ON HOMOGENIZATION FOR A PERIODIC ELLIPTIC OPERATOR IN A STRIP

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Dedicated to my dear teacher Mikhail Shlemovich Birman with love and gratitude

ABSTRACT. In a strip $\Pi = \mathbb{R} \times (0, a)$, the operator

$$A_\varepsilon = D_1 g_1(x_1/\varepsilon, x_2) D_1 + D_2 g_2(x_1/\varepsilon, x_2) D_2$$

is considered, where g_1, g_2 are periodic with respect to the first variable. Periodic boundary conditions are put on the boundary of the strip. The behavior of the operator A_ε in the limit $\varepsilon \rightarrow 0$ is studied. It is proved that, with respect to the operator norm in $L_2(\Pi)$, the resolvent $(A_\varepsilon + I)^{-1}$ tends to the resolvent of the effective operator A^0 . A sharp order estimate for the norm of the difference of the resolvents is obtained. The operator A^0 is of the same type, but its coefficients depend only on x_2 .

§0. INTRODUCTION

0.1. The problem of homogenization for periodic operators of mathematical physics is of significant interest both from the general-theoretic viewpoint and for applications. An extensive body of literature is devoted to homogenization problems. At the first place, the books [1]–[3] should be mentioned. A typical homogenization problem looks like this. Consider the family of operators

$$(0.1) \quad \mathcal{A}_\varepsilon = -\operatorname{div} g(\mathbf{x}/\varepsilon) \operatorname{grad}, \quad \mathbf{x} \in \mathbb{R}^d, \quad \varepsilon > 0,$$

where $g(\mathbf{x})$ is a positive definite bounded matrix-valued function periodic with respect to some lattice in $\mathbb{R}^d$. The problem is to find a *constant* positive matrix g^0 such that the resolvent $(\mathcal{A}_\varepsilon + I)^{-1}$ converges (in some sense) as $\varepsilon \rightarrow 0$ to the resolvent of the operator $\mathcal{A}^0 = -\operatorname{div} g^0 \operatorname{grad}$. The operator $\mathcal{A}^0$ is called the *effective* or homogenized operator.

A possible approach to homogenization of periodic differential operators in $\mathbb{R}^d$ consists in application of the Floquet theory. In this connection, we mention the papers [4]–[8]. In [7, 8], it was proposed to treat the homogenization procedure as a threshold effect near the lower edge of the spectrum of a periodic differential operator. Namely, let Ω be the cell of the lattice of periods, and let $\mathcal{A}(\mathbf{k})$ be the family of operators acting in $L_2(\Omega)$ that appear in the direct integral decomposition for the operator $\mathcal{A} = -\operatorname{div} g(\mathbf{x}) \operatorname{grad}$. The parameter $\mathbf{k} \in \mathbb{R}^d$ is called the *quasimomentum*. The operator $\mathcal{A}(\mathbf{k})$ is defined by the expression $(\mathbf{D} + \mathbf{k})^* g(\mathbf{x})(\mathbf{D} + \mathbf{k})$ and periodic boundary conditions. The homogenization problem for $\mathcal{A}$ reduces to the study of a (two-parameter) family of operators $(\mathcal{A}(\mathbf{k}) + \varepsilon^2 I)^{-1}$ acting in $L_2(\Omega)$. Here, an important role is played by distinguishing the *spectral germ* of the family $\mathcal{A}(\mathbf{k})$ at $\mathbf{k} = 0$. In terms of the spectral germ, we can find a good approximation for the resolvent $(\mathcal{A}(\mathbf{k}) + \varepsilon^2 I)^{-1}$. Since the effective operator $\mathcal{A}^0$ has the same spectral germ, this

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leads to approximation of the resolvent $(\mathcal{A}(\mathbf{k}) + \varepsilon^2 I)^{-1}$ by the resolvent $(\mathcal{A}^0(\mathbf{k}) + \varepsilon^2 I)^{-1}$ of the effective operator. As a result, for the initial operator (0.1) acting in $L_2(\mathbb{R}^d)$, the facts established in [7, 8] give convergence of the resolvent in the operator norm with the following sharp order estimate:

$$\|(\mathcal{A}_\varepsilon + I)^{-1} - (\mathcal{A}^0 + I)^{-1}\|_{L_2(\mathbb{R}^d) \rightarrow L_2(\mathbb{R}^d)} \leq C\varepsilon.$$

0.2. Along with problems in the entire space $\mathbb{R}^d$, periodic problems can be considered in domains of cylinder type, layer type, etc., where the coefficients of the operator are periodic only in part of variables. In such problems, the effective operator depends on the “nonperiodic” variables, while homogenization is related to the “periodic” variables.

In the present paper, for a simplest model, we try to apply the approach of [7, 8] to such problems. We study the two-dimensional operator $A = -\operatorname{div} g(x_1, x_2) \operatorname{grad}$ in the strip $\Pi = \{(x_1, x_2) : x_1 \in \mathbb{R}, x_2 \in (0, a)\}$. The matrix $g(\mathbf{x})$ is periodic along this strip (with respect to x_1) with period 1. On the boundary $\partial\Pi$, we put periodic boundary conditions. The problem is to study the behavior of the operator

$$(0.2) \quad A_\varepsilon = -\operatorname{div} g(x_1/\varepsilon, x_2) \operatorname{grad}$$

as $\varepsilon \rightarrow 0$. It is assumed that the matrix $g(\mathbf{x})$ is diagonal: $g(\mathbf{x}) = \operatorname{diag}\{g_1(\mathbf{x}), g_2(\mathbf{x})\}$. It turns out that the effective operator A^0 is of the form $A^0 = -\operatorname{div} g^0(x_2) \operatorname{grad}$, where $g^0(x_2) = \operatorname{diag}\{g_1^0(x_2), g_2^0(x_2)\}$. Using the standard rules of homogenization theory, it is not difficult to find (formally) the expression for $g^0(x_2)$.

Our *goal in the present paper* is to obtain an estimate of sharp order for the difference of the resolvents of A_ε and A^0 in the operator norm in $L_2(\Pi)$. The difficulties we encounter are related to the fact that the method of [7, 8] does not automatically yield the result.

In this paper, for the simple model problem just described, we show how the estimate

$$\|(A_\varepsilon + I)^{-1} - (A^0 + I)^{-1}\|_{L_2(\Pi) \rightarrow L_2(\Pi)} \leq C\varepsilon$$

can still be obtained. We employ the results of [7, 8], but in combination with other methods. This requires some smoothness assumptions on the matrix-valued function $g(\mathbf{x})$ (see condition (1.4) below).

Now for the cell Ω we take $(0, 1) \times (0, a)$, and the parameter k is one-dimensional. After reducing the problem to the cell, we need to study a (two-parameter) family

$$A(k, \varepsilon) = A_1(k) + \varepsilon^2 A_2$$

of operators acting in $L_2(\Omega)$. Here $A_1(k) = (D_1 + k)g_1(\mathbf{x})(D_1 + k)$ and $A_2 = D_2g_2(\mathbf{x})D_2$; the boundary conditions are periodic in both variables. The method of [7, 8] can be applied to the operator $A_1(k)$ “layerwise”. This allows us to obtain a good approximation for $A_1(k)$ for small $|k|$ in terms of the corresponding spectral germ. The main technical difficulties are related to adding the *unbounded* “perturbation” $\varepsilon^2 A_2$ and approximating the resolvent $(A(k, \varepsilon) + \varepsilon^2 I)^{-1}$.

0.3. In §1, we state the problem and the main result. In §2, the problem is reduced to the study of the resolvent $(A(k, \varepsilon) + \varepsilon^2 I)^{-1}$. In §3, the operator pencil $A_1(k)$ is treated in detail. In §4, approximation for the family $(A(k, \varepsilon) + \varepsilon^2 I)^{-1}$ is constructed in terms of the spectral germ. §§5 and 6 are devoted to auxiliary technical estimates. The proof of the main theorem is completed in §7. In §8, possible generalizations are discussed.

0.4. Notation. Let $\mathfrak{H}$ and $\mathfrak{G}$ be separable Hilbert spaces. The symbols $(\cdot, \cdot)_{\mathfrak{H}}$ and $\|\cdot\|_{\mathfrak{H}}$ stand for the inner product and the norm in $\mathfrak{H}$; the symbol $\|\cdot\|_{\mathfrak{H} \rightarrow \mathfrak{G}}$ denotes the norm of a bounded operator from $\mathfrak{H}$ to $\mathfrak{G}$. Sometimes we omit indices if this does not lead to confusion. By $I = I_{\mathfrak{H}}$ we denote the identity operator in $\mathfrak{H}$. If $\mathfrak{N}$ is a subspace in $\mathfrak{H}$, then $\mathfrak{N}^\perp := \mathfrak{H} \ominus \mathfrak{N}$. If P is the orthogonal projection of $\mathfrak{H}$ onto $\mathfrak{N}$, then we write $P^\perp$ for the

orthogonal projection onto $\mathfrak{N}^\perp$. By $H^s(\mathcal{D})$, $s \in \mathbb{R}$, we denote the Sobolev spaces in the domain $\mathcal{D} \subset \mathbb{R}^2$.

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§1. DEFINITION OF THE OPERATOR. MAIN RESULT

1.1. Let $a > 0$, and let

$$\Pi := \{\mathbf{x} = (x_1, x_2) \in \mathbb{R}^2 : 0 < x_2 < a\}.$$

We introduce the notation $\partial_j = \partial/\partial x_j$, $D_j = -i\partial_j$, $j = 1, 2$. Assume that $g(\mathbf{x}) = \text{diag}\{g_1(\mathbf{x}), g_2(\mathbf{x})\}$, where the $g_j(\mathbf{x})$, $j = 1, 2$, are real-valued measurable functions in the strip Π such that

$$(1.1) \quad 0 < c_0 \leq g_j(\mathbf{x}) \leq c_1 < \infty, \quad \mathbf{x} \in \Pi, \quad j = 1, 2.$$

The functions $g_j(x_1, x_2)$ are assumed to be periodic in x_1 with period 1:

$$(1.2) \quad g_j(x_1 + 1, x_2) = g_j(x_1, x_2), \quad \mathbf{x} \in \Pi, \quad j = 1, 2.$$

Next, suppose that the functions $g_j(x_1, x_2)$ are absolutely continuous with respect to the variable $x_2 \in [0, a]$ (with x_1 fixed), and that

$$(1.3) \quad g_j(x_1, 0) = g_j(x_1, a), \quad x_1 \in \mathbb{R}, \quad j = 1, 2,$$

$$(1.4) \quad \text{ess sup}_{\mathbf{x} \in \Pi} |\partial_2 g_j(\mathbf{x})| \leq c_2, \quad j = 1, 2.$$

In the space $L_2(\Pi)$, we consider the operator A given formally by the differential expression

$$\mathcal{A} = -\text{div } g(\mathbf{x}) \text{ grad} = D_1 g_1(\mathbf{x}) D_1 + D_2 g_2(\mathbf{x}) D_2$$

with periodic boundary conditions in x_2 . The precise definition of A as a selfadjoint operator in $L_2(\Pi)$ is given in terms of the closed semibounded quadratic form

$$a[u, u] = \int_{\Pi} (g_1(\mathbf{x}) |D_1 u|^2 + g_2(\mathbf{x}) |D_2 u|^2) d\mathbf{x}, \quad u \in \tilde{H}^1(\Pi).$$

Here $\tilde{H}^s(\Pi)$, $s \geq 0$, is the subspace in $H^s(\Pi)$ formed by all functions the $(a\mathbb{Z})$ -periodic extension of which with respect to variable x_2 belongs to $H_{\text{loc}}^s(\mathbb{R}^2)$.

1.2. The homogenization problem consists in the study of the behavior of a periodic operator in the limit of small period. In $L_2(\Pi)$, we consider the operator A_ε generated by the quadratic form

$$(1.5) \quad a_\varepsilon[u, u] = \int_{\Pi} (g_1(x_1/\varepsilon, x_2) |D_1 u|^2 + g_2(x_1/\varepsilon, x_2) |D_2 u|^2) d\mathbf{x},$$

$$u \in \tilde{H}^1(\Pi), \quad \varepsilon > 0.$$

Our goal is to find the *effective operator* A^0 of the same form as A , but with coefficients independent of the “periodic” variable x_1 and such that

$$(A_\varepsilon + I)^{-1} \rightarrow (A^0 + I)^{-1}, \quad \varepsilon \rightarrow 0.$$

On the formal level, it is easily discovered that the effective operator A^0 is the operator in $L_2(\Pi)$ generated by the form

$$(1.6) \quad a^0[u, u] = \int_{\Pi} (g_1^0(x_2) |D_1 u|^2 + g_2^0(x_2) |D_2 u|^2) d\mathbf{x}, \quad u \in \tilde{H}^1(\Pi),$$

where the coefficients $g_j^0(x_2)$ are given by the formulas

$$(1.7) \quad g_1^0(x_2) := \left(\int_0^1 (g_1(x_1, x_2))^{-1} dx_1 \right)^{-1},$$

$$(1.8) \quad g_2^0(x_2) := \int_0^1 g_2(x_1, x_2) dx_1.$$

We note that conditions (1.1), (1.4) imply the following inequalities for the coefficients (1.7), (1.8):

$$(1.9) \quad c_0 \leq g_j^0(x_2) \leq c_1, \quad x_2 \in [0, a], \quad j = 1, 2,$$

$$(1.10) \quad \operatorname{ess\,sup}_{x_2 \in [0, a]} |\partial_2 g_1^0(x_2)| \leq c_2(c_1/c_0)^2, \quad \operatorname{ess\,sup}_{x_2 \in [0, a]} |\partial_2 g_2^0(x_2)| \leq c_2.$$

1.3. The main result of the paper is the following theorem.

Theorem 1.1. *Suppose that conditions (1.1)–(1.4) are satisfied. Let A_ε be the operator in $L_2(\Pi)$ that corresponds to the form (1.5), and let the operator A^0 correspond to the form (1.6), where the coefficients are defined by (1.7), (1.8). Then the following estimate is true:*

$$(1.11) \quad \|(A_\varepsilon + I)^{-1} - (A^0 + I)^{-1}\|_{L_2(\Pi) \rightarrow L_2(\Pi)} \leq C\varepsilon, \quad 0 < \varepsilon \leq 1.$$

The constant C in (1.11) depends only on the constants c_0 and c_1 in (1.1) and on the constant c_2 in condition (1.4).

1.4. Interpolation. Let $\tilde{H}^{-s}(\Pi)$, $s \geq 0$, be the function space adjoint to $\tilde{H}^s(\Pi)$ with respect to the $L_2(\Pi)$ -duality. We estimate the resolvent $(A_\varepsilon + I)^{-1}$ as an operator from $\tilde{H}^{-1}(\Pi)$ to $\tilde{H}^1(\Pi)$. For $f \in \tilde{H}^{-1}(\Pi)$, we put $u_\varepsilon = (A_\varepsilon + I)^{-1}f$. Then

$$a_\varepsilon[u_\varepsilon, u_\varepsilon] + \|u_\varepsilon\|_{L_2(\Pi)}^2 = \int_\Pi f \overline{u_\varepsilon} d\mathbf{x}.$$

Combining this with the lower estimate (1.1), we arrive at the inequality

$$\min\{c_0, 1\} \|u_\varepsilon\|_{H^1(\Pi)} \leq \|f\|_{\tilde{H}^{-1}(\Pi)}.$$

In other words,

$$\|(A_\varepsilon + I)^{-1}\|_{\tilde{H}^{-1}(\Pi) \rightarrow \tilde{H}^1(\Pi)} \leq \max\{c_0^{-1}, 1\}.$$

A similar estimate for $(A^0 + I)^{-1}$ is also true. Consequently,

$$(1.12) \quad \|(A_\varepsilon + I)^{-1} - (A^0 + I)^{-1}\|_{\tilde{H}^{-1}(\Pi) \rightarrow \tilde{H}^1(\Pi)} \leq 2 \max\{c_0^{-1}, 1\}.$$

Standard interpolation between (1.11) and (1.12) yields the following result.

Theorem 1.2. *Under the conditions of Theorem 1.1, we have*

$$\|(A_\varepsilon + I)^{-1} - (A^0 + I)^{-1}\|_{\tilde{H}^{-s}(\Pi) \rightarrow \tilde{H}^s(\Pi)} \leq C_s \varepsilon^{1-s}, \quad 0 \leq s < 1, \quad 0 < \varepsilon \leq 1.$$

Here $C_s = C^{1-s} (2 \max\{c_0^{-1}, 1\})^s$.

§2. REDUCTION TO OPERATORS ON THE CELL

2.1. Direct integral expansion of the operator A is obtained with the help of the Gelfand transformation $\mathcal{U}$. Putting $\Omega = (0, 1) \times (0, a)$, we consider the Hilbert space

$$(2.1) \quad \mathcal{K} = \int_{[-\pi, \pi)} \oplus L_2(\Omega) dk = L_2([-\pi, \pi); L_2(\Omega)),$$

which is a direct integral with constant fibers. Initially, the operator $\mathcal{U} : L_2(\Pi) \rightarrow \mathcal{K}$ is defined on the functions $f \in C_0^\infty(\Pi)$ by the formula

$$(\mathcal{U}f)(\mathbf{x}, k) = (2\pi)^{-1/2} \sum_{m \in \mathbb{Z}} e^{-ik(x_1+m)} f(x_1+m, x_2), \quad \mathbf{x} \in \Omega, \quad k \in [-\pi, \pi),$$

and then $\mathcal{U}$ extends by continuity to a *unitary mapping from $L_2(\Pi)$ onto $\mathcal{K}$* .

In $L_2(\Omega)$, we consider the operator $A(k)$ generated by the quadratic form

$$a(k)[u, u] = \int_{\Omega} (g_1(\mathbf{x})|(D_1 + k)u|^2 + g_2(\mathbf{x})|D_2 u|^2) d\mathbf{x}, \quad u \in \tilde{H}^1(\Omega), \quad k \in \mathbb{R}.$$

Here $\tilde{H}^1(\Omega)$ is the subspace in $H^1(\Omega)$ formed by the functions the $(\mathbb{Z} \times (a\mathbb{Z}))$ -periodic extensions of which belong to $H_{\text{loc}}^1(\mathbb{R}^2)$.

The operator A is unitarily equivalent to the operator of multiplication by $A(k)$ in $\mathcal{K}$, which is expressed by the relation

$$\mathcal{U}A\mathcal{U}^{-1} = \int_{[-\pi, \pi)} \oplus A(k) dk.$$

2.2. We denote by T_ε , $\varepsilon > 0$, the unitary scale transformation in $L_2(\Pi)$ defined by the formula

$$(T_\varepsilon u)(y_1, x_2) = \varepsilon^{1/2} u(\varepsilon y_1, x_2), \quad y_1 \in \mathbb{R}, \quad 0 < x_2 < a.$$

Let $A(k, \varepsilon)$ be the operator in $L_2(\Omega)$ corresponding to the form

$$(2.2) \quad a(k, \varepsilon)[u, u] = \int_{\Omega} (g_1(\mathbf{x})|(D_1 + k)u|^2 + \varepsilon^2 g_2(\mathbf{x})|D_2 u|^2) d\mathbf{x},$$

$$u \in \tilde{H}^1(\Omega), \quad k \in \mathbb{R}, \quad \varepsilon > 0.$$

We put

$$(2.3) \quad B(k, \varepsilon) := \varepsilon^2 (A(k, \varepsilon) + \varepsilon^2 I)^{-1}, \quad k \in \mathbb{R}, \quad \varepsilon > 0.$$

The operator of multiplication by the operator-valued function (2.3) acting in the direct integral (2.1) is denoted by B_ε :

$$B_\varepsilon := \int_{[-\pi, \pi)} \oplus B(k, \varepsilon) dk, \quad \varepsilon > 0.$$

It is easily seen [7, (5.6)] that

$$(2.4) \quad (A_\varepsilon + I)^{-1} = T_\varepsilon^* \mathcal{U}^* B_\varepsilon \mathcal{U} T_\varepsilon, \quad \varepsilon > 0.$$

Relation (2.4) reduces the problem to the study of the operator family (2.3) acting in $L_2(\Omega)$.

2.3. Let $A^0(k, \varepsilon)$ be the operator in $L_2(\Omega)$ corresponding to the form

$$a^0(k, \varepsilon)[u, u] = \int_{\Omega} (g_1^0(x_2)|D_1 + k|u|^2 + \varepsilon^2 g_2^0(x_2)|D_2 u|^2) dx, \quad u \in \tilde{H}^1(\Omega).$$

We put

$$(2.5) \quad B^0(k, \varepsilon) := \varepsilon^2(A^0(k, \varepsilon) + \varepsilon^2 I)^{-1}, \quad \varepsilon > 0,$$

$$B_{\varepsilon}^0 := \int_{[-\pi, \pi)} \oplus B^0(k, \varepsilon) dk, \quad \varepsilon > 0.$$

As in (2.4), we have

$$(2.6) \quad (A^0 + I)^{-1} = T_{\varepsilon}^* \mathcal{U}^* B_{\varepsilon}^0 \mathcal{U} T_{\varepsilon}, \quad \varepsilon > 0.$$

Theorem 1.1 is a direct consequence of relations (2.4), (2.6) and the following theorem.

Theorem 2.1. *Suppose that conditions (1.1), (1.3), and (1.4) are satisfied. Let $B(k, \varepsilon)$ and $B^0(k, \varepsilon)$ be the operators defined by (2.3) and (2.5). Then*

$$(2.7) \quad \max_{k \in [-\pi, \pi)} \|B(k, \varepsilon) - B^0(k, \varepsilon)\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq C\varepsilon, \quad 0 < \varepsilon \leq 1.$$

The constant C in (2.7) depends only on the constants c_0 and c_1 in (1.1) and on the constant c_2 in (1.4).

§3. THE OPERATOR PENCIL $A_1(k)$

3.1. If we formally put $\varepsilon = 0$, then the operator $A(k, \varepsilon)$ defined in Subsection 2.2 will turn into the operator $A_1(k)$ given by the expression $(D_1 + k)g_1(\mathbf{x})(D_1 + k)$ and periodic boundary conditions in x_1 , and the boundary conditions in x_2 “disappear”. Strictly speaking, $A_1(k)$ is the selfadjoint operator in $L_2(\Omega)$ generated by the quadratic form

$$(3.1) \quad a_1(k)[u, u] = \int_{\Omega} g_1(\mathbf{x})|(D_1 + k)u|^2 dx,$$

$$(3.2) \quad \text{Dom } a_1(k) = \int_{(0, a)} \oplus \tilde{H}^1(0, 1) dx_2 = L_2((0, a); \tilde{H}^1(0, 1)).$$

Here $\tilde{H}^1(0, 1)$ is the subspace in $H^1(0, 1)$ formed by the functions whose $\mathbb{Z}$ -periodic extensions belong to $H_{\text{loc}}^1(\mathbb{R})$.

In this section we study the operator pencil $A_1(k)$. Here, it suffices to impose only condition (1.1) with $j = 1$. We shall not use conditions (1.2)–(1.4). Now it is convenient to interpret the space $L_2(\Omega)$ as a direct integral in x_2 with constant fibers $L_2(0, 1)$:

$$(3.3) \quad L_2(\Omega) = \int_{(0, a)} \oplus L_2(0, 1) dx_2.$$

The operator $A_1(k)$ acts in the direct integral (3.3) “layerwise”, i.e., it is the operator of multiplication by the operator-valued function $A_1(k; x_2)$:

$$(3.4) \quad A_1(k) = \int_{(0, a)} \oplus A_1(k; x_2) dx_2,$$

where $A_1(k; x_2)$ is the selfadjoint operator in $L_2(0, 1)$ generated by the form

$$(3.5) \quad a_1(k; x_2)[v, v] = \int_0^1 g_1(x_1, x_2)|(D_1 + k)v|^2 dx_1, \quad v \in \tilde{H}^1(0, 1), \quad x_2 \in (0, a).$$

3.2. The method of the papers [7, 8] is applicable to the operator family $A_1(k; x_2)$ acting in $L_2(0, 1)$. This family admits the following factorization:

$$(3.6) \quad A_1(k; x_2) = (X(k; x_2))^* X(k; x_2),$$

$$(3.7) \quad X(k; x_2) = X_0(x_2) + kX_1(x_2),$$

where $X_0(x_2)$ is the (closed) operator in $L_2(0, 1)$ given by the expression $(g_1(x_1, x_2))^{\frac{1}{2}} D_1$ on the domain $\tilde{H}^1(0, 1)$, and $X_1(x_2)$ is the (bounded) operator of multiplication by the function $(g_1(x_1, x_2))^{1/2}$. Then the operator (3.7) is closed in $L_2(0, 1)$ on the domain $\tilde{H}^1(0, 1)$. Obviously, the kernel

$$(3.8) \quad \tilde{\mathfrak{N}} = \text{Ker } A_1(0; x_2) = \text{Ker } X_0(x_2) = \{u \in L_2(0, 1) : u = \text{const}\}$$

is one-dimensional. Let $d_0(x_2)$ denote the distance between the point $\lambda = 0$ and the rest of the spectrum of $A_1(0; x_2)$. We estimate $d_0(x_2)$ from below.

By condition (1.1) (with $j = 1$), the quadratic form (3.5) with $k = 0$ satisfies the estimate

$$(3.9) \quad a_1(0; x_2)[v, v] \geq c_0 \int_0^1 |D_1 v(x_1)|^2 dx_1, \quad v \in \tilde{H}^1(0, 1).$$

Expanding the function $v \in \tilde{H}^1(0, 1)$ in the Fourier series

$$v(x_1) = \sum_{m \in \mathbb{Z}} \hat{v}_m e^{2\pi i m x_1}$$

and using (3.9), we obtain

$$a_1(0; x_2)[v, v] \geq c_0 \sum_{m \in \mathbb{Z}} 4\pi^2 m^2 |\hat{v}_m|^2, \quad v \in \tilde{H}^1(0, 1).$$

If $v \perp \tilde{\mathfrak{N}}$ in $L_2(0, 1)$, then $\hat{v}_0 = 0$, whence

$$a_1(0; x_2)[v, v] \geq 4\pi^2 c_0 \|v\|_{L_2(0, 1)}^2, \quad v \in \tilde{H}^1(0, 1), \quad v \perp \tilde{\mathfrak{N}}.$$

This implies the inequality

$$d_0(x_2) \geq d_0 := 4\pi^2 c_0.$$

In accordance with [7, §1], we choose δ such that

$$0 < \delta < d_0/8 = \pi^2 c_0/2.$$

For instance, we can put

$$\delta = \pi^2 c_0/4.$$

Next, we estimate the norm of the operator $X_1(x_2)$. By the upper estimate (1.1) (with $j = 1$), we have

$$(3.10) \quad \|X_1(x_2)\|_{L_2(0, 1) \rightarrow L_2(0, 1)} \leq \sup_{0 < x_1 < 1} (g_1(x_1, x_2))^{1/2} \leq c_1^{1/2}$$

for all $x_2 \in (0, a)$. As in [7, (1.5)], we choose a number $t_0 = t_0(\delta)$ such that $t_0(\delta) \leq \delta^{1/2} \|X_1(x_2)\|^{-1}$ for all $x_2 \in (0, a)$. It suffices to put

$$(3.11) \quad t_0 = \delta^{1/2} c_1^{-1/2} = \pi c_0^{1/2} / 2 c_1^{1/2}.$$

Let $F(k; x_2; \sigma)$ be the spectral projection of the operator $A_1(k; x_2)$ corresponding to the closed interval $[0, \sigma]$. From [7, Proposition 1.2] it follows that

$$(3.12) \quad \dim F(k; x_2; \delta) L_2(0, 1) = \dim F(k; x_2; 3\delta) L_2(0, 1) = 1, \quad |k| = t \leq t_0.$$

Therefore, for all $x_2 \in (0, a)$ the interval $(\delta, 3\delta)$ is free of the spectrum of the operator $A_1(k; x_2)$ if $|k| \leq t_0$. We put $F(k; x_2) := F(k; x_2; \delta)$.

3.3. The spectral germ $S(x_2)$ of the operator family $A_1(k; x_2)$. In accordance with [7, 8], we introduce the operator $S(x_2)$ acting in the one-dimensional space $\tilde{\mathfrak{N}}$ (see (3.8)) and called the spectral germ of the operator family $A_1(k; x_2)$ at $k = 0$. In our case (i.e., for the one-dimensional elliptic operator of the form $-\frac{d}{dx}g(x)\frac{d}{dx}$), the germ is calculated explicitly (see [7, Subsection 6.1]): the operator $S(x_2)$ is multiplication by the number $g_1^0(x_2)$ (see (1.7)).

Let $\tilde{P}$ be the orthoprojection in $L_2(0, 1)$ onto the one-dimensional subspace $\tilde{\mathfrak{N}}$. By [7, (2.13)–(2.15)], the following approximations are valid:

$$(3.13) \quad F(k; x_2) = \tilde{P} + t\Phi(k; x_2), \quad t = |k| \leq t_0, \quad x_2 \in (0, a),$$

$$(3.14) \quad A_1(k; x_2)F(k; x_2) = t^2S(x_2)\tilde{P} + t^3\Psi(k; x_2), \quad t = |k| \leq t_0, \quad x_2 \in (0, a).$$

Here, the operators $\Phi(k; x_2)$ and $\Psi(k; x_2)$ admit the estimates

$$(3.15) \quad \|\Phi(k; x_2)\|_{L_2(0,1) \rightarrow L_2(0,1)} \leq C_1, \quad t = |k| \leq t_0, \quad x_2 \in (0, a),$$

$$(3.16) \quad \|\Psi(k; x_2)\|_{L_2(0,1) \rightarrow L_2(0,1)} \leq C_2, \quad t = |k| \leq t_0, \quad x_2 \in (0, a).$$

The constants C_1, C_2 are controlled explicitly. By [7, (1.4.3), (1.4.6)], and (3.10), we have

$$C_1 = \beta_1 \delta^{-1/2} \max_{x_2 \in (0, a)} \|X_1(x_2)\| \leq \beta_1 \delta^{-1/2} c_1^{1/2},$$

$$C_2 = \beta_2 \delta^{-1/2} \max_{x_2 \in (0, a)} \|X_1(x_2)\|^3 \leq \beta_2 \delta^{-1/2} c_1^{3/2},$$

where β_1 and β_2 are absolute constants.

3.4. The spectral germ of the operator family $A_1(k)$. We return to the family (3.4) in $L_2(\Omega)$. First, we note that the interval $(\delta, 3\delta)$ is free of the spectrum of $A_1(k)$ for $|k| \leq t_0$ (this follows from (3.4) and (3.12)).

By (3.4) and (3.6), we have

$$A_1(k) = X(k)^* X(k),$$

where

$$X(k) = \int_{(0, a)} \oplus X(k; x_2) dx_2.$$

Here, by (3.7),

$$X(k) = X_0 + kX_1,$$

$$X_0 := \int_{(0, a)} \oplus X_0(x_2) dx_2,$$

$$X_1 := \int_{(0, a)} \oplus X_1(x_2) dx_2.$$

Clearly, X_0 is the operator in $L_2(\Omega)$ given by the expression $(g_1(\mathbf{x}))^{1/2} D_1$ on the domain (3.2), and X_1 is the operator of multiplication by the function $(g_1(\mathbf{x}))^{1/2}$.

For the operator $A_1(0)$ with $k = 0$, we have

$$(3.17) \quad \mathfrak{N} = \text{Ker } A_1(0) = \text{Ker } X_0 = \{u \in L_2(\Omega) : u = u(x_2)\}.$$

Thus, the kernel $\mathfrak{N}$ is isomorphic to $L_2(0, a)$ and consists of functions depending only on the second variable. Obviously,

$$(3.18) \quad \mathfrak{N} = \int_{(0, a)} \oplus \tilde{\mathfrak{N}} dx_2,$$

which means that the space $\mathfrak{N} (= L_2(0, a))$ can be viewed as a direct integral of the one-dimensional spaces $\tilde{\mathfrak{N}} (= \mathbb{C})$. It is also clear that the orthoprojection P of $L_2(\Omega)$ onto the kernel $\mathfrak{N}$ can be represented as the direct integral of the orthoprojections $\tilde{P}$ of $L_2(0, 1)$ onto $\tilde{\mathfrak{N}}$:

$$(3.19) \quad P = \int_{(0,a)} \oplus \tilde{P} dx_2.$$

In the space (3.18), we introduce the decomposable operator

$$S = \int_{(0,a)} \oplus S(x_2) dx_2,$$

which is the operator of multiplication by the function $g_1^0(x_2)$. The operator S is called the *spectral germ of the operator family $A_1(k)$ at $k = 0$* .

Let $F(k)$ denote the spectral projection of the operator $A_1(k)$ corresponding to the interval $[0, \delta]$. Since $A_1(k)$ can be represented as the direct integral (3.4), we have

$$(3.20) \quad F(k) = \int_{(0,a)} \oplus F(k; x_2) dx_2.$$

From (3.13), (3.19), and (3.20) it follows that

$$(3.21) \quad F(k) = P + t\Phi(k), \quad t = |k| \leq t_0,$$

$$(3.22) \quad \Phi(k) = \int_{(0,a)} \oplus \Phi(k; x_2) dx_2.$$

Similarly, (3.4), (3.14), (3.19), and (3.20) imply the relations

$$(3.23) \quad A_1(k)F(k) = t^2 SP + t^3 \Psi(k), \quad t = |k| \leq t_0,$$

$$(3.24) \quad \Psi(k) = \int_{(0,a)} \oplus \Psi(k; x_2) dx_2.$$

By (3.15) and (3.16), the operators (3.22) and (3.24) satisfy the estimates

$$(3.25) \quad \|\Phi(k)\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq C_1, \quad t = |k| \leq t_0,$$

$$(3.26) \quad \|\Psi(k)\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq C_2, \quad t = |k| \leq t_0.$$

§4. APPROXIMATION OF THE OPERATOR FAMILY $B(k; \varepsilon)$

4.1. Now we return to the study of the operator family $A(k, \varepsilon)$ introduced in Subsection 2.2. Our goal is to approximate the operator $B(k, \varepsilon)$ defined by (2.3).

In the space $\mathfrak{N} = L_2(0, a)$ (see (3.17)), we consider the quadratic form

$$(4.1) \quad s(k, \varepsilon)[\omega, \omega] = \int_0^a (g_1^0(x_2)t^2|\omega(x_2)|^2 + g_2^0(x_2)\varepsilon^2|D_2\omega(x_2)|^2) dx_2,$$

$$k \in [-\pi, \pi), \quad \varepsilon > 0,$$

$$(4.2) \quad \text{Dom } s(k, \varepsilon) = \tilde{H}^1(0, a).$$

Let $\tilde{H}^s(0, a)$ denote the subspace in $H^s(0, a)$ formed by the functions the $(a\mathbb{Z})$ -periodic extensions of which belong to $H_{\text{loc}}^s(\mathbb{R})$. The form (4.1) is closed and semibounded on the domain (4.2). The selfadjoint operator in $L_2(0, a)$ corresponding to this form is denoted by $S(k, \varepsilon)$. Clearly, under conditions (1.9) and (1.10), the operator $S(k, \varepsilon)$ is given by the differential expression

$$(4.3) \quad S(k, \varepsilon) = g_1^0(x_2)t^2 + \varepsilon^2 D_2 g_2^0(x_2) D_2$$

on the domain

$$(4.4) \quad \text{Dom } S(k, \varepsilon) = \tilde{H}^2(0, a).$$

Note that the first summand in (4.3) corresponds to the operator $t^2 S$.

The key role in the proof of Theorem 2.1 is played by the following theorem.

Theorem 4.1. *Suppose that conditions (1.1), (1.3), and (1.4) are satisfied. Let P be the orthoprojection of $L_2(\Omega)$ onto the subspace $\mathfrak{N}$ of functions depending only on x_2 . Let $A(k, \varepsilon)$ be the operator in $L_2(\Omega)$ corresponding to the quadratic form (2.2), and let $S(k, \varepsilon)$ be the operator in $\mathfrak{N}$ defined by (4.3) and (4.4). Then*

$$(4.5) \quad \varepsilon^2 \| (A(k, \varepsilon) + \varepsilon^2 I)^{-1} - (S(k, \varepsilon)P + \varepsilon^2 I)^{-1} P \|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq \tilde{C} \varepsilon, \\ |k| \leq t_0, \quad 0 < \varepsilon \leq 1.$$

The number t_0 is defined as in (3.11). The constant $\tilde{C}$ depends only on the constants c_0 and c_1 in (1.1) and on the constant c_2 in (1.4).

4.2. Deduction of Theorem 2.1 from Theorem 4.1. Let $A^0(k, \varepsilon)$ be the operator family corresponding to the effective operator A^0 defined in Subsection 2.3. For $A^0(k, \varepsilon)$, we define the corresponding operator $S^0(k, \varepsilon)$ by the same rule as for $A(k, \varepsilon)$. Clearly, $S^0(k, \varepsilon) = S(k, \varepsilon)$. Applying Theorem 4.1 to the operator $A^0(k, \varepsilon)$, we obtain

$$(4.6) \quad \varepsilon^2 \| (A^0(k, \varepsilon) + \varepsilon^2 I)^{-1} - (S(k, \varepsilon)P + \varepsilon^2 I)^{-1} P \|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq \tilde{C}' \varepsilon, \\ |k| \leq t_0, \quad 0 < \varepsilon \leq 1.$$

Since the coefficients $g_j^0(x_2)$, $j = 1, 2$, satisfy estimate (1.9) with the same constants c_0 and c_1 as in (1.1), the number t_0 (see (3.11)) in (4.5) and in (4.6) is one and the same. Moreover, combining (1.9) and (1.10) with the statement of Theorem 4.1 concerning the constant $\tilde{C}$, we see that the constant $\tilde{C}'$ in (4.6) depends only on c_0 , c_1 , and c_2 .

From (4.5) and (4.6) it follows that

$$(4.7) \quad \varepsilon^2 \| (A(k, \varepsilon) + \varepsilon^2 I)^{-1} - (A^0(k, \varepsilon) + \varepsilon^2 I)^{-1} \|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq (\tilde{C} + \tilde{C}') \varepsilon, \\ |k| \leq t_0, \quad 0 < \varepsilon \leq 1.$$

For $|k| > t_0$, the required estimate can easily be obtained by variational arguments. By (1.1) and (2.2), we have

$$(4.8) \quad a(k, \varepsilon)[u, u] \geq c_0 \int_{\Omega} (|(D_1 + k)u|^2 + \varepsilon^2 |D_2 u|^2) d\mathbf{x}, \\ u \in \tilde{H}^1(\Omega), \quad k \in [-\pi, \pi), \quad \varepsilon > 0.$$

Expanding the function $u \in \tilde{H}^1(\Omega)$ in the Fourier series with respect to variable x_1 ,

$$u(x_1, x_2) = \sum_{m \in \mathbb{Z}} \hat{u}_m(x_2) e^{2\pi i m x_1},$$

we obtain

$$(4.9) \quad a(k, \varepsilon)[u, u] \geq c_0 \sum_{m \in \mathbb{Z}} (2\pi m + k)^2 \int_0^a |\hat{u}_m(x_2)|^2 dx_2, \\ u \in \tilde{H}^1(\Omega), \quad k \in [-\pi, \pi), \quad \varepsilon > 0.$$

Obviously, if $k \in [-\pi, -t_0] \cup [t_0, \pi)$ (see (3.11)), then

$$|2\pi m + k| \geq t_0, \quad m \in \mathbb{Z}.$$

Then (4.9) implies that

$$a(k, \varepsilon)[u, u] \geq c_0 t_0^2 \|u\|_{L_2(\Omega)}^2, \quad u \in \tilde{H}^1(\Omega), \quad k \in [-\pi, -t_0] \cup [t_0, \pi), \quad \varepsilon > 0.$$

Consequently,

$$\|(A(k, \varepsilon) + \varepsilon^2 I)^{-1}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq c_0^{-1} t_0^{-2}, \quad k \in [-\pi, -t_0] \cup [t_0, \pi], \quad \varepsilon > 0.$$

The same estimate is true with $A(k, \varepsilon)$ replaced by $A^0(k, \varepsilon)$. Combining this with (4.7), we obtain (2.7) with the constant $C = \max\{\tilde{C} + \tilde{C}', 2c_0^{-1}t_0^{-2}\}$ depending only on c_0, c_1 , and c_2 .

Theorem 4.1 is proved in §§5–7.

§5. PRELIMINARY ESTIMATES

5.1. By (4.9), we have

$$(5.1) \quad \begin{aligned} a(k, \varepsilon)[u, u] + \varepsilon^2 \|u\|_{L_2(\Omega)}^2 &\geq (c_0 t^2 + \varepsilon^2) \|u\|_{L_2(\Omega)}^2 \geq c_*(t^2 + \varepsilon^2) \|u\|_{L_2(\Omega)}^2, \\ u &\in \tilde{H}^1(\Omega), \quad k \in [-\pi, \pi], \quad \varepsilon > 0, \end{aligned}$$

where $c_* = \min\{c_0, 1\}$.

As in (4.8), from (1.1) and (3.1) it follows that

$$(5.2) \quad a_1(k)[u, u] \geq c_0 \int_{\Omega} |(D_1 + k)u|^2 d\mathbf{x}, \quad u \in \tilde{H}^1(\Omega), \quad k \in [-\pi, \pi],$$

whence

$$(5.3) \quad \begin{aligned} a_1(k)[u, u] + \varepsilon^2 \|u\|_{L_2(\Omega)}^2 &\geq (c_0 t^2 + \varepsilon^2) \|u\|_{L_2(\Omega)}^2 \geq c_*(t^2 + \varepsilon^2) \|u\|_{L_2(\Omega)}^2, \\ u &\in \text{Dom } a_1(k), \quad k \in [-\pi, \pi], \quad \varepsilon > 0. \end{aligned}$$

By (1.9), the form (4.1) satisfies the estimate

$$(5.4) \quad \begin{aligned} s(k, \varepsilon)[\omega, \omega] &\geq c_0 \int_0^a (t^2 |\omega|^2 + \varepsilon^2 |D_2 \omega|^2) dx_2, \\ \omega &\in \tilde{H}^1(0, a), \quad k \in [-\pi, \pi], \quad \varepsilon > 0. \end{aligned}$$

Then

$$(5.5) \quad \begin{aligned} s(k, \varepsilon)[\omega, \omega] + \varepsilon^2 \|\omega\|_{L_2(0, a)}^2 &\geq (c_0 t^2 + \varepsilon^2) \|\omega\|_{L_2(0, a)}^2 \geq c_*(t^2 + \varepsilon^2) \|\omega\|_{L_2(0, a)}^2, \\ \omega &\in \tilde{H}^1(0, a), \quad k \in [-\pi, \pi], \quad \varepsilon > 0. \end{aligned}$$

Relations (5.1), (5.3), and (5.5) imply that

$$(5.6) \quad \|(A(k, \varepsilon) + \varepsilon^2 I)^{-1}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq c_*^{-1} (t^2 + \varepsilon^2)^{-1},$$

$$\|(A_1(k) + \varepsilon^2 I)^{-1}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq c_*^{-1} (t^2 + \varepsilon^2)^{-1},$$

$$(5.7) \quad \|(S(k, \varepsilon)P + \varepsilon^2 I)^{-1}P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq c_*^{-1} (t^2 + \varepsilon^2)^{-1}$$

for all $k \in [-\pi, \pi]$ and $\varepsilon \geq 0$.

Next, from the estimate (see (4.8))

$$\begin{aligned} c_0 \int_{\Omega} (|(D_1 + k)u|^2 + \varepsilon^2 |D_2 u|^2) d\mathbf{x} \\ \leq a(k, \varepsilon)[u, u] + \varepsilon^2 \|u\|_{L_2(\Omega)}^2 = \|(A(k, \varepsilon) + \varepsilon^2 I)^{1/2} u\|_{L_2(\Omega)}^2, \quad u \in \tilde{H}^1(\Omega), \end{aligned}$$

it follows that

$$(5.8) \quad \|(D_1 + k)(A(k, \varepsilon) + \varepsilon^2 I)^{-1/2}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq c_0^{-1/2},$$

$$(5.9) \quad \|\varepsilon D_2(A(k, \varepsilon) + \varepsilon^2 I)^{-1/2}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq c_0^{-1/2}$$

for all $k \in [-\pi, \pi]$ and $\varepsilon > 0$. Similarly, (5.2) implies

$$(5.10) \quad \|(D_1 + k)A_1(k)^{-1/2}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq c_0^{-1/2}, \quad k \in [-\pi, \pi].$$

Finally, from (5.4) we obtain

$$(5.11) \quad \|\varepsilon D_2(S(k, \varepsilon)P + \varepsilon^2 I)^{-1/2} P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq c_0^{-1/2}$$

for all $k \in [-\pi, \pi)$ and $\varepsilon > 0$.

5.2. Recall that $F(k)$ denotes the spectral projection of the operator $A_1(k)$ for the interval $[0, \delta]$, and P denotes the orthoprojection of $L_2(\Omega)$ onto the subspace $\mathfrak{N}$ of functions depending only on x_2 .

Lemma 5.1. *Let $t = |k| \leq t_0$, $\varepsilon > 0$. Then*

$$(5.12) \quad \begin{aligned} & \|F(k)^\perp (A(k, \varepsilon) + \varepsilon^2 I)^{-1/2}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ & = \|(A(k, \varepsilon) + \varepsilon^2 I)^{-1/2} F(k)^\perp\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq \delta^{-1/2}; \end{aligned}$$

$$(5.13) \quad \|(A(k, \varepsilon) + \varepsilon^2 I)^{-1/2} P^\perp\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq C_3 := \delta^{-1/2} + C_1 c_*^{-1/2};$$

$$(5.14) \quad \|F(k)^\perp (S(k, \varepsilon)P + \varepsilon^2 I)^{-1/2} P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq C_4 := C_1 c_*^{-1/2}.$$

Proof. By (2.2) and (3.1), we have

$$a_1(k)[u, u] + \varepsilon^2 \|u\|_{L_2(\Omega)}^2 \leq a(k, \varepsilon)[u, u] + \varepsilon^2 \|u\|_{L_2(\Omega)}^2, \quad u \in \tilde{H}^1(\Omega).$$

Consequently,

$$\|(A_1(k) + \varepsilon^2 I)^{1/2} (A(k, \varepsilon) + \varepsilon^2 I)^{-1/2}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq 1.$$

Therefore,

$$\begin{aligned} & \|F(k)^\perp (A(k, \varepsilon) + \varepsilon^2 I)^{-1/2}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ & \leq \|F(k)^\perp (A_1(k) + \varepsilon^2 I)^{-1/2}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ & \quad \times \|(A_1(k) + \varepsilon^2 I)^{1/2} (A(k, \varepsilon) + \varepsilon^2 I)^{-1/2}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ & \leq \delta^{-1/2}, \end{aligned}$$

which proves (5.12).

In order to prove (5.13), we use (5.12), (5.6) and (3.21), (3.25):

$$\begin{aligned} & \|(A(k, \varepsilon) + \varepsilon^2 I)^{-1/2} P^\perp\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ & = \|(A(k, \varepsilon) + \varepsilon^2 I)^{-1/2} (F(k)^\perp + t\Phi(k))\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ & \leq \delta^{-1/2} + C_1 c_*^{-1/2} t(t^2 + \varepsilon^2)^{-1/2} \leq C_3. \end{aligned}$$

Inequality (5.14) follows from (3.21), (3.25), (5.7), and the obvious relation

$$P^\perp (S(k, \varepsilon)P + \varepsilon^2 I)^{-1/2} P = 0.$$

We have

$$\begin{aligned} & \|F(k)^\perp (S(k, \varepsilon)P + \varepsilon^2 I)^{-1/2} P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ & = \|t\Phi(k) (S(k, \varepsilon)P + \varepsilon^2 I)^{-1/2} P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ & \leq C_1 c_*^{-1/2} t(t^2 + \varepsilon^2)^{-1/2} \leq C_4. \quad \square \end{aligned}$$

Lemma 5.1 and inequalities (5.6) and (5.7) imply the following statement.

Corollary 5.2. *Let $t = |k| \leq t_0$, $\varepsilon > 0$. Then*

$$(5.15) \quad \varepsilon^2 \|(A(k, \varepsilon) + \varepsilon^2 I)^{-1} P^\perp\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq C_3 c_*^{-1/2} \varepsilon,$$

$$(5.16) \quad \varepsilon^2 \|F(k)^\perp (S(k, \varepsilon)P + \varepsilon^2 I)^{-1} P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq C_4 c_*^{-1/2} \varepsilon.$$

Relations (5.15) and (5.16) show that, for the proof of Theorem 4.1, it suffices to estimate the operator $\varepsilon^2 J(k, \varepsilon)$ with

$$(5.17) \quad J(k, \varepsilon) := (A(k, \varepsilon) + \varepsilon^2 I)^{-1} P - F(k)(S(k, \varepsilon)P + \varepsilon^2 I)^{-1} P.$$

Thus, *Theorem 4.1 is a consequence of the following proposition.*

Proposition 5.3. *Suppose that conditions of Theorem 4.1 are satisfied. Let $J(k, \varepsilon)$ be the operator defined by (5.17). Then*

$$(5.18) \quad \varepsilon^2 \|J(k, \varepsilon)\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq \widehat{C}\varepsilon, \quad |k| \leq t_0, \quad 0 < \varepsilon \leq 1.$$

The constant $\widehat{C}$ depends only on c_0 , c_1 , and c_2 .

Proposition 5.3 will be proved in §7.

§6. ESTIMATES OF COMMUTATORS

6.1. For the proof of Proposition 5.3, we need to estimate the commutator of the operator D_2 with the resolvent $(A(k, \varepsilon) + \varepsilon^2 I)^{-1}$ and with the spectral projection $F(k)$ of the operator $A_1(k)$.

Lemma 6.1. *Let $|k| \leq t_0$, $\varepsilon > 0$. The operator $\text{clos}[(A(k, \varepsilon) + \varepsilon^2 I)^{-1}, D_2]$ is bounded on $L_2(\Omega)$, and*

$$(6.1) \quad \begin{aligned} & \|[(A(k, \varepsilon) + \varepsilon^2 I)^{-1}, D_2]P^\perp\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ & = \|P^\perp[(A(k, \varepsilon) + \varepsilon^2 I)^{-1}, D_2]\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq C_5(t^2 + \varepsilon^2)^{-1/2}. \end{aligned}$$

The constant C_5 depends on c_0 , c_1 , and c_2 .

Proof. Under conditions (1.1), (1.3), and (1.4), for the functions $u \in \text{Dom } A(k, \varepsilon)$ we have

$$u \in \widetilde{H}^1(\Omega), \quad D_2 u \in \widetilde{H}^1(\Omega).$$

This property expresses a partial smoothness improvement (with respect to x_2) of the solutions of the elliptic equation $-\text{div } g(\mathbf{x})\nabla u = f(\mathbf{x})$ with periodic boundary conditions and $f \in L_2(\Omega)$. This is ensured by the fact that the matrix of coefficients $g(\mathbf{x})$ is periodic and satisfies the Lipschitz condition (with respect to x_2). The corresponding theorem about smoothness improvement can easily be established by analogy with the proof of Theorem 10.1 in Chapter III of the book [9] or Theorem 8.8 in [10].

Let $u, v \in \text{Dom } A(k, \varepsilon)$. Then $D_2 u, D_2 v \in \text{Dom } a(k, \varepsilon)$, and

$$(6.2) \quad \begin{aligned} & a(k, \varepsilon)[D_2 u, v] + \varepsilon^2(D_2 u, v)_{L_2(\Omega)} - a(k, \varepsilon)[u, D_2 v] - \varepsilon^2(u, D_2 v)_{L_2(\Omega)} \\ & = \int_{\Omega} \left(g_1(\mathbf{x})((D_1 + k)D_2 u) \overline{(D_1 + k)v} + \varepsilon^2 g_2(\mathbf{x})(D_2^2 u) \overline{D_2 v} \right) d\mathbf{x} \\ & \quad + \varepsilon^2 \int_{\Omega} ((D_2 u)\overline{v} - u \overline{(D_2 v)}) d\mathbf{x} \\ & \quad - \int_{\Omega} \left(g_1(\mathbf{x})((D_1 + k)u) \overline{(D_1 + k)D_2 v} + \varepsilon^2 g_2(\mathbf{x})(D_2 u) \overline{D_2^2 v} \right) d\mathbf{x} \\ & = - \int_{\Omega} \left((D_2 g_1)(\mathbf{x})((D_1 + k)u) \overline{(D_1 + k)v} + \varepsilon^2 (D_2 g_2)(\mathbf{x})(D_2 u) \overline{D_2 v} \right) d\mathbf{x}. \end{aligned}$$

We have integrated by parts and used the periodic boundary conditions (in x_2) on u and v , and also the periodicity of the coefficients $g_j(\mathbf{x})$ with respect to x_2 (see (1.3)).

Now we plug the functions $u = (A(k, \varepsilon) + \varepsilon^2 I)^{-1} \varphi$, $v = (A(k, \varepsilon) + \varepsilon^2 I)^{-1} \psi$ in the relation obtained. Here $\varphi, \psi \in L_2(\Omega)$. Then the left-hand side of (6.2) takes the form

$$\begin{aligned} & (D_2(A(k, \varepsilon) + \varepsilon^2 I)^{-1} \varphi, \psi)_{L_2(\Omega)} - (\varphi, D_2(A(k, \varepsilon) + \varepsilon^2 I)^{-1} \psi)_{L_2(\Omega)} \\ &= (\text{clos}[D_2, (A(k, \varepsilon) + \varepsilon^2 I)^{-1}] \varphi, \psi)_{L_2(\Omega)}, \end{aligned}$$

and the right-hand side looks like this:

$$\begin{aligned} & - \int_{\Omega} (D_2 g_1)(\mathbf{x}) ((D_1 + k)(A(k, \varepsilon) + \varepsilon^2 I)^{-1} \varphi) \overline{((D_1 + k)(A(k, \varepsilon) + \varepsilon^2 I)^{-1} \psi)} d\mathbf{x} \\ & - \int_{\Omega} \varepsilon^2 (D_2 g_2)(\mathbf{x}) (D_2(A(k, \varepsilon) + \varepsilon^2 I)^{-1} \varphi) \overline{(D_2(A(k, \varepsilon) + \varepsilon^2 I)^{-1} \psi)} d\mathbf{x}. \end{aligned}$$

Thus,

$$\begin{aligned} & \text{clos}[D_2, (A(k, \varepsilon) + \varepsilon^2 I)^{-1}] \\ (6.3) \quad &= -((D_1 + k)(A(k, \varepsilon) + \varepsilon^2 I)^{-1})^* ((D_2 g_1)(D_1 + k)(A(k, \varepsilon) + \varepsilon^2 I)^{-1}) \\ & - (\varepsilon D_2(A(k, \varepsilon) + \varepsilon^2 I)^{-1})^* ((D_2 g_2) \varepsilon D_2(A(k, \varepsilon) + \varepsilon^2 I)^{-1}). \end{aligned}$$

By (1.4), (5.6), (5.8), and (5.9), the operator on the right is bounded.

From (1.4), (5.6), (5.8), (5.9), and (5.13) it follows that the operator (6.3) multiplied by $P^\perp$ from the right satisfies the estimate

$$\|[D_2, (A(k, \varepsilon) + \varepsilon^2 I)^{-1}] P^\perp\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq 2c_0^{-1} c_*^{-1/2} C_3 c_2 (t^2 + \varepsilon^2)^{-1/2},$$

which proves (6.1) with $C_5 = 2c_0^{-1} c_*^{-1/2} C_3 c_2$. $\square$

6.2. For calculating the commutator $[D_2, F(k)]$, we represent the spectral projection $F(k)$ as the integral of the resolvent $(A_1(k) - zI)^{-1}$ over the *contour* Γ that envelopes the real interval $[0, \delta]$ equidistantly with distance δ . We recall that, for $|k| \leq t_0$, the interval $(\delta, 3\delta)$ is free of the spectrum of the operator $A_1(k)$; therefore,

$$(6.4) \quad \|(A_1(k) - zI)^{-1}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq \delta^{-1}, \quad |k| \leq t_0, \quad z \in \Gamma.$$

We have

$$(6.5) \quad F(k) = -\frac{1}{2\pi i} \int_{\Gamma} (A_1(k) - zI)^{-1} dz.$$

Lemma 6.2. *Let $|k| \leq t_0$, $\varepsilon > 0$. The operator $\text{clos}[D_2, F(k)]$ is bounded on $L_2(\Omega)$, and*

$$(6.6) \quad \|[D_2, F(k)]\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq C_6,$$

$$(6.7) \quad \|[D_2, F(k)]F(k)\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq C_7 t,$$

where the constants C_6 and C_7 depend on c_0 , c_1 , and c_2 .

Proof. First, we calculate the commutator $\text{clos}[D_2, (A_1(k) - zI)^{-1}]$. Suppose functions $u, v \in \text{Dom } a_1(k)$ (see (3.2)) are such that $D_2 u, D_2 v \in \text{Dom } a_1(k)$ and $u(\mathbf{x}), v(\mathbf{x})$ are equal to zero in some neighborhood of the set $\{\mathbf{x} \in \partial\Omega : x_2 = 0 \text{ or } x_2 = a\}$. Then

$$\begin{aligned} & a_1(k)[D_2 u, v] - z(D_2 u, v)_{L_2(\Omega)} - a_1(k)[u, D_2 v] + z(u, D_2 v)_{L_2(\Omega)} \\ &= \int_{\Omega} \left(g_1(\mathbf{x}) ((D_1 + k) D_2 u) \overline{(D_1 + k) v} - g_1(\mathbf{x}) ((D_1 + k) u) \overline{(D_1 + k) D_2 v} \right) d\mathbf{x} \\ (6.8) \quad & + z \int_{\Omega} (u \overline{D_2 v} - (D_2 u) \overline{v}) d\mathbf{x} \\ &= - \int_{\Omega} (D_2 g_1)(\mathbf{x}) (D_1 + k) u \overline{(D_1 + k) v} d\mathbf{x}. \end{aligned}$$

We have integrated by parts and used the fact that u and v vanish near the corresponding part of the boundary. We put $u = (A_1(k) - zI)^{-1}\varphi$, $v = (A_1(k) - zI)^{-1}\psi$, where $\varphi, \psi \in C_0^\infty(\Omega)$. By using (1.4), it is easily seen that u and v satisfy the required conditions. From (6.8) it follows that

$$\begin{aligned} & (D_2(A_1(k) - zI)^{-1}\varphi, \psi)_{L_2(\Omega)} - (\varphi, D_2(A_1(k) - zI)^{-1}\psi)_{L_2(\Omega)} \\ &= - \int_{\Omega} (D_2 g_1)(\mathbf{x}) ((D_1 + k)(A_1(k) - zI)^{-1}\varphi) \overline{((D_1 + k)(A_1(k) - zI)^{-1}\psi)} d\mathbf{x}, \\ & \quad \varphi, \psi \in C_0^\infty(\Omega). \end{aligned}$$

Since $C_0^\infty(\Omega)$ is dense in $L_2(\Omega)$, we obtain

$$\begin{aligned} (6.9) \quad & \text{clos}[D_2, (A_1(k) - zI)^{-1}] \\ &= -((D_1 + k)(A_1(k) - zI)^{-1})^* ((D_2 g_1)(D_1 + k)(A_1(k) - zI)^{-1}). \end{aligned}$$

Now (1.4) and (5.10) yield the estimate

$$\begin{aligned} (6.10) \quad & \| [D_2, (A_1(k) - zI)^{-1}] \|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ & \leq c_2 \| (D_1 + k)A_1(k)^{-1/2} \|^2 \| A_1(k)^{1/2}(A_1(k) - zI)^{-1} \|^2 \\ & \leq c_2 c_0^{-1} \| A_1(k)^{1/2}(A_1(k) - zI)^{-1} \|^2. \end{aligned}$$

Finally, we use the inequality

$$(6.11) \quad \| A_1(k)^{1/2}(A_1(k) - zI)^{-1} \| \leq c(\delta), \quad z \in \Gamma, \quad |k| \leq t_0,$$

where $c(\delta)$ is the maximum value of the function $f(\lambda) = \lambda^{1/2}|\lambda - z|^{-1}$ for $\lambda \in [0, \delta] \cup [3\delta, \infty)$ and $z \in \Gamma$. Inequalities (6.10) and (6.11) imply

$$(6.12) \quad \| [D_2, (A_1(k) - zI)^{-1}] \|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq c_2 c_0^{-1} c(\delta)^2, \quad z \in \Gamma, \quad |k| \leq t_0.$$

By (6.5) and (6.12), we have

$$(6.13) \quad \| [D_2, F(k)] \|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq (2\delta + 2\pi\delta)(2\pi)^{-1} c_2 c_0^{-1} c(\delta)^2 =: C_6, \quad |k| \leq t_0,$$

which proves (6.6).

Relations (1.4), (5.10), (6.4), and (6.11) show that the commutator (6.9) multiplied by $F(k)$ from the right satisfies the estimate

$$\begin{aligned} (6.14) \quad & \| [D_2, (A_1(k) - zI)^{-1}] F(k) \|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ & \leq c_2 \| (D_1 + k)A_1(k)^{-1/2} \|^2 \| A_1(k)^{1/2}(A_1(k) - zI)^{-1} \| \\ & \quad \times \| A_1(k)^{1/2} F(k) \| \| (A_1(k) - zI)^{-1} \| \\ & \leq c_2 c_0^{-1} c(\delta) \delta^{-1} \| A_1(k)^{1/2} F(k) \|, \quad z \in \Gamma, \quad |k| \leq t_0. \end{aligned}$$

Recalling (3.23) and (3.26), we obtain

$$\begin{aligned} & \| A_1(k)^{1/2} F(k) u \|_{L_2(\Omega)}^2 = (A_1(k) F(k) u, u)_{L_2(\Omega)} = ((t^2 S P + t^3 \Psi(k)) u, u)_{L_2(\Omega)} \\ & \leq (t^2 \| S \|_{L_2(\Omega) \rightarrow L_2(\Omega)} + t^3 \| \Psi(k) \|_{L_2(\Omega) \rightarrow L_2(\Omega)}) \| u \|_{L_2(\Omega)}^2 \\ & \leq (c_1 + C_2 t_0) t^2 \| u \|_{L_2(\Omega)}^2, \quad t \leq t_0 \end{aligned}$$

(we have also used the fact that $\| S \|_{L_2(\Omega) \rightarrow L_2(\Omega)} = \max |g_1^0(x_2)| \leq c_1$ by (1.9)). Thus,

$$(6.15) \quad \| A_1(k)^{1/2} F(k) \|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq (c_1 + C_2 t_0)^{1/2} t, \quad |k| \leq t_0.$$

Combining this with (6.5) and (6.14), we arrive at (6.7) with

$$C_7 = (2\delta + 2\pi\delta)(2\pi)^{-1} c_2 c_0^{-1} c(\delta) \delta^{-1} (c_1 + C_2 t_0)^{1/2}. \quad \square$$

6.3. In the sequel, we shall need the following statement.

Lemma 6.3. *We have*

$$(6.16) \quad F(k)\tilde{H}^1(\Omega) \subset \tilde{H}^1(\Omega).$$

Proof. Let $u \in \tilde{H}^1(\Omega)$. Then, obviously, $u \in \text{Dom } a_1(k) = \int_{(0,a)} \oplus \tilde{H}^1(0,1) dx_2$. Since $F(k)\text{Dom } a_1(k) \subset \text{Dom } a_1(k)$, *a fortiori*, we have $F(k)u \in \text{Dom } a_1(k)$. This means that $D_1(F(k)u) \in L_2(\Omega)$, and $F(k)u$ satisfies periodic boundary conditions with respect to x_1 . Next,

$$D_2(F(k)u) = F(k)D_2u + [D_2, F(k)]u \in L_2(\Omega)$$

because $D_2u \in L_2(\Omega)$, and the operators $F(k)$ and $[D_2, F(k)]$ are bounded in $L_2(\Omega)$. The function $F(k)u$ satisfies periodic boundary conditions in x_2 because $u(\mathbf{x})$ is periodic in x_2 , the operator $F(k)$ acts in the direct integral (3.3) layerwise, and the coefficient $g_1(\mathbf{x})$ is periodic with respect to x_2 (see (1.3)). $\square$

§7. PROOF OF PROPOSITION 5.3

7.1. For the study of the operator (5.17), we wish to apply the identity

$$(7.1) \quad \begin{aligned} J(k, \varepsilon) &= (A(k, \varepsilon) + \varepsilon^2 I)^{-1} (S(k, \varepsilon)P + \varepsilon^2 P - (A(k, \varepsilon) + \varepsilon^2 I)F(k)) \\ &\quad \times (S(k, \varepsilon)P + \varepsilon^2 I)^{-1} P. \end{aligned}$$

However, strictly speaking, the operator in the central brackets makes no sense. In order to obtain a precise version of (7.1), we need the following statement of abstract nature.

Lemma 7.1. *Let $\mathfrak{t}_1$ and $\mathfrak{t}_2$ be two closed and densely defined sesquilinear forms in a separable Hilbert space $\mathfrak{H}$. Suppose that both quadratic forms $\mathfrak{t}_1[u, u]$ and $\mathfrak{t}_2[u, u]$ are positive definite. Let T_1 and T_2 denote the selfadjoint operators in $\mathfrak{H}$ corresponding to the forms $\mathfrak{t}_1$ and $\mathfrak{t}_2$, respectively, and let $\mathcal{P}$ and $\mathcal{Q}$ be orthoprojections in $\mathfrak{H}$. Assume the following conditions:*

- 1° $\mathfrak{d} := \text{Dom } \mathfrak{t}_1 \subset \text{Dom } \mathfrak{t}_2$;
- 2° $\mathcal{P} \text{Dom } \mathfrak{t}_2 \subset \text{Dom } \mathfrak{t}_1$;
- 3° $\mathcal{Q}\mathfrak{d} \subset \mathfrak{d}$;
- 4° $\mathcal{P}$ commutes with T_2 ;
- 5° the following representation is true:

$$(7.2) \quad \mathfrak{t}_2[\mathcal{P}u, v] - \mathfrak{t}_1[\mathcal{Q}u, v] = (G_0u, Gv)_{\mathfrak{G}}, \quad u, v \in \mathfrak{d},$$

where $G_0, G : \mathfrak{H} \rightarrow \mathfrak{G}$ are densely defined linear operators acting from $\mathfrak{H}$ into some separable Hilbert space $\mathfrak{G}$. Moreover, $\mathfrak{d} \subset \text{Dom } G_0$, $\mathfrak{d} \subset \text{Dom } G$, and the operators

$$GT_1^{-1} : \mathfrak{H} \rightarrow \mathfrak{G}, \quad G_0T_2^{-1}\mathcal{P} : \mathfrak{H} \rightarrow \mathfrak{G}$$

are bounded.

Then we have

$$(7.3) \quad T_1^{-1}\mathcal{P} - \mathcal{Q}T_2^{-1}\mathcal{P} = (GT_1^{-1})^*(G_0T_2^{-1}\mathcal{P}).$$

Proof. In (7.2) we take $u = \mathcal{P}T_2^{-1}\varphi$ and $v = T_1^{-1}\psi$, where $\varphi, \psi \in \mathfrak{H}$. Then $u \in \mathfrak{d}$ by condition 2°, and obviously $v \in \mathfrak{d}$. Thus,

$$\mathfrak{t}_2[\mathcal{P}T_2^{-1}\varphi, T_1^{-1}\psi] - \mathfrak{t}_1[\mathcal{Q}\mathcal{P}T_2^{-1}\varphi, T_1^{-1}\psi] = (G_0\mathcal{P}T_2^{-1}\varphi, GT_1^{-1}\psi)_{\mathfrak{G}}, \quad \varphi, \psi \in \mathfrak{H}.$$

Using condition 4°, we rewrite this as

$$(\mathcal{P}\varphi, T_1^{-1}\psi)_{\mathfrak{H}} - (\mathcal{Q}T_2^{-1}\mathcal{P}\varphi, \psi)_{\mathfrak{H}} = (G_0T_2^{-1}\mathcal{P}\varphi, GT_1^{-1}\psi)_{\mathfrak{G}}, \quad \varphi, \psi \in \mathfrak{H},$$

which implies (7.3). $\square$

7.2. Let $|k| \leq t_0$, $\varepsilon > 0$. We are going to apply Lemma 7.1 with $\mathfrak{H} = L_2(\Omega)$, $\mathcal{P} = P$, $\mathcal{Q} = F(k)$,

$$\begin{aligned}\mathfrak{t}_1[u, v] &= a(k, \varepsilon)[u, v] + \varepsilon^2(u, v)_{L_2(\Omega)}, \\ \mathfrak{t}_2[u, v] &= s(k, \varepsilon)[Pu, Pv] + \varepsilon^2(u, v)_{L_2(\Omega)},\end{aligned}$$

$\mathfrak{d} = \tilde{H}^1(\Omega)$, and $\text{Dom } \mathfrak{t}_2 = \{u \in L_2(\Omega) : Pu \in \tilde{H}^1(0, a)\}$. Obviously, conditions 1°, 2°, and 4° are satisfied. Condition 3° is also satisfied by Lemma 6.3. We check condition 5°. The left-hand side of (7.2) takes the form

$$(7.4) \quad \begin{aligned}\mathfrak{A}[u, v] &:= s(k, \varepsilon)[Pu, Pv] + \varepsilon^2(Pu, v)_{L_2(\Omega)} - a(k, \varepsilon)[F(k)u, v] - \varepsilon^2(F(k)u, v)_{L_2(\Omega)} \\ &= \mathcal{I}_1[u, v] + \mathcal{I}_2[u, v] + \mathcal{I}_3[u, v],\end{aligned}$$

where, in accordance with (2.2), (3.1), and (4.1),

$$(7.5) \quad \mathcal{I}_1[u, v] := \varepsilon^2((P - F(k))u, v)_{L_2(\Omega)},$$

$$(7.6) \quad \mathcal{I}_2[u, v] := t^2(SP u, Pv)_{L_2(\Omega)} - a_1(k)[F(k)u, v],$$

$$(7.7) \quad \mathcal{I}_3[u, v] := (g_2^0 \varepsilon D_2 Pu, \varepsilon D_2 Pv)_{L_2(\Omega)} - (g_2 \varepsilon D_2 F(k)u, \varepsilon D_2 v)_{L_2(\Omega)}.$$

The forms $\mathcal{I}_1$ and $\mathcal{I}_2$. By (3.21), for the form (7.5) we have

$$(7.8) \quad \mathcal{I}_1[u, v] = -\varepsilon^2 t(\Phi(k)u, v)_{L_2(\Omega)} = (G_{01}u, G_1v)_{L_2(\Omega)},$$

where $G_{01} = -\varepsilon t^{1/2}\Phi(k)$, $G_1 = \varepsilon t^{1/2}I$. Using (3.25), (5.6), and (5.7), we obtain

$$(7.9) \quad \varepsilon \|G_{01}(S(k, \varepsilon)P + \varepsilon^2 I)^{-1}P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq C_1 \varepsilon^2 t^{1/2} c_*^{-1} (t^2 + \varepsilon^2)^{-1} \leq C_1 c_*^{-1} \varepsilon^{1/2},$$

$$(7.10) \quad \varepsilon \|G_1(A(k, \varepsilon) + \varepsilon^2 I)^{-1}P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq \varepsilon^2 t^{1/2} c_*^{-1} (t^2 + \varepsilon^2)^{-1} \leq c_*^{-1} \varepsilon^{1/2}.$$

As to the form (7.6), we apply (3.23):

$$(7.11) \quad \mathcal{I}_2[u, v] = -t^3(\Psi(k)u, v)_{L_2(\Omega)} = (G_{02}u, G_2v)_{L_2(\Omega)},$$

where $G_{02} = -t^{3/2}\Psi(k)$, $G_2 = t^{3/2}I$. Then, by (3.26) and (5.7),

$$(7.12) \quad \begin{aligned}\varepsilon \|G_{02}(S(k, \varepsilon)P + \varepsilon^2 I)^{-1}P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ \leq C_2 \varepsilon t^{3/2} c_*^{-1} (t^2 + \varepsilon^2)^{-1} \leq C_2 c_*^{-1} \varepsilon^{1/2}.\end{aligned}$$

From (5.6) we deduce the estimate

$$(7.13) \quad \varepsilon \|G_2(A(k, \varepsilon) + \varepsilon^2 I)^{-1}P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq \varepsilon t^{3/2} c_*^{-1} (t^2 + \varepsilon^2)^{-1} \leq c_*^{-1} \varepsilon^{1/2}.$$

The form $\mathcal{I}_3$. We rewrite the form (7.7) as

$$(7.14) \quad \mathcal{I}_3[u, v] = \mathcal{I}_3'[u, v] + \mathcal{I}_4[u, v],$$

where

$$(7.15) \quad \mathcal{I}_3'[u, v] = -(g_2 \varepsilon D_2(F(k) - P)u, \varepsilon D_2 v)_{L_2(\Omega)},$$

$$(7.16) \quad \mathcal{I}_4[u, v] = (g_2^0 \varepsilon D_2 Pu, \varepsilon D_2 Pv)_{L_2(\Omega)} - (g_2 \varepsilon D_2 Pu, \varepsilon D_2 v)_{L_2(\Omega)}.$$

For the form (7.15), we have

$$(7.17) \quad \mathcal{I}_3'[u, v] = (G_{03}u, G_3v)_{L_2(\Omega)},$$

where $G_{03} = -g_2 \varepsilon^{1/2} D_2(F(k) - P)$, $G_3 = \varepsilon^{3/2} D_2$. By (3.21), we can represent the operator G_{03} as

$$(7.18) \quad G_{03} = \tilde{G}_{03} + \hat{G}_{03} + \check{G}_{03},$$

where

$$\begin{aligned}\tilde{G}_{03} &= -g_2\varepsilon^{1/2}[D_2, F(k)]F(k), \\ \hat{G}_{03} &= -g_2\varepsilon^{1/2}[D_2, F(k)]F(k)^\perp, \\ \check{G}_{03} &= -g_2\varepsilon^{1/2}(F(k) - P)D_2 = -g_2\varepsilon^{1/2}t\Phi(k)D_2.\end{aligned}$$

From (1.1), (5.7), and (6.7) it follows that

$$(7.19) \quad \varepsilon\|\tilde{G}_{03}(S(k, \varepsilon)P + \varepsilon^2 I)^{-1}P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq \varepsilon^{3/2}c_1 C_7 t c_*^{-1}(t^2 + \varepsilon^2)^{-1} \leq c_1 C_7 c_*^{-1} \varepsilon^{1/2}.$$

Relations (1.1), (5.16), and (6.6) yield the estimate

$$(7.20) \quad \varepsilon\|\hat{G}_{03}(S(k, \varepsilon)P + \varepsilon^2 I)^{-1}P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \leq c_1 C_6 C_4 c_*^{-1/2} \varepsilon^{1/2}.$$

Finally, relations (1.1), (3.25), (5.7), and (5.11) imply

$$(7.21) \quad \begin{aligned}\varepsilon\|\check{G}_{03}(S(k, \varepsilon)P + \varepsilon^2 I)^{-1}P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ \leq \varepsilon^{1/2}c_1 C_1 t c_0^{-1/2} c_*^{-1/2}(t^2 + \varepsilon^2)^{-1/2} \leq c_1 C_1 c_0^{-1/2} c_*^{-1/2} \varepsilon^{1/2}.\end{aligned}$$

As a result, from (7.18)–(7.21) we obtain

$$(7.22) \quad \begin{aligned}\varepsilon\|G_{03}(S(k, \varepsilon)P + \varepsilon^2 I)^{-1}P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} &\leq C_8 \varepsilon^{1/2}, \\ C_8 &= c_1 c_*^{-1/2}(C_7 c_*^{-1/2} + C_4 C_6 + C_1 c_0^{-1/2}).\end{aligned}$$

Next, by (5.6) and (5.9),

$$(7.23) \quad \begin{aligned}\varepsilon\|G_3(A(k, \varepsilon) + \varepsilon^2 I)^{-1}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} &= \varepsilon^{3/2}\|\varepsilon D_2(A(k, \varepsilon) + \varepsilon^2 I)^{-1}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ &\leq \varepsilon^{3/2}c_0^{-1/2} c_*^{-1/2}(t^2 + \varepsilon^2)^{-1/2} \leq c_0^{-1/2} c_*^{-1/2} \varepsilon^{1/2}.\end{aligned}$$

The form (7.16) can be rewritten as

$$(7.24) \quad \begin{aligned}\mathcal{I}_4[u, v] &= (g_2^0 \varepsilon D_2 P u, \varepsilon D_2 P v)_{L_2(\Omega)} \\ &\quad - (g_2 \varepsilon D_2 P u, P \varepsilon D_2 v)_{L_2(\Omega)} - (g_2 \varepsilon D_2 P u, P^\perp \varepsilon D_2 v)_{L_2(\Omega)}.\end{aligned}$$

Since $P g_2 P = g_2^0$ (see (1.8)), the first two terms on the right in (7.24) compensate each other. Consequently,

$$(7.25) \quad \mathcal{I}_4[u, v] = -(g_2 \varepsilon D_2 P u, P^\perp \varepsilon D_2 v)_{L_2(\Omega)} = (G_{04} u, G_4 v)_{L_2(\Omega)},$$

where $G_{04} = -g_2 \varepsilon^{3/2} D_2 P$, $G_4 = \varepsilon^{1/2} P^\perp D_2$. Relations (1.1), (5.7), and (5.11) imply that

$$(7.26) \quad \begin{aligned}\varepsilon\|G_{04}(S(k, \varepsilon)P + \varepsilon^2 I)^{-1}P\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ \leq \varepsilon^{3/2}c_1 c_0^{-1/2} c_*^{-1/2}(t^2 + \varepsilon^2)^{-1/2} \leq c_1 c_0^{-1/2} c_*^{-1/2} \varepsilon^{1/2}.\end{aligned}$$

Next, we have

$$G_4(A(k, \varepsilon) + \varepsilon^2 I)^{-1} = \varepsilon^{1/2} P^\perp [D_2, (A(k, \varepsilon) + \varepsilon^2 I)^{-1}] + \varepsilon^{1/2} P^\perp (A(k, \varepsilon) + \varepsilon^2 I)^{-1} D_2.$$

The first term on the right is estimated with the help of (6.1), and the second with the help of (5.13) and (5.9). As a result, we obtain

$$(7.27) \quad \begin{aligned}\varepsilon\|G_4(A(k, \varepsilon) + \varepsilon^2 I)^{-1}\|_{L_2(\Omega) \rightarrow L_2(\Omega)} \\ \leq \varepsilon^{3/2} C_5 (t^2 + \varepsilon^2)^{-1/2} + \varepsilon^{1/2} c_0^{-1/2} C_3 \leq (C_5 + c_0^{-1/2} C_3) \varepsilon^{1/2}.\end{aligned}$$

7.3. We summarize the study of the form (7.4). By (7.8), (7.11), (7.14), (7.17), and (7.25), this form can be represented as

$$\begin{aligned}\mathfrak{A}[u, v] &= (G_{01}u, G_1v)_{L_2(\Omega)} \\ &\quad + (G_{02}u, G_2v)_{L_2(\Omega)} + (G_{03}u, G_3v)_{L_2(\Omega)} + (G_{04}u, G_4v)_{L_2(\Omega)}.\end{aligned}$$

We put $\mathfrak{G} = L_2(\Omega; \mathbb{C}^4) = L_2(\Omega) \times L_2(\Omega) \times L_2(\Omega) \times L_2(\Omega)$ and introduce the operators

$$\begin{aligned}G_0 : L_2(\Omega) &\rightarrow \mathfrak{G}, & G_0 &= \text{col}\{G_{01}, G_{02}, G_{03}, G_{04}\}, \\ G : L_2(\Omega) &\rightarrow \mathfrak{G}, & G &= \text{col}\{G_1, G_2, G_3, G_4\}, \\ \text{Dom } G_0 &= \text{Dom } G = \{u \in L_2(\Omega) : D_2u \in L_2(\Omega)\}.\end{aligned}$$

The operators G_0 and G are densely defined, and we have $\mathfrak{d} = \tilde{H}^1(\Omega) \subset \text{Dom } G_0 = \text{Dom } G$ and

$$\mathfrak{A}[u, v] = (G_0u, Gv)_{\mathfrak{G}}, \quad u, v \in \mathfrak{d}.$$

Thus, we have checked condition (7.2) in the case in question. Relations (7.10), (7.13), (7.23), and (7.27) show that the operator $G(A(k, \varepsilon) + \varepsilon^2 I)^{-1} : L_2(\Omega) \rightarrow \mathfrak{G}$ is bounded, and

$$(7.28) \quad \varepsilon \|G(A(k, \varepsilon) + \varepsilon^2 I)^{-1}\|_{L_2(\Omega) \rightarrow \mathfrak{G}} \leq C_9 \varepsilon^{1/2},$$

where $C_9^2 = 2c_*^{-2} + c_0^{-1}c_*^{-1} + (C_5 + c_0^{-1/2}C_3)^2$. Inequalities (7.9), (7.12), (7.22), and (7.26) imply that the operator $G_0(S(k, \varepsilon)P + \varepsilon^2 I)^{-1}P : L_2(\Omega) \rightarrow \mathfrak{G}$ is bounded, and

$$(7.29) \quad \varepsilon \|G_0(S(k, \varepsilon)P + \varepsilon^2 I)^{-1}P\|_{L_2(\Omega) \rightarrow \mathfrak{G}} \leq C_{10} \varepsilon^{1/2},$$

where $C_{10}^2 = (C_1^2 + C_2^2)c_*^{-2} + C_8^2 + c_1^2c_0^{-1}c_*^{-1}$. Observe that the constants C_9 and C_{10} can be expressed in terms of c_0, c_1, c_2 .

Thus, we see that condition 5° of Lemma 7.1 is satisfied. Applying Lemma 7.1, we obtain the representation

$$(7.30) \quad \begin{aligned} &(A(k, \varepsilon) + \varepsilon^2 I)^{-1}P - F(k)(S(k, \varepsilon)P + \varepsilon^2 I)^{-1}P \\ &= (G(A(k, \varepsilon) + \varepsilon^2 I)^{-1})^*(G_0(S(k, \varepsilon)P + \varepsilon^2 I)^{-1}P). \end{aligned}$$

Relations (7.28)–(7.30) imply estimate (5.18) with $\hat{C} = C_9 C_{10}$. *This completes the proof of Proposition 5.3 and, with it, of Theorem 4.1. Also, this completes (see Subsection 4.2) the proof of Theorem 2.1, and, thereby (see Subsection 2.3), that of Theorem 1.1.*

§8. ON GENERALIZATIONS OF THEOREM 1.1

In this section, we briefly discuss possible generalizations of the main result of the paper.

8.1. The two-dimensional operator of the form (0.2) in $L_2(\Pi)$ with a nondiagonal matrix $g(\mathbf{x})$ can also be studied. The statement of Theorem 1.1 carries over to the case where the *antidiagonal entries of the matrix $g(\mathbf{x})$ do not depend on x_1* , i.e., $g(\mathbf{x})$ is of the form

$$g(\mathbf{x}) = \begin{pmatrix} g_1(x_1, x_2) & \hat{g}(x_2) \\ \hat{g}(x_2) & g_2(x_1, x_2) \end{pmatrix}.$$

We suppose that the matrix $g(\mathbf{x})$ is positive definite and bounded, and that the coefficients g_1, g_2 are subject to the same conditions as before. Also, we assume that the coefficient $\hat{g}(x_2)$ is real-valued, periodic (with respect to x_2), and Lipschitz. In this case the effective operator coincides with $A^0 = -\text{div } g^0 \text{ grad}$, where

$$g^0 = \begin{pmatrix} g_1^0(x_2) & \hat{g}(x_2) \\ \hat{g}(x_2) & g_2^0(x_2) \end{pmatrix},$$

and the $g_j^0(x_2)$, $j = 1, 2$, are as in (1.7) and (1.8). *Under the above conditions, an analog of estimate (1.11) is valid.*

In the case where $g(\mathbf{x})$ is a nondiagonal matrix of general type, additional technical difficulties arise, and the author has not yet succeeded in overcoming them.

8.2. An analog of the operator (0.2) can also be studied in the domain $\Xi = \mathbb{R}^{d_1} \times (0, a)^{d_2}$ of an arbitrary dimension $d = d_1 + d_2$. We use the notation $\mathbf{x} = (\mathbf{x}', \mathbf{x}'') \in \Xi$, $\mathbf{x}' \in \mathbb{R}^{d_1}$, $\mathbf{x}'' \in (0, a)^{d_2}$, $\mathbf{D}_{\mathbf{x}'} = -i \operatorname{grad}_{\mathbf{x}'}$, $\mathbf{D}_{\mathbf{x}'}^* = -i \operatorname{div}_{\mathbf{x}'}$, and similarly for $\mathbf{D}_{\mathbf{x}''}$, $\mathbf{D}_{\mathbf{x}''}^*$. In $L_2(\Xi)$, we consider the operator A corresponding to the differential expression

$$(8.1) \quad A = \mathbf{D}_{\mathbf{x}'}^* g_1(\mathbf{x}', \mathbf{x}'') \mathbf{D}_{\mathbf{x}'} + \mathbf{D}_{\mathbf{x}''}^* g_2(\mathbf{x}', \mathbf{x}'') \mathbf{D}_{\mathbf{x}''},$$

with periodic boundary conditions on $\partial\Xi$. Here $g_1(\mathbf{x})$ is a $(d_1 \times d_1)$ -matrix-valued function, and $g_2(\mathbf{x})$ is a $(d_2 \times d_2)$ -matrix-valued function in Ξ . Both matrices are real-valued, measurable, positive definite, and periodic in $\mathbf{x}'$ with respect to a lattice $\Gamma \subset \mathbb{R}^{d_1}$. We assume that

$$g_j, g_j^{-1} \in L_\infty(\Xi), \quad j = 1, 2.$$

Moreover, we assume that g_1, g_2 belong to the Lipschitz class with respect to $\mathbf{x}''$,

$$\operatorname{ess\,sup}_{\mathbf{x} \in \Xi} |\mathbf{D}_{\mathbf{x}''} g_j(\mathbf{x})| \leq c, \quad j = 1, 2,$$

and satisfy periodic boundary conditions in $\mathbf{x}''$ (the values of g_1, g_2 on the opposite faces of Ξ coincide). We consider the operator

$$(8.2) \quad A_\varepsilon = \mathbf{D}_{\mathbf{x}'}^* g_1(\mathbf{x}'/\varepsilon, \mathbf{x}'') \mathbf{D}_{\mathbf{x}'} + \mathbf{D}_{\mathbf{x}''}^* g_2(\mathbf{x}'/\varepsilon, \mathbf{x}'') \mathbf{D}_{\mathbf{x}''}, \quad \varepsilon > 0.$$

Theorem 1.1 can be carried over to the operator (8.2). The effective operator has the form

$$A^0 = \mathbf{D}_{\mathbf{x}'}^* g_1^0(\mathbf{x}'') \mathbf{D}_{\mathbf{x}'} + \mathbf{D}_{\mathbf{x}''}^* g_2^0(\mathbf{x}'') \mathbf{D}_{\mathbf{x}''}.$$

We present the rule of calculation of the coefficients. Let $\Omega' \subset \mathbb{R}^{d_1}$ be the cell of the lattice Γ . Then $g_2^0(\mathbf{x}'')$ is equal to the mean value of g_2 :

$$(8.3) \quad g_2^0(\mathbf{x}'') = |\Omega'|^{-1} \int_{\Omega'} g_2(\mathbf{x}', \mathbf{x}'') d\mathbf{x}'.$$

The matrix $g_1(\mathbf{x})$ is homogenized by the rule corresponding to the d_1 -dimensional elliptic operator $\mathbf{D}_{\mathbf{x}'}^* g_1 \mathbf{D}_{\mathbf{x}'}$. This rule can be described as follows. For a given $\mathbf{C} \in \mathbb{C}^{d_1}$, let $v(\mathbf{x})$ be the periodic (in $\mathbf{x}'$) solution of the equation $\mathbf{D}_{\mathbf{x}'}^* g_1(\mathbf{x}', \mathbf{x}'') (\mathbf{D}_{\mathbf{x}'} v + \mathbf{C}) = 0$. (Here $\mathbf{x}''$ is a parameter.) Then

$$(8.4) \quad g_1^0(\mathbf{x}'') \mathbf{C} = |\Omega'|^{-1} \int_{\Omega'} g_1(\mathbf{x}', \mathbf{x}'') (\mathbf{D}_{\mathbf{x}'} v + \mathbf{C}) d\mathbf{x}', \quad \mathbf{C} \in \mathbb{C}^{d_1}.$$

Note that, for $d_1 = 1$, the rule (8.4) corresponds to calculating the harmonic average value of g_1 , and then (8.4) turns into (1.7). The following analog of Theorem 1.1 is true.

Theorem 8.1. *Under the above conditions, we have*

$$(8.5) \quad \|(A_\varepsilon + I)^{-1} - (A^0 + I)^{-1}\|_{L_2(\Xi) \rightarrow L_2(\Xi)} \leq C\varepsilon, \quad 0 < \varepsilon \leq 1.$$

8.3. We can include “antidiagonal” terms in the operator (8.1), but only with coefficients independent of $\mathbf{x}'$. In this case, $A = -\operatorname{div} g(\mathbf{x}) \operatorname{grad}$, and the matrix $g(\mathbf{x})$ is of the form

$$g(\mathbf{x}) = \begin{pmatrix} g_1(\mathbf{x}', \mathbf{x}'') & \widehat{g}(\mathbf{x}'') \\ (\widehat{g}(\mathbf{x}''))^* & g_2(\mathbf{x}', \mathbf{x}'') \end{pmatrix},$$

where g_1, g_2 are as above, the $(d \times d)$ -matrix-valued function $g(\mathbf{x})$ is bounded and positive definite, and the $(d_1 \times d_2)$ -matrix-valued function $\widehat{g}(\mathbf{x}'')$ is periodic (in $\mathbf{x}''$) and belongs to the Lipschitz class. In this case, the effective operator is of the form $A^0 = -\operatorname{div} g^0(\mathbf{x}'') \operatorname{grad}$, where

$$g^0(\mathbf{x}'') = \begin{pmatrix} g_1^0(\mathbf{x}'') & \widehat{g}(\mathbf{x}'') \\ (\widehat{g}(\mathbf{x}''))^* & g_2^0(\mathbf{x}'') \end{pmatrix}.$$

Here g_1^0 and g_2^0 are defined by the same rules (8.4) and (8.3), respectively. *Under the above conditions, an analog of estimate (8.5) is valid.*

REFERENCES

- [1] A. Bensoussan, J. L. Lions, and G. Papanicolaou, *Asymptotic analysis for periodic structures*, Stud. Math. Appl., vol. 5, North-Holland Publishing Co., Amsterdam–New York, 1978. MR0503330 (82h:35001)
- [2] N. S. Bakhvalov and G. P. Panasenko, *Homogenization: averaging processes in periodic media. Mathematical problems in mechanics of composite materials*, “Nauka,” Moscow, 1984; English transl., Math. Appl. (Soviet Ser.), vol. 36, Kluwer Acad. Publ. Group, Dordrecht, 1989. MR1112788 (92d:73002)
- [3] V. V. Zhikov, S. M. Kozlov, and O. A. Oleĭnik, *Homogenization of differential operators*, Fizmatgiz, Moscow, 1993; English transl., Springer-Verlag, Berlin, 1994. MR1329546 (96h:35003b)
- [4] V. V. Zhikov, *Spectral approach to asymptotic diffusion problems*, Differentsial’nye Uravneniya **25** (1989), no. 1, 44–50; English transl., Differential Equations **25** (1989), no. 1, 33–39. MR0986395 (90a:35107)
- [5] C. Conca and M. Vanninathan, *Homogenization of periodic structures via Bloch decomposition*, SIAM J. Appl. Math. **57** (1997), no. 6, 1639–1659. MR1484944 (98j:35017)
- [6] E. V. Sevost’yanova, *Asymptotic expansion of the solution of a second-order elliptic equation with periodic rapidly oscillating coefficients*, Mat. Sb. (N.S.) **115** (1981), no. 2, 204–222; English transl., Math. USSR-Sb. **43** (1982), no. 2, 181–198. MR0622145 (83d:35038)
- [7] M. Birman and T. Suslina, *Threshold effects near the lower edge of the spectrum for periodic differential operators of mathematical physics*, Systems, Approximation, Singular Integral Operators, and Related Topics (Bordeaux, 2000), Oper. Theory Adv. Appl., vol. 129, Birkhäuser, Basel, 2001, pp. 71–107. MR1882692 (2003f:35220)
- [8] ———, *Second order periodic differential operators. Threshold properties and homogenization*, Algebra i Analiz **15** (2003), no. 5, 1–108; English transl., St. Petersburg Math. J. **15** (2004), no. 5, 1–77. MR2068790
- [9] O. A. Ladyzhenskaya and N. N. Ural’tseva, *Linear and quasilinear equations of elliptic type*, 2nd ed., “Nauka”, Moscow, 1973; English transl. of 1st ed., Acad. Press, New York–London, 1968. MR0509265 (58:23009)
- [10] D. Gilbarg and N. S. Trudinger, *Elliptic partial differential equations of second order*, Springer-Verlag, Berlin, 1998. MR0737190 (86c:35035)

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