpha, \beta) \to S_a(\alpha, \beta)$ is a conformal map with h(0) = 0 and $h(z) \sim z + c_2 z^2 + c_3 z^3 + \cdots$, as $z \to 0$, and where $f(\omega)$ is chosen as in the theorem with asymptotic development $f(\omega) \sim d_0 + d_1 \omega + d_2 \omega^2 + \cdots$, as $\omega \to 0$, with the d_i 's chosen in such a way that, formally,

$$\sum_{j=0}^{\infty} d_j \left(z + \sum_{n=2}^{\infty} c_n z^n \right)^j = \sum_{n=0}^{\infty} b_n z^n.$$

2. Solution of the moment problem

In this section we use the theorem of Ritt to give a proof of the existence of solutions in the class $\mathcal{S}(0,\infty)$ of the problem

(2.1)
$$\int_0^\infty x^n \phi(x) \, dx = \mu_n, \qquad n \in \mathbb{N},$$

where $\{\mu_n\}_{n\in\mathbb{N}}$ is an arbitrary sequence. This result was obtained by Durán [6], who gave a different proof.

We first recall [2] that a function $\psi \in \mathcal{S}(\mathbb{R})$ is the Fourier transform

(2.2)
$$\psi(u) = \hat{\phi}(u) = \int_0^\infty e^{ixu} \phi(x) \, dx$$

of a function ϕ of the class $\mathcal{S}(0,\infty)$ if and only if it can be extended to a bounded continuous function $\Psi(z)$ in the upper half plane $\operatorname{Im} z \geq 0$, analytic in $\operatorname{Im} z > 0$ and vanishing as $z \to \infty$.

Since

(2.3)
$$\frac{d^n\hat{\phi}(0)}{du^n} = i^n \int_0^\infty x^n \phi(x) \, dx \,,$$

problem (2.1) is equivalent to that of finding $\psi = \hat{\phi}$ in the class $\mathscr{F}(\mathscr{S}(0,\infty))$ that satisfies

(2.4)
$$\psi^{(n)}(0) = i^n \mu_n, \qquad n \in \mathbb{N}.$$

We construct ψ as follows. Let

(2.5)
$$G(z) = e^{(1-i)(z+i)^{1/2}}.$$

where the branch of the square root is chosen so as to assure that $G(z) \to 0$ as $z \to \infty$ in the upper half plane Im $z \ge 0$. Let the sequence a_0, a_1, a_2, \ldots be defined as

(2.6)
$$\frac{1}{G(z)} = a_0 + a_1 z + a_2 z^2 + \cdots, \qquad |z| < 1.$$

Let $\delta \in (0, \pi/2)$, and let F(z) be a bounded analytic function in the sector $S: -\delta < \operatorname{Arg} z < \pi + \delta$ with the asymptotic power series

$$(2.7) F(z) \sim b_0 + b_1 z + b_2 z^2 + \cdots, z \to 0, z \in S,$$

where

(2.8)
$$b_n = \sum_{k=0}^n \frac{i^k}{k!} \mu_k a_{n-k}.$$

Let $\Psi(z) = F(z)G(z)$. Then Ψ is analytic in the sector S, and $\Psi(z) = o(z^{-n})$ for each $n \in \mathbb{N}$ as $z \to \infty$ within S. Also,

(2.9)
$$\Psi(z) \sim \sum_{n=0}^{\infty} \frac{i^n \mu_n z^n}{n!}, \qquad z \to 0, z \in S.$$

It follows that, if ψ is the restriction of Ψ to the real axis, then $\psi \in \mathscr{F}(\mathcal{S}(0,\infty))$ and ψ satisfies (2.4), as required.

3. The strong moment problem

In this section we consider the following strong moment problem: Find $\phi \in \mathcal{S}(\mathbb{R})$ that satisfies $\phi^{(n)}(0) = 0$, $n \in \mathbb{N}$, and

(3.1)
$$\int_{-\infty}^{\infty} x^n \phi(x) dx = \mu_n, \qquad n \in \mathbb{Z},$$

where $\{\mu_n\}_{n\in\mathbb{Z}}$ is an arbitrary sequence. We refer to [3, 11, 15] for the study of this problem in the class of the positive Radon measures.

Let $\psi(u) = \hat{\phi}(u)$. Then $\psi \in \mathcal{S}(\mathbb{R})$, and (3.1) is equivalent to the two sets of conditions

(3.2a)
$$\psi^{(n)}(0) = i^n \mu_n, \qquad n = 0, 1, 2, ...,$$

(3.2b)
$$\int_0^\infty \mu^n \psi(\mu) \, d\mu = n! i^{n+1} \mu_{-n-1}, \qquad n = 0, 1, 2, \dots.$$

Using the construction of §2, we can find $\psi_1 \in \mathcal{S}(\mathbb{R})$ such that $\psi_1^{(n)}(0) = i^n \mu_n$, $n = 0, 1, 2, \ldots$. We can also find $\psi_2 \in \mathcal{S}(0, \infty)$ such that

(3.3)
$$\int_0^\infty \mu^n \psi_2(\mu) \, d\mu = n! i^{n+1} \mu_{-n-1} - \int_0^\infty \mu^n \psi_1(\mu) \, d\mu.$$

Then $\psi = \psi_1 + \psi_2$ belongs to $\mathscr{S}(\mathbb{R})$ and satisfies (3.2a,b). However, if $\hat{\phi} = \psi$, the condition $\phi^{(n)}(0) = 0$, $n \in \mathbb{N}$, might not be satisfied. Following a suggestion of Professor C. Berg, we solve this by choosing $\psi_3 \in \mathscr{S}(-\infty, 0)$ such that

(3.4)
$$\int_{-\infty}^{\infty} y^n \psi_3(y) \, dy = -n! i^{n+1} \mu_{-n-1} - \int_{-\infty}^{0} y^n \psi_1(y) \, dy \,,$$

for $n \in \mathbb{N}$, and define $\psi = \psi_1 + \psi_2 + \psi_3$. This function ψ satisfies (3.2a, b) but also

(3.5)
$$\int_{-\infty}^{\infty} y^n \psi(y) \, dy = 0, \qquad n \in \mathbb{N},$$

and thus if $\hat{\phi} = \psi$, then $\phi^{(n)}(0) = 0$, $n \in \mathbb{N}$.

Actually, the strong moment problem can be solved in $\mathcal{S}(0, \infty)$; that is, there are $\sigma \in \mathcal{S}(0, \infty)$ with

(3.6)
$$\int_0^\infty x^n \sigma(x) \, dx = \mu_n, \qquad n \in \mathbb{Z}.$$

To see it, let $\{\mu_n\}_{n\in\mathbb{Z}}$ be any sequence, and let $\phi\in\mathscr{S}(\mathbb{R})$ with $\phi^{(n)}(0)=0$, $n\in\mathbb{N}$, be a solution of the strong problem

(3.7)
$$\int_{-\infty}^{\infty} x^n \phi(x) dx = \begin{cases} \mu_{n/2}, & n \text{ even,} \\ 0, & n \text{ odd.} \end{cases}$$

Let $\rho \in \mathcal{S}(0, \infty)$ be given by $\rho(x) = \phi(x) + \phi(-x)$, $x \ge 0$. Then

(3.8)
$$\int_0^\infty u^{2n} \rho(u) du = \mu_n, \qquad n \in \mathbb{Z}.$$

Set

(3.9)
$$\sigma(x) = \frac{\rho(x^{1/2})}{2x^{1/2}}, \qquad x \ge 0.$$

Then $\sigma \in \mathcal{S}(0, \infty)$, and it satisfies (3.6).

Summarizing, we have shown the following result:

Theorem. Let $\{\mu_n\}_{n\in\mathbb{Z}}$ be an arbitrary sequence. Then there exist functions $\phi \in \mathcal{S}(0,\infty)$ that satisfy

(3.10)
$$\int_0^\infty x^n \phi(x) \, dx = \mu_n$$

for each $n \in \mathbb{Z}$.

4. FURTHER REMARKS

Our results imply the existence of smooth functions $\phi_k \in \mathcal{S}(0, \infty)$, $k = 0, 1, 2, \ldots$, that satisfy the moment problem

(4.1)
$$\int_0^\infty x^n \phi_k(x) \, dx = \delta_{nk} \,,$$

where $\delta_{nk} = 1$ if n = k, $\delta_{nk} = 0$ if $n \neq k$ is the Kronecker delta. The functions $\phi_k(x)$ can be used as "smooth versions" of the Dirac delta functions $(-1)^k \delta^{(k)}(x)/k!$, which also satisfy

(4.2)
$$\left\langle \frac{(-1)^k \delta^{(k)}(x)}{k!}, x^n \right\rangle = \delta_{nk}.$$

In particular, if $\psi \in \mathcal{S}(\mathbb{R})$ has moments $\mu_n = \langle \psi(x), x^n \rangle$, $n = 0, 1, 2, \ldots$, then we could expect some relationship between $\psi(x)$ and the series

 $\sum_{n=0}^{\infty} \mu_n \phi_n(x)$. Although the series is not necessarily equal to ψ (it could even be divergent), we do have the asymptotic relation

$$(4.3) \psi(\lambda x) \sim \mu_0 \phi_0(\lambda x) + \mu_1 \phi_1(\lambda x) + \mu_2 \phi_2(\lambda x) + \cdots , \quad \text{as } \lambda \to \infty$$

in the space $\mathscr{S}'(\mathbb{R})$. The interpretation of (4.3) is in the distributional sense; namely, that for each $\rho \in \mathscr{S}(\mathbb{R})$ we have (4.4)

$$\int_{-\infty}^{\infty} \psi(\lambda x) \rho(x) \, dx = \sum_{j=0}^{n} \mu_j \int_{-\infty}^{\infty} \phi_j(\lambda x) \rho(x) \, dx + O(\lambda^{-n-2}), \quad \text{as } \lambda \to \infty.$$

Relation (4.3) can be interpreted as saying that $\{\phi_k\}_{k\in\mathbb{N}}$ is an "asymptotic basis" of the space $\mathcal{S}(\mathbb{R})$. Actually, (4.3) is a smooth version of the moment asymptotic expansion [10]

(4.5)
$$f(\lambda x) \sim \sum_{n=0}^{\infty} \frac{(-1)^n \mu_n \delta^{(n)}(x)}{n! \lambda^{n+1}}, \quad \text{as } \lambda \to \infty,$$

valid if f is a generalized function of distributional rapid decay at infinity, where the $\mu_n = \langle f(x), x^n \rangle$ are the moments.

We would like to remark that our results imply that the class \mathcal{M} of entire functions of the form

(4.6)
$$\Phi(z) = \int_0^\infty x^z \phi(x) \, dx,$$

where $\phi \in \mathscr{S}(0,\infty)$, satisfies the following interpolation property: for each sequence $\{\mu_n\}_{n\in\mathbb{Z}}$ there are $\Phi \in \mathscr{M}$ that satisfy

$$\Phi(n) = \mu_n, \qquad n \in \mathbb{Z}.$$

The class \mathcal{M} is larger than the class of entire functions of exponential type [5]. Indeed, if $\Phi \in \mathcal{M}$ then there exists $\sigma \in C^{\infty}(\mathbb{R})$ such that $\sigma(t) = o(e^{-n|t|})$ as $|t| \to \infty$ for each $n \in \mathbb{N}$, with

(4.8)
$$\Phi(z) = \int_{-\infty}^{\infty} e^{zt} \sigma(t) dt.$$

For functions of exponential type, the $\sigma(t)$ has compact support. Actually, there are sequences $\{\mu_n\}_{n\in\mathbb{Z}}$ for which (4.7) does not have solutions of exponential type.

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