# UNIQUENESS OF SOLUTIONS OF MEAN FIELD EQUATIONS IN $R^{2}$ 

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#### Abstract

In this paper, we prove uniqueness of solutions of mean field equations with general boundary conditions for the critical and subcritical total mass regime, extending the earlier results for null Dirichlet boundary condition. The proof is based on new Bol's inequalities for weak radial solutions obtained from rearrangement of the solutions.


## 1. Introduction

Let $\Omega \subset \mathbb{R}^{2}$ be an open bounded domain and consider the mean field equation

$$
\begin{cases}\Delta u+\rho \frac{e^{u}}{\int_{\Omega} e^{u}}=0 & \text { in } \Omega  \tag{1}\\ u=0 & \text { on } \partial \Omega\end{cases}
$$

Suzuki [23] proved that if $\Omega$ is simply-connected, then for $0<\rho<8 \pi$ the equation (11) has a unique solution. Later in [8] the authors extended this result to the case $\rho=8 \pi$. Recently in [5] Bartolucci and Lin extended the result to multiplyconnected domains. Indeed they proved the following.

Theorem A (Theorem 2 in [5). Let $\Omega$ be an open, bounded, and multiply-connected domain of class $C^{1}$. Then equation (11) admits at most one solution for $0<\rho \leq 8 \pi$.

The proof relies on a generalization of the classical Bol's inequality for multiplyconnected domains (see Theorem C below). A necessary and sufficient condition for the existence of a solution at the critical parameter $\rho=8 \pi$ is also provided in [5].

In this paper, among other results, we study uniqueness of solutions of the general mean field equation

$$
\begin{cases}\Delta u+\rho \frac{K(x) e^{u}}{\int_{\Omega} K(x) e^{u}}=f & \text { in } \Omega,  \tag{2}\\ u=g & \text { on } \partial \Omega,\end{cases}
$$

on simply-connected domains, where $K$ is a prescribed positive $C^{2}$ function. We shall prove the following uniqueness results for $\rho \leq 8 \pi$.

[^0]Theorem 1.1. Let $\Omega$ be an open, bounded and simply-connected domain, and let $K \in C^{2}(\Omega) \cap C(\bar{\Omega})$ be positive. Assume that $v_{i} \in C^{2}(\Omega) \cap C(\bar{\Omega}), i=1,2$, satisfy

$$
\begin{equation*}
\Delta v_{i}+K e^{v_{i}}=f_{i}(x) \tag{3}
\end{equation*}
$$

where $f_{2} \geq f_{1} \geq-\Delta \ln (K)$ in $\Omega$. If $v_{2}-v_{1}>c$ on $\Omega, v_{2}-v_{1}=c$ on $\partial \Omega$ for some $c \in \mathbb{R}$, and

$$
\int_{\Omega} K e^{v_{1}} d x=\int_{\Omega} K e^{v_{2}} d x=\rho
$$

then $\rho>8 \pi$.
The above theorem is equivalent to the next uniqueness result. Indeed Theorem 1.1 follows from Theorem 1.2 by letting $w_{i}=\ln K+v_{i}, \mathrm{i}=1,2$, and Theorem 1.2 follows from Theorem 1.1 by letting $K \equiv 1$.
Theorem 1.2. Let $\Omega$ be an open, bounded, and simply-connected domain. Assume that $w_{i} \in C^{2}(\Omega) \cap C(\bar{\Omega}), i=1,2$, satisfy

$$
\begin{equation*}
\Delta w_{i}+e^{w_{i}}=f_{i}(x) \tag{4}
\end{equation*}
$$

where $f_{2} \geq f_{1} \geq 0$ in $\Omega$. If $w_{2}-w_{1}>c$ on $\Omega, w_{2}-w_{1}=c$ on $\partial \Omega$ for some $c \in \mathbb{R}$, and

$$
\int_{\Omega} e^{w_{1}} d x=\int_{\Omega} e^{w_{2}} d x=\rho
$$

then $\rho>8 \pi$.

Corollary 1.3. Let $\Omega$ be an open, bounded and simply-connected domain, and let $K \in C^{2}(\Omega) \cap C(\bar{\Omega})$ be positive. Assume that $u_{i} \in C^{2}(\Omega) \cap C(\bar{\Omega}), i=1,2$ satisfy

$$
\begin{equation*}
\Delta u_{i}+\rho \frac{K e^{u_{i}}}{\int_{\Omega} K e^{u_{i}}}=f_{i}(x) \tag{5}
\end{equation*}
$$

where $f_{2} \geq f_{1} \geq-\Delta \ln (K)$ in $\Omega$, and $u_{2}-u_{1}>c$ on $\Omega$. If $u_{2}-u_{1}=c$ on $\partial \Omega$ for some $c \in \mathbb{R}$, then $\rho>8 \pi$.

We also present the following uniqueness results on multiply-connected domains.
Theorem 1.4. Let $\Omega$ be an open, bounded and multiply-connected domain, and let $K \in C^{2}(\Omega) \cap C(\bar{\Omega})$ be positive. Assume that $v_{i} \in C^{2}(\Omega) \cap C(\bar{\Omega}), i=1,2$, satisfy

$$
\begin{equation*}
\Delta v_{i}+K e^{v_{i}}=f_{i}(x) \tag{6}
\end{equation*}
$$

where $f_{2} \geq f_{1} \geq-\Delta \ln (K)$. If $v_{2}-v_{1}>c_{2}-c_{1}$ on $\Omega, v_{1}+\ln (K)=c_{1}$ and $v_{2}+\ln (K)=c_{2}$ on $\partial \Omega$ for some $c_{1}, c_{2} \in \mathbb{R}$, and

$$
\int_{\Omega} K e^{v_{1}} d x=\int_{\Omega} K e^{v_{2}} d x=\rho
$$

then $\rho>8 \pi$.
Corollary 1.5. Let $\Omega$ be an open, bounded and multiply-connected domain, and let $K \in C^{2}(\Omega) \cap C(\bar{\Omega})$ be positive. Assume that $u_{i} \in C^{2}(\Omega) \cap C(\bar{\Omega}), i=1,2$ satisfy

$$
\begin{equation*}
\Delta u_{i}+\rho \frac{K e^{u_{i}}}{\int_{\Omega} K e^{u_{i}}}=f_{i}(x) \tag{7}
\end{equation*}
$$

where $f_{2} \geq f_{1} \geq-\Delta \ln (K)$. If $u_{2}-u_{1}>c_{2}-c_{1}$ on $\Omega, u_{1}+\ln (K)=c_{1}$ and $u_{2}+\ln (K)=c_{2}$ on $\partial \Omega$ for some $c_{1}, c_{2} \in \mathbb{R}$, then $\rho>8 \pi$.

We should mention that Theorems 1.1 and 1.2 , and Corollary 1.3 are known in some special cases and also when the wight $K$ is singular (see [4]).

We shall use a new idea from [14] regarding the rearrangement of the solutions according to the standard metric on a sphere (projected to $\mathbb{R}^{2}$ ) and compare the total masses of the solutions. In particular, we shall show a reversed Bol's inequality in exterior domain for weak radial solutions (Proposition 3.1).

## 2. Preliminaries

Bol's isoperimetric inequality plays a crucial role in the proof of our main results. In this section we gather some results on Bol's inequality that will be used in subsequent sections. Let us first recall the classical Bol's isoperimetric inequality; see [2,4, 6, 8, 21, 23, and [7] for a detailed history of the Bol's inequality.

Theorem B. Let $\Omega \subseteq \mathbb{R}^{2}$ be a simply-connected and assume $u \in C^{2}(\Omega) \cap C(\bar{\Omega})$ satisfies

$$
\begin{equation*}
\Delta u+e^{u} \geq 0, \quad \int_{\Omega} e^{u} d x \leq 8 \pi \tag{8}
\end{equation*}
$$

Then for every $\omega \Subset \Omega$ of class $C^{1}$ the following inequality holds:

$$
\begin{equation*}
\left(\int_{\partial \omega} e^{\frac{u}{2}}\right)^{2} \geq \frac{1}{2}\left(\int_{\omega} e^{u}\right)\left(8 \pi-\int_{\omega} e^{u}\right) . \tag{9}
\end{equation*}
$$

Moreover the inequality in (19) is strict if $\Delta u+e^{u}>0$ somewhere in $\omega$ or $\omega$ is not simply-connected.

For $\lambda>0$ the function $U_{\lambda}$ defined by

$$
\begin{equation*}
U_{\lambda}:=-2 \ln \left(1+\frac{\lambda^{2}|y|^{2}}{8}\right)+2 \ln (\lambda) \tag{10}
\end{equation*}
$$

satisfies

$$
\Delta U_{\lambda}+e^{U_{\lambda}}=0
$$

and

$$
\left(\int_{\partial B_{r}} e^{\frac{U_{\lambda}}{2}}\right)^{2}=\frac{1}{2}\left(\int_{B_{r}} e^{U_{\lambda}}\right)\left(8 \pi-\int_{B_{r}} e^{U_{\lambda}}\right)
$$

for all $r>0$ and $\lambda>0$, where $B_{r}$ denotes the ball of radius $r$ centered at the origin in $\mathbb{R}^{2}$.

Note that Theorem B requires $\Omega$ to be simply-connected but $\omega$ can be multiplyconnected. Recently in 5 Theorem B is extended to the case where $\Omega \subset \mathbb{R}^{2}$ is multiply-connected and $u$ is constant on $\partial \Omega$.
Theorem C (Theorem 3 in [5). Let $\Omega$ be an open and bounded domain of class $C^{1}$ in $\mathbb{R}^{2}$ and assume $u \in C^{2}(\Omega) \cap C(\bar{\Omega})$ satisfies (8) and $u=c$ on $\partial \Omega$, for some constant $c \in \mathbb{R}$. Then (9) holds for every $\omega \Subset \Omega$. Moreover the inequality is strict if $\Delta u+e^{u}>0$ somewhere in $\omega$ or $\omega$ is not simply-connected.

Let $\Omega$ be an open, bounded, and multiply-connected domain of class $C^{1}$ in $\mathbb{R}^{2}$, and $\overline{\Omega^{*}}$ be the closure of the union of the bounded components of $\mathbb{R}^{2} \backslash \partial \Omega$ and $\Omega^{*}=\overline{\Omega^{*}} \backslash \partial \overline{\Omega^{*}}$. It is easy to see that $\Omega \subset \Omega^{*}$. Suppose $g \in C(\partial \Omega)$ satisfies

$$
\begin{equation*}
g=\left.G\right|_{\partial \Omega} \tag{11}
\end{equation*}
$$

where $G$ is Lipschitz continuous in $\overline{\Omega^{*}}, G$ is subharmonic in $\overline{\Omega^{*}}$ and harmonic in $\Omega$. The following more general result is also proved in [5].

Theorem D (Theorem 4 in [5]). Let $\Omega$ be an open, bounded, and multiply-connected domain of class $C^{1}$ in $\mathbb{R}^{2}$. Suppose $u \in C^{2}(\Omega) \cap C(\bar{\Omega})$ satisfies (8) with $u=g$ on $\partial \Omega$, and $g \in C(\partial \Omega)$ satisfies (11). Then (9) holds for every $\omega \Subset \Omega$. Moreover the inequality is strict if $\Delta u+e^{u}>0$ somewhere in $\omega$ or $\omega$ is not simply-connected.

Next we shall recall some facts about rearrangements according to the metric on $\mathbb{R}^{2}$ which is the stereographic projection of the standard metric on the unit sphere. Such rearrangments are discussed in detail in [14], but we also include it here for the sake of the readers. Let $\Omega \subset \mathbb{R}^{2}$ and $\lambda>0$, and suppose that $u \in C^{2}(\bar{\Omega})$ satisfies

$$
\Delta u+e^{u} \geq 0
$$

Then any function $\phi \in C^{2}(\bar{\Omega})$ which is constant on $\partial \Omega$ can be equimeasurably rearranged with respect to the measures $e^{u} d y$ and $e^{U_{\lambda}} d y$ (see [2], 4], [8, [21], [23]), where $U_{\lambda}$ is defined in (10). More precisely, for $t>\min _{y \in \bar{\Omega}} \phi$ define

$$
\Omega_{t}:=\{\phi>t\} \subset \subset \Omega,
$$

and let $\Omega_{t}^{*}$ be the ball centered at the origin in $\mathbb{R}^{2}$ such that

$$
\int_{\Omega_{t}^{*}} e^{U_{\lambda}} d y=\int_{\Omega_{t}} e^{u} d y:=a(t)
$$

Then $a(t)$ is a right-continuous function, and $\phi^{*}: \Omega^{*} \rightarrow \mathbb{R}$ defined by $\phi^{*}(y):=$ $\sup \left\{t \in \mathbb{R}: y \in \Omega_{t}^{*}\right\}$ provides an equimeasurable rearrangement of $\phi$ with respect to the measure $e^{u} d y$ and $e^{U_{\lambda}} d y$, i.e.

$$
\begin{equation*}
\int_{\left\{\phi^{*}>t\right\}} e^{U_{\lambda}} d y=\int_{\{\phi>t\}} e^{u} d y, \quad \forall t>\min _{y \in \bar{\Omega}} \phi . \tag{12}
\end{equation*}
$$

We shall need the following lemma.
Proposition 2.1. Let $u, \varphi \in C^{1}(\bar{\Omega})$ and assume that $\phi$ is constant on $\partial \Omega$. Let $\phi^{*}(r)$ be the equimeasurable rearrangement of $\phi$ with respect to the measure $e^{u} d y$ and $e^{U_{\lambda}} d y$. Then $\phi^{*}$ is Lipschitz continuous on $(\epsilon, R-\epsilon)$, for every $\epsilon>0$, where $R$ is the radius of $\Omega^{*}$.
Proof. First note that the function $\phi^{*}$ is decreasing and the set

$$
\mathcal{T}:=\left\{t \geq \min _{\bar{\Omega}} \phi: \quad\left(\phi^{*}\right)^{-1}(t) \text { is not a singleton }\right\}
$$

has Lebesgue measure zero. Indeed $\left(\phi^{*}\right)^{-1}(t)$ is a connected closed interval for all $t \in \mathcal{T}$. Let $0<r_{1}<r_{2}<R$ and

$$
a(t)=\int_{\left\{\phi^{*}>t\right\}} e^{U_{\lambda}} d y=\int_{\{\phi>t\}} e^{u} d y, \quad \forall t>\min _{y \in \bar{\Omega}} \phi .
$$

For $\phi^{*}\left(r_{1}\right), \phi^{*}\left(r_{2}\right) \notin \mathcal{T}$, we have

$$
\begin{aligned}
a\left(\phi^{*}\left(r_{2}\right)\right)-a\left(\phi^{*}\left(r_{1}\right)\right) & =\int_{\left\{\phi^{*}(|y|)>\phi^{*}\left(r_{2}\right)\right\}} e^{U_{\lambda}} d y-\int_{\left\{\phi^{*}(|y|)>\phi^{*}\left(r_{2}\right)\right\}} e^{U_{\lambda}} d y \\
& =\int_{\left\{\phi(y)>\phi^{*}\left(r_{2}\right)\right\}} e^{u} d y-\int_{\left\{\phi(y)>\phi^{*}\left(r_{2}\right)\right\}} e^{u} d y \\
& =\int_{\left\{\phi^{*}\left(r_{2}\right) \leq \phi(y) \leq \phi^{*}\left(r_{1}\right)\right\}} e^{u} d y \\
& =\int_{\left\{\phi^{*}\left(r_{2}\right) \leq \phi^{*}(|y|) \leq \phi^{*}\left(r_{1}\right)\right\}} e^{U_{\lambda}} d y .
\end{aligned}
$$

Now let $m:=\min _{\bar{\Omega}} e^{u(y)}, M_{1}:=\max _{\overline{\Omega^{*}}} e^{U_{\lambda}(y)}$, and $M_{2}:=\max _{\bar{\Omega}}|\nabla \phi|$. Then it follows from the above equality that

$$
\begin{aligned}
a\left(\phi^{*}\left(r_{2}\right)\right)-a\left(\phi^{*}\left(r_{1}\right)\right) & \leq M_{1} \mu\left(\left\{\phi^{*}\left(r_{2}\right) \leq \phi^{*}(|y|) \leq \phi^{*}\left(r_{1}\right)\right\}\right) \\
& =M_{1} \mu\left(r_{1} \leq|y| \leq r_{2}\right)=M_{1} \pi\left(r_{2}^{2}-r_{1}^{2}\right) \\
& \leq 2 \pi R M_{1}\left(r_{2}-r_{1}\right)
\end{aligned}
$$

On the other hand,

$$
\begin{aligned}
a\left(\phi^{*}\left(r_{2}\right)\right)-a\left(\phi^{*}\left(r_{1}\right)\right) & \geq m \mu\left(\left\{\phi^{*}\left(r_{2}\right) \leq \phi(y) \leq \phi^{*}\left(r_{1}\right)\right\}\right) \\
& \geq \frac{m}{M_{2}} \int_{\left\{\phi^{*}\left(r_{2}\right) \leq \phi(y) \leq \phi^{*}\left(r_{1}\right)\right\}}|\nabla \phi| d y \\
& \geq \frac{m}{M_{2}} \int_{\phi^{*}\left(r_{2}\right)}^{\phi^{*}\left(r_{1}\right)} \int_{\left\{\phi^{-1}(t)\right\}} d s d t \\
& \geq \frac{m}{M_{2}}\left(\phi^{*}\left(r_{1}\right)-\phi^{*}\left(r_{2}\right)\right) K\left(r_{1}, r_{2}\right),
\end{aligned}
$$

where

$$
K\left(r_{1}, r_{2}\right)=\min _{\left\{\phi^{*}\left(r_{2}\right) \leq t \leq \phi^{*}\left(r_{1}\right)\right\}} \mathcal{H}^{n-1}\left(\phi^{-1}(t)\right)>0, \quad 0<r_{1}<r_{2}<R
$$

Since $\left\{\phi^{-1}(t)\right\}=\partial\{x: \phi(x)>t\}$, it follows from the isoperimetric inequality that if $\phi^{*}\left(r_{1}\right)<\max _{y \in \bar{\Omega}} \phi-\delta$ and $\phi^{*}\left(r_{2}\right)>\min _{y \in \bar{\Omega}} \phi+\delta$, for some $\alpha \in(0,1)$, then

$$
K\left(r_{1}, r_{2}\right)>C>0, \quad \forall r_{2} \quad \text { with } \quad r_{1}<r_{2}<R,
$$

for some $C>0$ independent of $\phi$. Hence we have

$$
\begin{equation*}
0 \leq \frac{a\left(\phi^{*}\left(r_{2}\right)\right)-a\left(\phi^{*}\left(r_{1}\right)\right)}{r_{2}-r_{1}} \leq \frac{2 \pi R M_{1} M_{2}}{m K\left(r_{1}, r_{2}\right)} \leq \frac{2 \pi R M_{1} M_{2}}{m C} \tag{13}
\end{equation*}
$$

By approximation the above also holds for $\epsilon<r_{1}<r_{2}<R-\epsilon$. Thus $\phi^{*}$ is Lipschitz continuous on $(\epsilon, R-\epsilon)$ for every $\epsilon>0$.

Now let

$$
\begin{aligned}
j(t) & :=\int_{\{\phi>t\}}|\nabla \phi|^{2} d y, \quad j^{*}(t):=\int_{\{\phi>t\}}\left|\nabla \phi^{*}\right|^{2} d y, \quad \forall t>\min _{y \in \bar{\Omega}} \phi ; \\
J(t) & :=\int_{\{\phi>t\}}|\nabla \phi| d y, \quad J^{*}(t):=\int_{\left\{\phi^{*}>t\right\}}\left|\nabla \phi^{*}\right| d y, \quad \forall t>\min _{y \in \bar{\Omega}} \phi .
\end{aligned}
$$

It is easy to see that both $j(t)$ and $J(t)$ are absolutely continuous and decreasing in $t>\min _{y \in \bar{\Omega}} \phi$. If $\phi \equiv C$ on $\partial \Omega$, it can be shown that

$$
\begin{equation*}
\int_{\{\phi=t\}}|\nabla \phi| d s \geq \int_{\left\{\phi^{*}=t\right\}}\left|\nabla \phi^{*}\right| d s, \quad \text { for } \text { a.e. } t>\min _{y \in \bar{\Omega}} \phi \tag{14}
\end{equation*}
$$

Indeed it follows from Cauchy-Schwarz and Bol's inequalities that

$$
\begin{aligned}
\int_{\{\phi=t\}}|\nabla \phi| d s & \geq\left(\int_{\{\phi=t\}} e^{\frac{u}{2}}\right)^{2}\left(\int_{\{\phi=t\}} \frac{e^{u}}{|\nabla \phi|}\right)^{-1} \\
& =\left(\int_{\{\phi=t\}} e^{\frac{u}{2}}\right)^{2}\left(-\frac{d}{d t} \int_{\Omega_{t}} e^{u}\right)^{-1} \\
& \geq \frac{1}{2}\left(\int_{\Omega_{t}} e^{u}\right)\left(8 \pi-\int_{\Omega_{t}} e^{u}\right)\left(-\frac{d}{d t} \int_{\Omega_{t}} e^{u}\right)^{-1} \\
& =\frac{1}{2}\left(\int_{\Omega_{t}^{*}} e^{U_{\lambda}}\right)\left(8 \pi-\int_{\Omega_{t}^{*}} e^{U_{\lambda}}\right)\left(-\frac{d}{d t} \int_{\Omega_{t}^{*}} e^{U_{\lambda}}\right)^{-1} \\
& =\int_{\left\{\phi^{*}=t\right\}}\left|\nabla \phi^{*}\right| d s, \text { for } a . e . t>\min _{y \in \bar{\Omega}} \phi .
\end{aligned}
$$

It also follows that $j^{*}(t), J^{*}(t)$ are absolutely continuous and decreasing in $t>$ $\min _{y \in \bar{\Omega}} \phi$, since both functions are right-continuous by definition and

$$
\begin{aligned}
& 0 \leq j^{*}(t-0)-j^{*}(t) \leq j(t-0)-j(t)=\int_{\{\phi=t\}}|\nabla \phi|^{2} d y=0, \quad t>\min _{y \in \bar{\Omega}} \phi \\
& 0 \leq J^{*}(t-0)-J^{*}(t) \leq J(t-0)-J(t)=\int_{\{\phi=t\}}|\nabla \phi| d y=0, \quad t>\min _{y \in \bar{\Omega}} \phi
\end{aligned}
$$

Therefore we have the following proposition.
Proposition 2.2. Let $u \in C^{2}(\bar{\Omega})$ satisfy

$$
\Delta u+e^{u} \geq 0 \quad \text { in } \Omega,
$$

and let $U_{\lambda}$ be given by (10). Suppose $\phi \in C^{1}(\bar{\Omega})$ and $\phi \equiv C$ on $\partial \Omega$. Define the equimeasurable symmetric rearrangement $\phi^{*}$ of $\phi$, with respect to the measures $e^{u} d y$ and $e^{U_{\lambda}} d y$, by (12). Then $\phi^{*}$ is Lipschitz continuous on $(\epsilon, R-\epsilon)$ for every $\epsilon>0$, and $j^{*}(t), J^{*}(t)$ are absolutely continuous and decreasing in $t>\min _{y \in \bar{\Omega}} \phi$ and (14) holds.

## 3. Bol's type inequalities

We first prove the following lemma.
Lemma 3.1. Let $\psi \in C\left(\mathbb{R}^{2} \backslash B_{R}\right)$ be a decreasing radial function and

$$
\int_{\left(\mathbb{R}^{2} \backslash B_{R}\right)} e^{\psi} d x<\infty
$$

for some $R>0$. Then

$$
\lim _{s \rightarrow-\infty} e^{s} \int_{\{\psi>s\}} d x=0
$$

Proof. Since $\psi$ is decreasing,

$$
\frac{3 \pi}{4} r^{2} e^{\psi(r)} \leq \int_{\left(B_{r} \backslash B_{r / 2}\right)} e^{\psi} d x
$$

for $r>2 R$. Letting $r \rightarrow \infty$ we obtain,

$$
\lim _{r \rightarrow \infty} r^{2} e^{\psi(r)}=0 .
$$

Define

$$
r(s):=\sup \{r \geq R: \psi(r)>s\}, \quad s \in \mathbb{R}
$$

Then $r(s)$ is well defined for $s<\psi(R)$ and $\lim _{s \rightarrow-\infty} r(s)=\infty$. Since

$$
e^{s} \int_{\{\psi>s\}} d x \leq \pi\left(r(s)^{2}-R^{2}\right) e^{\psi(r(s))}
$$

we obtain

$$
\lim _{s \rightarrow-\infty} e^{s} \int_{\{\psi>s\}} d x=0
$$

The proof is complete.
For the proof of our main results, we shall need the following reversed Bol's inequality.
Proposition 3.1. Let $B_{R}$ be the ball of radius $R$ in $\mathbb{R}^{2} \psi \in C^{0,1}\left(\mathbb{R}^{2} \backslash B_{R}\right)$ be a strictly decreasing radial function satisfying
(15) $\int_{\partial B_{r}}|\nabla \psi| d s \leq 8 \pi-\int_{\mathbb{R}^{2} \backslash B_{r}} e^{\psi}$ for a.e. $r \in(R, \infty)$, and $\int_{\mathbb{R}^{2} \backslash B_{R}} e^{\psi}<8 \pi$.

Then the following inequality holds:

$$
\begin{equation*}
\left(\int_{\partial B_{R}} e^{\frac{\psi}{2}}\right)^{2} \leq \frac{1}{2}\left(\int_{\mathbb{R}^{2} \backslash B_{R}} e^{\psi}\right)\left(8 \pi-\int_{\mathbb{R}^{2} \backslash B_{R}} e^{\psi}\right) \tag{16}
\end{equation*}
$$

Moreover if $\int_{\partial B_{r}}|\nabla \psi| d s \not \equiv 8 \pi-\int_{\mathbb{R}^{2} \backslash B_{r}} e^{\psi}$ on $(R, \infty)$, then the inequality in (16) is strict.

Proof. Let $\beta:=\psi(R)$ and define

$$
k(s)=8 \pi-\int_{\{\psi<s\}} e^{\psi} d x, \text { and } \mu(s)=\int_{\{\psi>s\}} d x+\pi R^{2},
$$

for $s<\beta$. Then

$$
-k^{\prime}(s)=\int_{\{\psi=s\}} \frac{e^{\psi}}{|\nabla \psi|}=-e^{s} \mu^{\prime}(s) .
$$

Hence

$$
\begin{align*}
-k(s) k^{\prime}(s) & \geq \int_{\{\psi=s\}}|\nabla \psi| \cdot \int_{\{\psi=s\}} \frac{e^{\psi}}{|\nabla \psi|}  \tag{17}\\
& =\left(\int_{\{\psi=s\}} e^{\psi / 2}\right)^{2}=e^{s}\left(\int_{\{\psi=s\}} d s\right)^{2} \\
& =e^{s} \cdot 4 \pi\left(\int_{\{\psi>s\}} d x+\pi R^{2}\right)=4 \pi e^{s} \mu(s),
\end{align*}
$$

for a.e. $s<\beta$. Therefore

$$
\frac{d}{d s}\left[e^{s} \mu(s)-k(s)+\frac{1}{8 \pi} k^{2}(s)\right]=\mu(s)+\frac{1}{4 \pi} k^{\prime}(s) k(s) \leq 0
$$

for a.e. $s<\beta$. Integrating on $(-\infty, \beta)$ and using Lemma 3.1 we get

$$
\begin{equation*}
\left[e^{s} \mu(s)-k(s)+\frac{1}{8 \pi} k^{2}(s)\right]_{-\infty}^{\beta}=e^{\beta} \mu(\beta)-k(\beta)+\frac{1}{8 \pi} k^{2}(\beta) \leq 0 . \tag{18}
\end{equation*}
$$

Now notice that

$$
k(\beta)=8 \pi-\int_{\mathbb{R}^{2} \backslash B_{R}} e^{\psi} d x
$$

and

$$
e^{\beta} \mu(\beta)=e^{\beta} \int_{B_{R}} d x=\frac{1}{4 \pi} e^{\beta}\left(\int_{\partial B_{R}} d s\right)^{2}=\frac{1}{4 \pi}\left(\int_{\partial B_{R}} e^{\frac{\psi}{2}} d s\right)^{2} .
$$

Thus (16) follows from the inequality (18). Finally if $\int_{\partial B_{r}}|\nabla \psi| d s \not \equiv 8 \pi-\int_{\mathbb{R}^{2} \backslash B_{r}} e^{\psi}$ on $(R, \infty)$, then the inequality (17) will be strict, and consequently (16) will also be strict.

Similarly one can prove the following proposition (see Proposition 2.2 in [14).
Proposition 3.2. Let $B_{R}$ be the ball of radius $R$ in $\mathbb{R}^{2}$ and $u \in C^{0,1}\left(\overline{B_{R}}\right)$ be a strictly decreasing radial function satisfying

$$
\begin{equation*}
\int_{\partial B_{r}}|\nabla u| d s \leq \int_{B_{r}} e^{u} \quad \text { for a.e. } \quad r \in(0, R), \quad \text { and } \quad \int_{B_{R}} e^{u}<8 \pi . \tag{19}
\end{equation*}
$$

Then the following inequality holds:

$$
\begin{equation*}
\left(\int_{\partial B_{R}} e^{\frac{u}{2}}\right)^{2} \geq \frac{1}{2}\left(\int_{B_{R}} e^{u}\right)\left(8 \pi-\int_{B_{R}} e^{u}\right) \tag{20}
\end{equation*}
$$

Moreover if $\int_{\partial B_{r}}|\nabla \psi| d s \not \equiv \int_{B_{r}} e^{\psi}$ on $(0, R)$, then the inequality in (20) is strict.

## 4. Proof of the main results

Lemma 4.1. Let $R>0$ and assume that $\psi \in C^{0,1}\left(\mathbb{R}^{2} \backslash B_{R}\right)$ is a strictly decreasing radial function that satisfies

$$
\begin{equation*}
\int_{\partial B_{r}}|\nabla \psi| \leq 8 \pi-\int_{\mathbb{R}^{2} \backslash B_{r}} e^{\psi} \tag{21}
\end{equation*}
$$

for a.e. $r \in(R, \infty)$ and $\psi=U_{\lambda_{1}}=U_{\lambda_{2}}$ on $\partial B_{R}$ for some $\lambda_{2}>\lambda_{1}$. Then

$$
\begin{equation*}
\int_{\mathbb{R}^{2} \backslash B_{R}} e^{{U \lambda_{2}}_{2}} \leq \int_{\mathbb{R}^{2} \backslash B_{R}} e^{\psi} \leq \int_{\mathbb{R}^{2} \backslash B_{R}} e^{U_{\lambda_{1}}} \tag{22}
\end{equation*}
$$

Moreover if $\int_{\partial B_{r}}|\nabla \psi| \not \equiv \int_{\mathbb{R}^{2} \backslash B_{r}} e^{\psi}$ on $r \in(R, \infty)$, then the inequalities in (22) are also strict.
Proof. Let $m_{1}:=\int_{\mathbb{R}^{2} \backslash B_{R}} e^{U_{\lambda_{1}}}, m_{2}:=\int_{\mathbb{R}^{2} \backslash B_{R}} e^{U_{\lambda_{2}}}$, and $m:=\int_{\mathbb{R}^{2} \backslash B_{R}} e^{\psi}$. Define

$$
\beta:=\left(\int_{\partial B_{R}} e^{\frac{\psi}{2}}\right)^{2}=\left(\int_{\partial B_{R}} e^{\frac{U_{\lambda_{1}}}{2}}\right)^{2}=\left(\int_{\partial B_{R}} e^{\frac{U \lambda_{2}}{2}}\right)^{2}
$$

It follows from Proposition 3.1 that

$$
\beta \leq \frac{1}{2} m(8 \pi-m) .
$$

On the other hand

$$
\beta=\frac{1}{2} m_{1}\left(8 \pi-m_{1}\right)=\frac{1}{2} m_{2}\left(8 \pi-m_{2}\right),
$$

i.e. $m_{1}$ and $m_{2}$ are roots of the quadratic equation

$$
x^{2}-8 \pi x+2 \beta=0
$$

Since $m$ satisfies

$$
m^{2}-8 \pi m+2 \beta \leq 0
$$

we have

$$
m_{2} \leq m \leq m_{1}
$$

Similarly the following lemma holds (see Lemma 3.3 in [14]).
Lemma 4.2. Assume that $\psi \in C^{0,1}\left(\overline{B_{R}}\right)$ is a strictly decreasing, radial, Lipschitz function, and satisfies

$$
\begin{equation*}
\int_{\partial B_{r}}|\nabla \psi| \leq \int_{B_{r}} e^{\psi} \tag{23}
\end{equation*}
$$

a.e. $r \in(0, R)$ and $\psi=U_{\lambda_{1}}=U_{\lambda_{2}}$ for some $\lambda_{2}>\lambda_{1}$ on $\partial B_{R}$, and $R>0$. Then there holds

$$
\begin{equation*}
\text { either } \int_{B_{R}} e^{\psi} \leq \int_{B_{R}} e^{U_{\lambda_{1}}} \quad \text { or } \quad \int_{B_{R}} e^{\psi} \geq \int_{B_{R}} e^{U_{\lambda_{2}}} \tag{24}
\end{equation*}
$$

Moreover if the inequality in (23) is strict in a set with positive measure in $(0, R)$, then the inequalities in (24) are also strict.

We shall also need the following lemma.
Lemma 4.3. Assume that $\psi \in C^{0,1}\left(\overline{B_{R}}\right)$ is a strictly decreasing radial function satisfying (23) for a.e. $r \in(0, R)$. If

$$
\rho=\int_{B_{R}} e^{\psi} d x=\int_{B_{R}} e^{U_{\lambda}}<8 \pi
$$

then $U_{\lambda}(R) \leq \psi(R)$.
Proof. By Proposition 3.2 we have

$$
\begin{aligned}
\left(\int_{\partial B_{R}} e^{\frac{U_{\lambda}}{2}}\right)^{2} & =\frac{1}{2}\left(\int_{B_{R}} e^{U_{\lambda}}\right)\left(8 \pi-\int_{B_{R}} e^{U_{\lambda}}\right) \\
& =\frac{1}{2}\left(\int_{B_{R}} e^{\psi}\right)\left(8 \pi-\int_{B_{R}} e^{\psi}\right) \\
& \leq\left(\int_{\partial B_{R}} e^{\frac{\psi}{2}}\right)^{2}
\end{aligned}
$$

and hence $U_{\lambda}(R) \leq \psi(R)$.
Now we are ready to prove the main result of this paper, Theorem 1.2.
Proof of Theorem 1.2. First we prove that $\rho \geq 8 \pi$. Suppose $w_{1}$ and $w_{2}$ satisfy the assumptions of Theorem 1.2, Then

$$
\Delta\left(w_{2}-w_{1}\right)+e^{w_{2}}-e^{w_{1}}=f_{2}-f_{1} \geq 0
$$

Now choose $\lambda>0$ and $R \in(0, \infty)$ such that

$$
\begin{equation*}
\int_{\Omega} e^{w_{1}}=\int_{B_{R}} e^{U_{\lambda}} \tag{25}
\end{equation*}
$$

and let $\phi$ be the symmetrization of $w_{2}-w_{1}$ with respect to the measures $e^{w_{1}} d y$ and $e^{U_{\lambda}} d y$. Then it follows from Proposition 2.2 and Fubini's theorem that

$$
\begin{aligned}
\int_{\{\phi=t\}}|\nabla \phi| & \leq \int_{\left\{w_{2}-w_{1}=t\right\}}\left|\nabla\left(w_{2}-w_{1}\right)\right| \\
& \leq \int_{\Omega_{t}} e^{w_{2}}-e^{w_{1}} d x \\
& =\int_{\{\phi>t\}} e^{U_{\lambda}+\phi}-\int_{\{\phi>t\}} e^{U_{\lambda}} \\
& =\int_{\{\phi>t\}} e^{U_{\lambda}+\phi}-\int_{\{\phi=t\}}\left|\nabla U_{\lambda}\right|,
\end{aligned}
$$

for a.e. $t>\inf _{\Omega}\left(w_{2}-w_{1}\right)$. Hence

$$
\begin{equation*}
\int_{\{\phi=t\}}\left|\nabla\left(U_{\lambda}+\phi\right)\right| \leq \int_{\{\phi>t\}} e^{\left(U_{\lambda}+\phi\right)} d \tag{26}
\end{equation*}
$$

for all $t>\inf _{\Omega}\left(w_{2}-w_{1}\right)$. Since $\phi$ is decreasing in $r, \psi:=U_{\lambda}+\phi$ is a strictly decreasing function, and

$$
\begin{equation*}
\int_{\partial B_{r}}|\nabla \psi| \leq \int_{B_{r}} e^{\psi} d y, \quad \text { a.e. } \quad r \in(0, R), \tag{27}
\end{equation*}
$$

by Proposition 2.2 and the above inequality we see that $\psi \in W^{1, \infty}\left(B_{R}\right)$ and thus by Morrey's inequality $\psi \in C^{0,1}\left(B_{R}\right)$.

Since $w_{1} \not \equiv w_{2}$ and $\int_{\Omega} e^{w_{1}}=\int_{\Omega} e^{w_{2}}$, then $w_{2}<w_{1}$ on a subset of $\Omega$ with positive measure. Hence $\phi(R)<0$ and consequently $\psi(R)=U_{\lambda}(R)+\phi(R)<U_{\lambda}(R)$. This is a contradiction in view of Lemma 4.3 and therefore we must have $\rho \geq 8 \pi$.

Next we prove that $\rho>8 \pi$. Suppose $\rho=8 \pi$ and let $\lambda_{1}>0$. With an argument similar to the one above we may show that there exists $\psi=U_{\lambda_{1}}+\phi \in C^{0,1}\left(\mathbb{R}^{2}\right)$ such that

$$
\int_{\Omega} e^{w_{1}} d x=\int_{\mathbb{R}^{2}} e^{U_{\lambda_{1}}}=8 \pi=\int_{\Omega} e^{w_{2}} d x=\int_{\mathbb{R}^{2}} e^{\psi} d x
$$

and

$$
\begin{equation*}
\int_{\partial B_{r}}|\nabla \psi| \leq \int_{B_{r}} e^{\psi} d x \tag{28}
\end{equation*}
$$

for a.e. $r \in(0, \infty)$. Since $\int_{\mathbb{R}^{2}} e^{\psi}=\int_{\mathbb{R}^{2}} e^{U_{\lambda_{1}}}$, there exists $r_{0} \in(0, \infty)$ such that $\psi\left(r_{0}\right)=U_{\lambda_{1}}\left(r_{0}\right)$. There exists a positive constant $\lambda_{2} \neq \lambda_{1}$ such that $U_{\lambda_{2}}\left(r_{0}\right)=$ $U_{\lambda_{1}}\left(r_{0}\right)=\psi\left(r_{0}\right)$. Since $\psi>U_{\lambda_{1}}$ in $B_{r_{0}}$, it follows from Proposition4.2 that $\lambda_{1}<\lambda_{2}$ and

$$
\int_{B_{r_{0}}} e^{\psi} \geq \int_{B_{r_{0}}} e^{U_{\lambda_{2}}}
$$

On the other hand $\psi<U_{\lambda_{1}}$ in $\mathbb{R}^{2} \backslash B_{r_{0}}$ and consequently it follows from Proposition 4.1 that

$$
\int_{\mathbb{R}^{2} \backslash B_{r_{0}}} e^{\psi} \geq \int_{\mathbb{R}^{2} \backslash B_{r_{0}}} e^{{U \lambda_{2}}^{2}}
$$

Hence

$$
\begin{equation*}
8 \pi=\rho=\int_{\mathbb{R}^{2}} e^{\psi} \geq \int_{\mathbb{R}^{2}} e^{U_{\lambda_{2}}}=8 \pi . \tag{29}
\end{equation*}
$$

Note that if $f_{1} \not \equiv 0$ or $f_{2} \not \equiv f_{1}$, then the inequality in (29) will be strict, which is a contradiciton. Suppose $f_{1} \equiv f_{2} \equiv 0$. We may assume without loss of generality that $c=w_{2}-w_{1} \geq 0$ on $\partial \Omega$, since otherwise we can switch $w_{1}$ and $w_{2}$. By (29) we conclude that the equality in (28) holds for a.e. $r \in(0, \infty)$ and $\psi=U_{\lambda_{2}}$. It also yields that the equality in (14) must be true for $\phi=w_{2}-w_{1}$ and $t \geq \inf _{\Omega} \phi$. By the proof of Proposition (2.2, we also know that Bol's inequality (9) on $\omega=\{\phi>t\}$ must be equality, and therefore $\{\phi>t\}$ must be simply-connected for $t \geq \inf _{\Omega} \phi$ by Theorem D. This is a contradiction since $\{\phi>t\}$ is not simply-connected when $\inf _{\Omega} \phi<t<0$. The contradiction implies $\rho>8 \pi$.

Proof of Theorem 1.4. The proof follows from Theorem D and the same argument used in the proof of Theorem 1.2,

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