A TANGENTIAL CONVERGENCE FOR BOUNDED HARMONIC FUNCTIONS ON A RANK ONE SYMMETRIC SPACE

BY

JACEK CYGAN

ABSTRACT. Let u be a bounded harmonic function on a noncompact rank one symmetric space $M = G/K \approx N^-A$, N^-AK being a fixed Iwasawa decomposition of G . We prove that if for an $a_0 \in A$ there exists a limit $\lim u(na_0) \equiv c_0$, as $n \in N^-$ goes to infinity, then for any $a \in A$, $\lim u(na) = c_0$. For $M = SU(n, 1)/S(U(n) \times U(1)) = B^n$, the unit ball in $\mathbb{C}^n$ with the Bergman metric, this is a result of Hulanicki and Ricci, and in this case it reads (via the Cayley transformation) as a theorem on convergence of a bounded harmonic function to a boundary value at a fixed boundary point, along appropriate, tangent to ∂B^n , surfaces.

0. Introduction. Let M be a noncompact symmetric space of rank one. M can be expressed as a homogeneous space G/K where G is a semisimple group of isometries of M and K is a maximal compact subgroup of G . Let $\mathfrak{g}, \mathfrak{k}$ denote the Lie algebras of G and K , B the Killing form of $\mathfrak{g}$, and $\mathfrak{p}$ the orthogonal complement of $\mathfrak{k}$ in $\mathfrak{g}$ relative to B . If $\pi: G \rightarrow G/K$ denotes the canonical projection, its differential at the identity, π_* , identifies the subspace $\mathfrak{p}$ of $\mathfrak{g}$ with $T_0(M)$, the tangent space of M at the origin $o = \pi(e)$, and the invariant metric g on M can be chosen so that g_o corresponds to the restriction of B to $\mathfrak{p} \times \mathfrak{p}$ under the above identification. We denote by Δ the corresponding (G -invariant) Laplace-Beltrami operator on M . A function $u \in C^\infty(M)$ is called *harmonic* if $\Delta u = 0$. Let $\mathfrak{a}$ be a maximal (one-dimensional) abelian subspace of $\mathfrak{p}$, α and possibly 2α in $\mathfrak{a}^*$, the corresponding system of positive restricted roots relative to the fixed choice of a "positive part" $\mathfrak{a}^+$ in $\mathfrak{a}$. Let $\mathfrak{g}_{-\alpha}$ and $\mathfrak{g}_{-2\alpha}$ denote the root spaces corresponding to $-\alpha$ and -2α . Then $\mathfrak{n}^- = \mathfrak{g}_{-\alpha} \oplus \mathfrak{g}_{-2\alpha}$ is a nilpotent subalgebra of $\mathfrak{g}$ and one has the Iwasawa decomposition $G = N^-AK$, with $N^- = \exp \mathfrak{n}^-$, $A = \exp \mathfrak{a}$. The above decomposition shows that every $p \in M$ can be uniquely written as $p = na \cdot o$ ($n \in N^-$, $a \in A$). We regard the nilpotent group N^- as a boundary for the symmetric space M in the following sense. The bounded harmonic functions u on M have boundary values on N^- , i.e. $\lim_{\log a \rightarrow \infty} u(na \cdot o) \equiv \varphi(n)$ exists a.e. (relative to the Haar measure on N^-) and $\varphi \in L^\infty(N^-)$. $\log a \rightarrow \infty$ is understood with respect to the ordering induced on $\mathfrak{a}$ by $\mathfrak{a}^+$. Moreover,

$$u(na \cdot o) = \varphi * P_a(n) = \int_{N^-} \varphi(n_1) P_a(n_1^{-1}n) \, dn_1.$$

Received by the editors April 28, 1980.

AMS (MOS) subject classifications (1970). Primary 43A85, 53C35; Secondary 43A20, 43A90.

© 1981 American Mathematical Society
0002-9947/81/0000-0255/\$04.50

The function $P_a(n)$ on $N^- \times A$ is called the Poisson kernel for the symmetric space M and is given by (Helgason [4])

$$P_a(n) = c\epsilon^{d/2} \left[\left(\epsilon + \frac{1}{2} Q(X_{-\alpha}) \right)^2 + 2Q(X_{-2\alpha}) \right]^{-d/2},$$

where

$$\begin{aligned} n &= \exp(X_{-\alpha} + X_{-2\alpha}), \quad X_{-\alpha} \in \mathfrak{g}_{-\alpha}, \quad X_{-2\alpha} \in \mathfrak{g}_{-2\alpha}; \\ \epsilon &= e^{-\alpha(\log a)}; \quad Q(X) = (X, X)_\theta / 2(m_\alpha + 4m_{2\alpha}) \end{aligned}$$

with $(X, X)_\theta = -B(X, \theta X)$ for $X \in \mathfrak{g}$, θ denoting the Cartan involution associated with the pair $(\mathfrak{g}, \mathfrak{k})$; $m_\alpha = \dim \mathfrak{g}_{-\alpha}$, $m_{2\alpha} = \dim \mathfrak{g}_{-2\alpha}$, $d = m_\alpha + 2m_{2\alpha}$. The constant c is such that the integral of P_a over N^- is equal to 1.

The following theorem on "tangential" convergence for bounded harmonic functions on the Siegel domain

$$D_{r-1} = \left\{ (z_1, \dots, z_r) \in \mathbb{C}^r : \operatorname{Im} z_r > \sum_{j=1}^{r-1} |z_j|^2 \right\},$$

$r \geq 2$, (or, equivalently, on $M = SU(r, 1)/S(U(r) \times U(1))$ —the complex hyperbolic space) has been obtained by Hulanicki and Ricci [5]. We formulate it below in terms of a homogeneous space M .

THEOREM. *Let u be a bounded harmonic function on a noncompact rank one symmetric space M . In the notation above, assume that for an $a_0 \in A$, $\lim_{N^- \ni n \rightarrow \infty} u(na_0 \cdot o) = c_0$. Then for any $a \in A$, $\lim_{N^- \ni n \rightarrow \infty} u(na \cdot o) = c_0$.*

Our aim here is to prove the above Theorem and the proof is based on the classification of symmetric spaces. That is, we discuss separately the cases of M being the real, complex (to see how the $M = D_{r-1}$ case fits to our scheme), quaternion and octonion hyperbolic space, which corresponds respectively to G being the classical group $SO_0(r, 1)$, $SU(r, 1)$, $Sp(r, 1)$ and the exceptional group $F_{4(-20)}$. Following the Hulanicki-Ricci method, for each case we construct a suitable commutative subalgebra $\mathcal{Q}$ of (multi) radial functions in $L^1(N^-)$, to which the Poisson kernel P_a belongs. We describe the set $\mathfrak{M}(\mathcal{Q})$ of the maximal ideals in $\mathcal{Q}$ and check that the Gel'fand transform $\hat{P}_a$ of P_a never vanishes on $\mathfrak{M}(\mathcal{Q})$. The Theorem may then be stated as a theorem on certain ideals in $L^1(N^-)$ and is a consequence of the Wiener property of the algebra $\mathcal{Q}$. To study the algebra $\mathcal{Q}$ we use the holomorphically induced (realizations of the irreducible unitary) representations of N^- .

1. Nilpotent group N^- . Let $\mathbf{F}$ denote the field $\mathbf{R}$, $\mathbf{C}$, $\mathbf{H}$ or the Cayley numbers $\mathbf{O}$ (octonions); $\mathbf{F}_0 = \{q \in \mathbf{F} : q + \bar{q} = 0\}$, $\bar{\cdot}$ being the usual conjugation in $\mathbf{F} = \mathbf{C}, \mathbf{H}, \mathbf{O}$ and $\bar{q} = q$ for $q \in \mathbf{R}$; $\operatorname{Im} q = \frac{1}{2}(q - \bar{q})$, $\sigma = 2s = \dim_{\mathbf{R}} \mathbf{F}$. According to the notation of the previous section, $\mathfrak{g} = \mathfrak{k} + \mathfrak{p}$ and for the classical G we have (cf., e.g., [3, pp. 348–351])

$$\begin{aligned} \mathfrak{k} &= \left\{ \begin{pmatrix} Z & 0 \\ 0 & p \end{pmatrix} : \begin{array}{l} Z \text{ an } r \times r \text{ skew-Hermitian matrix over } \mathbf{F}, \\ p \in \mathbf{F}_0, \operatorname{tr} Z = -p \text{ in case of } \mathbf{F} = \mathbf{C} \end{array} \right\}, \\ \mathfrak{p} &= \left\{ \begin{pmatrix} 0 & {}'q \\ \bar{q} & 0 \end{pmatrix} : q \in \mathbf{F}^r = \mathbf{F} \times \cdots \times \mathbf{F} \right\}, \\ \alpha &= \{tE_{1,r+1} + tE_{r+1,1} : t \in \mathbf{R}\}, \end{aligned}$$

where E_{kl} denotes the $(r+1) \times (r+1)$ matrix $(\delta_{ak}\delta_{bl})_{1 \leq a,b \leq r+1}$, $r \geq 2$. We choose a basis $H = E_{1,r+1} + E_{r+1,1}$ in α and fix an ordering so that $H \in \alpha^+$. Then $\alpha \in \alpha^*$ such that $\alpha(H) = 1$ is a positive restricted root, and we have

$$\begin{aligned} \mathfrak{g}_{-\alpha} &= \left\{ \begin{bmatrix} 0 & -\bar{q} & 0 \\ {}'q & 0 & {}'q \\ 0 & \bar{q} & 0 \end{bmatrix} : q = (q_2, \dots, q_r) \in \mathbf{F}^{r-1} \right\}, \\ \mathfrak{g}_{-2\alpha} &= \left\{ \begin{bmatrix} p & 0 & p \\ 0 & 0 & 0 \\ \bar{p} & 0 & \bar{p} \end{bmatrix} : p \in \mathbf{F}_0 \right\} \quad (= \{0\} \text{ for } \mathbf{F} = \mathbf{R}). \end{aligned}$$

We shall identify $\mathfrak{n}^- = \mathfrak{g}_{-\alpha} \oplus \mathfrak{g}_{-2\alpha}$ with $\mathbf{F}^{r-1} \times \mathbf{F}_0$ by the correspondence

$$\begin{bmatrix} p & -\bar{q} & p \\ {}'q & 0 & {}'q \\ \bar{p} & \bar{q} & \bar{p} \end{bmatrix} \leftrightarrow (q, -p).$$

In these coordinates on $\mathfrak{n}^-$ the commutator of $(q, p) = (q_1, \dots, q_{r-1}, p)$ and $(q', p') = (q'_1, \dots, q'_{r-1}, p')$ in $\mathbf{F}^{r-1} \times \mathbf{F}_0$ is given by

$$[(q, p), (q', p')] = (0, 2 \operatorname{Im}(\bar{q} \cdot q')), \quad (1)$$

where we have put $\bar{q} \cdot q'$ for $\sum_{i=1}^{r-1} \bar{q}_i q'_i$. We also have the formula (cf., e.g., [11, p. 39])

$$((q, p), (q', p'))_\theta = 4(m_\alpha + 4m_{2\alpha}) \operatorname{Re}(\bar{q} \cdot q' + \bar{p}p'). \quad (2)$$

For the exceptional G (cf., e.g., [10, pp. 522–530]), $\mathfrak{g} = \mathfrak{f}_{4(-20)}$ is isomorphic to the Lie algebra $\operatorname{Der}(\mathcal{J})$ of derivations of the Jordan algebra $(\mathcal{J}, \circ)$ of 3×3 octonion matrices A of the form

$$A = \begin{bmatrix} \alpha_1 & a_3 & a_2 \\ \bar{a}_3 & \alpha_2 & a_1 \\ -\bar{a}_2 & -\bar{a}_1 & \alpha_3 \end{bmatrix}, \quad a_i \in \mathbf{O}, \quad \alpha_i \in \mathbf{R}, \quad i = 1, 2, 3,$$

with multiplication given by $A \circ B = \frac{1}{2}(AB + BA)$, $A, B \in \mathcal{J}$, AB denoting the usual matrix multiplication. We have

$$\begin{aligned} \mathfrak{k} &= \{D \in \operatorname{Der}(\mathcal{J}) : D(E_{33}) = 0\}, \\ \mathfrak{p} &= \left\{ D_Q \in \operatorname{Der}(\mathcal{J}) : Q = \begin{pmatrix} 0 & {}'q \\ \bar{q} & 0 \end{pmatrix}, \quad q \in \mathbf{O}^2 \right\}, \end{aligned}$$

where $D_Q(B) = QB - BQ$, $B \in \mathcal{J}$.

$$\alpha = \{D_Q \in \text{Der}(\mathcal{G}): Q = tE_{13} + tE_{31}, \quad t \in \mathbf{R}\}.$$

We choose $H = D_Q$ with $Q = E_{13} + E_{31} \in \alpha^+$ and $\alpha \in \alpha^*$ such that $\alpha(H) = 1$. Then

$$\mathfrak{g}_{-\alpha} = \left\{ D_{Q(q)}: Q(q) = \begin{bmatrix} 0 & -\bar{q} & 0 \\ q & 0 & q \\ 0 & \bar{q} & 0 \end{bmatrix}, \quad q \in \mathbf{O} \right\},$$

$$\mathfrak{g}_{-2\alpha} = \left\{ D_{Q(p)}: Q(p) = \begin{bmatrix} p & 0 & p \\ 0 & 0 & 0 \\ \bar{p} & 0 & \bar{p} \end{bmatrix}, \quad p \in \mathbf{O}_0 \right\}.$$

Identifying $D_{Q(q)} + D_{Q(p)}$ in $\mathfrak{g}_{-\alpha} \oplus \mathfrak{g}_{-2\alpha}$ with $(q, -p)$ in $\mathbf{O} \times \mathbf{O}_0$, we obtain the same formulas for the commutator and the inner product of (q, p) and (q', p') in $\mathbf{O} \times \mathbf{O}_0$ as those given by (1) and (2) above.

Writing N^- as the manifold n^- with the group multiplication given by the Campbell-Hausdorff formula we obtain

PROPOSITION 1. *The underlying manifold for the nilpotent group N^- is $\mathbf{F}^{r-1} \times \mathbf{F}_0$ with $r \geq 2$ for $\mathbf{F} = \mathbf{R}, \mathbf{C}, \mathbf{H}$ ($\mathbf{F}_0 = \{0\}$ if $\mathbf{F} = \mathbf{R}$) and with $r = 2$ for $\mathbf{F} = \mathbf{O}$. The group law is*

$$(q, p)(q', p') = (q + q', p + p' + \text{Im}(\bar{q} \cdot q')).$$

The Haar measure on N^- is the ordinary Lebesgue measure on $\mathbf{R}^k \approx \mathbf{F}^{r-1} \times \mathbf{F}_0$, $k = r\sigma - 1$. We normalize it so that the volume of the unit cube in $\mathbf{R}^k$ is 1 and denote by $dqdp$. The Poisson kernel is given by

$$P_{\exp(tH)}(q, p) = c_{r, \mathbf{F}} \varepsilon^{d/2} [(|q|^2 + \varepsilon)^2 + 4|p|^2]^{-d/2},$$

where $\varepsilon = e^{-t}$, $d = (r+1)\sigma - 2$, $|q|^2 = \bar{q} \cdot q$, $c_{r, \mathbf{F}} = 2^{d-1} \pi^{-rs} \Gamma(rs)$ with $\sigma = 2s = \dim_{\mathbf{R}} \mathbf{F}$.

2. Holomorphically induced representations of N^- . The adjoint and coadjoint action of N^- on n^- and n^{-*} , respectively, is given by

$$\text{Ad}_{(q,p)}(q'', p'') = (q'', p'' + 2 \text{Im}(\bar{q} \cdot q'')),$$

$$\text{Ad}_{(q,p)}^*(q', p') = (q' + 2q\bar{p}', p'),$$

$(q, p) \in N^-$, $(q'', p'') \in n^-$, $(q', p') \in n^{-*}$, $q\bar{p}' = (q_1 \bar{p}', \dots, q_{r-1} \bar{p}')$, and we have identified n^{-*} , the dual space of n^- , with n^- by $\langle \cdot, \cdot \rangle = (\cdot, \cdot)_{\theta}/4(m_{\alpha} + 4m_{2\alpha})$. The single points $(q', 0) \in n^{-*}$ are 0-dimensional orbits of Ad^* on n^{-*} and the corresponding (1-dimensional) representations of N^- are given by the characters

$$\chi_{(q',0)}(q, p) = \exp(\sqrt{-1} \text{Re}(\bar{q}' \cdot q)), \quad (q, p) \in N^-.$$

The remaining (maximal dimensional) orbits, for $\mathbf{F} = \mathbf{C}, \mathbf{H}$ and $\mathbf{O}$, are of the form $\mathbf{F}^{r-1} \times \{p''\}$, $p'' \neq 0$, so they are parameterized, e.g., by the functionals $f = (0, p'') \in n^{-*}$ with $p'' \in F_0 \setminus \{0\}$. For such f and (q, p) , $(q', p') \in n^-$ we have

$$\langle f, [(q, p), (q', p')] \rangle = 2 \text{Re}((qp'')^{-} \cdot q') = -2 \text{Re}(\bar{q} \cdot (q'p'')),$$

i.e. the operator $R_p: q \rightarrow qp''$ is skew-symmetric on F^{r-1} with respect to the $\mathbf{R}$ -bilinear symmetric form $\langle \cdot, \cdot \rangle$ on F^{r-1} given by $\langle q, q' \rangle = 2 \operatorname{Re}(\bar{q} \cdot q') =$ the usual inner product on $\mathbf{R}^k$, $k = (r-1)\sigma$. Let $1, i$ denote the usual basis of $\mathbf{C}$ (over $\mathbf{R}$); $1, i, j, k$ the basis of $\mathbf{H}$, and $1, i, j, k, e, ie, je, ke$ the basis of $\mathbf{O} (= \mathbf{H} + \mathbf{H}e$ with the multiplication defined by $(ae)b = (a\bar{b})e$, $a(be) = (ba)e$, $(ae)(be) = -\bar{b}a$ for $a, b \in \mathbf{H}$). Suppose now that $f = (0, i\lambda) \in n^{-*}$ with λ positive real. In the above bases of $\mathbf{F}$ the matrix of $R_{i\lambda}$ acting on F^{r-1} with $r = 2$ is equal to $\lambda(E_{21} - E_{12})$ for $\mathbf{F} = \mathbf{C}$, to $\lambda(E_{21} - E_{12}) - \lambda(E_{43} - E_{34})$ for $\mathbf{F} = \mathbf{H}$ and to $\lambda(E_{21} - E_{12}) - \lambda(E_{43} - E_{34}) - \lambda(E_{65} - E_{56}) + \lambda(E_{87} - E_{78})$ for $\mathbf{F} = \mathbf{O}$. Put

$$\begin{aligned} e_1 &= \frac{1}{2}(1 + \sqrt{-1} i), & e_2 &= \frac{1}{2}(\sqrt{-1} j + k), \\ e_3 &= \frac{1}{2}(\sqrt{-1} e + ie), & e_4 &= \frac{1}{2}(je + \sqrt{-1} ke) \end{aligned}$$

for elements in $\mathbf{F}^{\mathbf{C}} = \mathbf{F} + \sqrt{-1} \mathbf{F}$ —the complexification of $\mathbf{F}$. Now define a subspace W of $\mathbf{F}^{\mathbf{C}}$ by

$$W = \begin{cases} \mathbf{C}e_1 & \text{if } \mathbf{F} = \mathbf{C}, \\ \mathbf{C}e_1 + \mathbf{C}e_2 & \text{if } \mathbf{F} = \mathbf{H}, \\ \mathbf{C}e_1 + \mathbf{C}e_2 + \mathbf{C}e_3 + \mathbf{C}e_4 & \text{if } \mathbf{F} = \mathbf{O}. \end{cases}$$

Thus

$$\mathfrak{h} = W^{r-1} \times \mathbf{F}_0^{\mathbf{C}} \quad (2.1)$$

is a positive polarization at $f = (0, i\lambda)$ such that

$$\mathfrak{h} + \bar{\mathfrak{h}} = n^{-\mathbf{C}}, \quad \mathfrak{h} \cap \bar{\mathfrak{h}} = \{0\} \times \mathbf{F}_0^{\mathbf{C}}, \quad \mathfrak{h} / \{0\} \times \mathbf{F}_0^{\mathbf{C}} = W, \quad (2.2)$$

where $\bar{z} = x - \sqrt{-1} y$ for $z = x + \sqrt{-1} y \in n^- + \sqrt{-1} n^- = n^{-\mathbf{C}}$. For arbitrary $f = (0, p'')$ with $p'' \in \mathbf{F}_0 \times \{0\}$, there exists an orthogonal transformation Ω on $\mathbf{R}^\sigma \approx \mathbf{F}$, such that

$$\langle R_{p''} q, q' \rangle = \langle R_{i|p''|} \Omega q, \Omega q' \rangle, \quad q, q' \in F^{r-1}.$$

Hence,

$$\langle f, [(q, p), (q', p')] \rangle = \langle (0, i|p''|), [(\Omega q, p), (\Omega q', p')] \rangle,$$

and $\mathfrak{h}' = \Omega \mathfrak{h}$, with $\mathfrak{h}$ as in (2.1), is a positive polarization at f and $\mathfrak{h}'$ satisfies (2.2) with $W' = \Omega W$ instead of W . Here $\Omega(q + \sqrt{-1} q')$ is understood as $\Omega q + \sqrt{-1} \Omega q'$, $q, q' \in F^{r-1}$, and $\Omega q = (\Omega q_1, \dots, \Omega q_{r-1}) \in F^{r-1}$. As in [1, pp. 158–162] one obtains that the space $\mathcal{H}(f, \mathfrak{h})$ of the representation $\rho(f, \mathfrak{h})$ corresponding to the chosen f and $\mathfrak{h}'$ may be realized as a space of complex C^∞ functions ψ on the complex space $\bar{W}'$, square integrable with respect to the measure $\exp(-\sqrt{-1} \langle f, [Y, \bar{Y}] \rangle) dY d\bar{Y}$ ($dY d\bar{Y}$ denoting the Lebesgue measure) on $\bar{W}'$ and satisfying the following functional equation:

$$[\tau(\sqrt{-1} \bar{X}) \psi](\bar{Y}) = \sqrt{-1} [\tau(\bar{X}) \psi](\bar{Y}), \quad X, Y \in W',$$

where $(\tau(\bar{X})\psi)(\bar{Y}) = (d/dt)\psi(\bar{Y} + t\bar{X})|_{t=0}$. The representation ρ is given by

$$\begin{aligned} & [\rho(f, \hbar)(\exp(\bar{X}_0 + X_0 + Z_0))\psi](\bar{X}) \\ &= \exp(\sqrt{-1} \langle f, [X_0, \bar{X}] \rangle - (\sqrt{-1}/2) \langle f, [X_0, \bar{X}_0] \rangle) \chi_f(Z_0) \cdot \psi(\bar{X} - \bar{X}_0), \end{aligned} \quad (2.3)$$

where $\bar{X}_0 + X_0 + Z_0$ is in n^- for $X_0 \in W'$, $Z_0 \in F_0$; $\chi_f(Z_0) = \exp(\sqrt{-1} \langle f, Z_0 \rangle)$. Passing to the complex coordinates $(z_1, \dots, z_{s(r-1)})$ on $\bar{W}'$, according to the identifications

$$\begin{aligned} \bar{X} &= (z_1 \bar{e}_1, \dots, z_{r-1} \bar{e}_1) \leftrightarrow (z_1, \dots, z_{r-1}) \quad \text{for } \mathbf{F} = \mathbf{C}, \\ \bar{X} &= (z_1 \bar{e}_1 + z_2 \bar{e}_2, \dots, z_{2(r-1)-1} \bar{e}_1 + z_{2(r-1)} \bar{e}_2) \leftrightarrow (z_1, \dots, z_{2(r-1)}) \quad \text{for } \mathbf{F} = \mathbf{H}, \\ \bar{X} &= (z_1 \bar{e}_1 + z_2 \bar{e}_2 + z_3 \bar{e}_3 + z_4 \bar{e}_4) \leftrightarrow (z_1, \dots, z_4) \quad \text{for } \mathbf{F} = \mathbf{O}, \end{aligned}$$

we have

$$\begin{aligned} \sqrt{-1} \langle f, [X, \bar{X}'] \rangle &= |p''| \bar{z} \cdot z', \quad X = (\bar{X})^-, \quad \bar{X}, \bar{X}' \in \bar{W}', \\ f &= (0, p''), \quad z = (z_1, \dots, z_{s(r-1)}) \in \mathbf{C}^{s(r-1)}. \end{aligned}$$

Rewriting (2.3) in these coordinates we obtain

PROPOSITION 2.1. *All the inequivalent irreducible unitary representations of N^- fall into two classes:*

(a) *a family of 1-dimensional characters χ_q parameterized by $q' \in \mathbf{F}^{-1}$ and given by*

$$\chi_q(q, p) = \exp(\sqrt{-1} \operatorname{Re}(\bar{q}' \cdot q)), \quad (q, p) \in N^-;$$

(b) *a family of infinite-dimensional representations ρ_p parameterized by $p' \in F_0 \setminus \{0\}$. The Hilbert space $\mathcal{H}_{p'}$ of the representation ρ_p consists of holomorphic functions ψ on $\mathbf{C}^{s(r-1)}$, such that*

$$\|\psi\|_{p'}^2 = \int_{\mathbf{C}^{s(r-1)}} |\psi(z)|^2 \exp(-|p'| |z|^2) dz d\bar{z} < \infty,$$

with

$$dz d\bar{z} = \prod_{j=1}^{s(r-1)} 2d \operatorname{Re} z_j d \operatorname{Im} z_j.$$

The action of ρ_p on $\psi \in \mathcal{H}_{p'}$ is given by

$$(\rho_p(q, p)\psi)(z) = \exp(\sqrt{-1} \operatorname{Re}(\bar{p}'p) + |p'|(\bar{z}_0 \cdot z - \frac{1}{2}|z_0|^2))\psi(z - z_0),$$

$z \in \mathbf{C}^{s(r-1)}$, $(q, p) \in N^-$ with $q = q(z_0)$, $z_0 \in \mathbf{C}^{s(r-1)}$, where

$$q(z) = \Omega(P(z_1, \dots, z_s), \dots, P(z_{s(r-2)+1}, \dots, z_{s(r-1)}))$$

and $P(z_{s(l-1)+1}, \dots, z_{sl})$, $l = 1, \dots, r-1$, is defined as

$$\operatorname{Re} z_l + i \operatorname{Im} z_l \quad \text{for } \mathbf{F} = \mathbf{C},$$

$$\operatorname{Re} z_{2l-1} + i \operatorname{Im} z_{2l-1} + j \operatorname{Im} z_{2l} + k \operatorname{Re} z_{2l} \quad \text{for } \mathbf{F} = \mathbf{H},$$

$$\operatorname{Re} z_1 + i \operatorname{Im} z_1 + j \operatorname{Im} z_2 + k \operatorname{Re} z_2 + e \operatorname{Im} z_3$$

$$+ i e \operatorname{Re} z_3 + j e \operatorname{Re} z_4 + k e \operatorname{Im} z_4 \quad \text{for } \mathbf{F} = \mathbf{O};$$

$\text{Im}(a + \sqrt{-1} b) = b$, $a, b \in \mathbf{R}$. The functions

$$\psi_n^{p'}(z) = (2\pi)^{-m/2} |p'|^{(|n|+m)/2} (n!)^{-1/2} z^n, \quad z \in \mathbf{C}^m,$$

$n = (n_1, \dots, n_m) \in \mathbf{N}^m$, with $n! = n_1! \cdots n_m!$, $z^n = z_1^{n_1} \cdots z_m^{n_m}$, $|n| = n_1 + \cdots + n_m$, $m = s(r-1)$, form an orthonormal basis of $\mathcal{H}_{p'}$, as n runs over $\mathbf{N}^m$.

We also note the following symmetry properties of χ_q and $\rho_{p'}$ relative to the orthogonal and the unitary transformations.

PROPOSITION 2.2. (a) Let $o_1, \dots, o_{r-1} \in O(\sigma, \mathbf{R})$; then for $q' \in \mathbf{F}^{r-1}$,

$$\chi_q(o_1 q_1, \dots, o_{r-1} q_{r-1}, p) = \chi_{(o_1 q'_1, \dots, o_{r-1} q'_{r-1})}(q, p), \quad (q, p) \in N^-.$$

(b) Let $u_1, \dots, u_{r-1} \in U(s)$; then for $p' \in F_0 \setminus \{0\}$,

$$A_u^{-1} \rho_{p'}(q, p) A_u = \rho_{p'}(q^u, p),$$

$(q, p) \in N^-$ with $q = q(z_0)$, $z_0 \in \mathbf{C}^{s(r-1)}$ and $q^u = q(uz_0)$ with $uz = (u_1(z_1, \dots, z_s), \dots, u_{r-1}(z_{s(r-2)+1}, \dots, z_{s(r-1)}))$; $((A_u)\psi)(z) = \psi(uz)$ for $\psi \in \mathcal{H}_{p'}$; $u = (u_1, \dots, u_{r-1})$, $\sigma = 2s = \dim_{\mathbf{R}} \mathbf{F}$.

3. Algebra of multiradial functions.

DEFINITION (cf. GELLER [2]). We say that a function F on $N^- = \mathbf{F}^{r-1} \times \mathbf{F}_0$, $\mathbf{F} = \mathbf{C}, \mathbf{H}, \mathbf{O}$, is *multiradial* if there is a function f on $\mathbf{R}_+^{r-1} \times \mathbf{F}_0$ such that

$$F(q, p) = f(|q_1|, \dots, |q_{r-1}|, p), \quad (q, p) \in N^-. \quad (3.0)$$

PROPOSITION 3.1. Let $\mathcal{Q}$ denote the space of multiradial functions in $L^1(N^-)$. Then $\mathcal{Q}$ is a commutative closed $*$ -subalgebra of $L^1(N^-)$ and $\mathcal{Q}$ is symmetric.

PROOF. 1°. If $F, G \in \mathcal{Q}$ then $F * G \in \mathcal{Q}$. For we have

$$\begin{aligned} F * G(q', p') &= \int f(|q_1|, \dots, |q_{r-1}|, p) \\ &\quad \times g(|q'_1 - q_1|, \dots, |q'_{r-1} - q_{r-1}|, p' - p - \text{Im}(\bar{q} \cdot q')) dq dp. \end{aligned} \quad (3.1)$$

Substituting $q = ((q'_1/|q'_1|)\tilde{q}_1, \dots, (q'_{r-1}/|q'_{r-1}|\tilde{q}_{r-1})$ we get (since $(a\bar{b})b = a(\bar{b}b)$ for $a, b \in \mathbf{F}$)

$$\begin{aligned} \int f(|\tilde{q}_1|, \dots, |\tilde{q}_{r-1}|, p) g\left(|q'_1| |1 - \tilde{q}_1/|q'_1||, \dots, \right. \\ \left. |q'_{r-1}| |1 - \tilde{q}_{r-1}/|q'_{r-1}||, p' - p - \text{Im}\left(\sum_{i=1}^{r-1} \tilde{q}_i |q'_i|\right)\right) d\tilde{q} dp, \end{aligned}$$

i.e. $F * G$ is multiradial. Obviously $\mathcal{Q}$ is closed.

2°. $\mathcal{Q}$ is commutative (cf. Kaplan and Putz [6, p. 377]). Under the orthogonal change of variables

$$q_l \mapsto q''_l = q'_l \cdot 2 \text{Re}(\bar{q}_l q'_l) / |q'_l|^2 - q_l, \quad l = 1, \dots, r-1,$$

one has $|q'_l - q''_l| = |q'_l - q_l|$ and $\text{Im}(\bar{q}_l q'_l) = -\text{Im}(\bar{q}_l'' q'_l)$. Thus (3.1) is equal to

$$\begin{aligned} & \int f(|q''_1|, \dots, |q''_{r-1}|, p) \\ & \quad \times g\left(|q'_1 - q''_1|, \dots, |q'_{r-1} - q''_{r-1}|, p' - p - \sum_{l=1}^{r-1} -\text{Im}(\bar{q}_l'' q'_l)\right) dq'' dp \\ & = \int_{N^-} F(q'', p) G((q', p')(q'', p)^{-1}) dq'' dp \\ & = G * F(q', p'). \end{aligned}$$

3°. Since $L^1(N^-)$ is symmetric (Leptin [8, p. 205]), its $*$ -subalgebra $\mathcal{Q}$ is also symmetric.

PROPOSITION 3.2. *For $F \in \mathcal{Q}$ and $u = (u_1, \dots, u_{r-1}) \in U(s) \times \dots \times U(s)$, the operators $\rho_p(F) = \int_{N^-} F(q, p) \rho_p(q, p) dq dp$ and A_u commute on $\mathcal{H}_p$.*

PROOF. By Proposition 2.2(b),

$$\begin{aligned} & A_u^{-1} \int \rho_p(q(z_0), p) F(q(z_0), p) dq(z_0) dp A_u \\ & = \int \rho_p(q(uz_0), p) F(q(z_0), p) dq(z_0) dp. \end{aligned}$$

Since

$$q_l(uz_0) = \Omega(P(u_l Z_l)) = (\Omega P u_l P^{-1} \Omega^{-1})(\Omega P Z_l)$$

with $Z_l = (z_{s(l-1)+1}^0, \dots, z_{sl}^0)$, $l = 1, \dots, r-1$, and since $\Omega P u_l P^{-1} \Omega^{-1}$ is an orthogonal transformation on $\mathbf{R}^s \approx \mathbf{F}$, and $\Omega P Z_l = q_l(z_0)$, the last integral is equal to $\rho_p(F)$.

REMARK. For $\mathbf{F} = \mathbf{R}$, the corresponding group N^- is $\mathbf{R}^{r-1}$, so the algebra $L^1(N^-)$ is already commutative, and, as in the case $M = \mathbf{R}^n \times \mathbf{R}_+$ with the Euclidean metric [5], we consider $\mathcal{Q} = L^1(N^-)$.

4. Multiplicative linear functionals on $\mathcal{Q}$. Let Φ be a nonzero multiplicative linear functional on $\mathcal{Q}$. Since $\mathcal{Q}$ is a symmetric $*$ -subalgebra of $L^1(N^-)$, there exist an irreducible $*$ -representation π of $L^1(N^-)$ and a unit vector ξ in the Hilbert space $\mathcal{H}_\pi$ such that

$$\pi(F)\xi = \Phi(F)\xi \quad \text{for all } F \text{ in } \mathcal{Q}. \quad (4.1)$$

If $\mathcal{H}_\pi$ is one dimensional, then

$$\pi(F)\xi = \int_{N^-} F(q, p) \chi_q(q, p) dq dp \xi \quad (4.2)$$

for some $q' \in \mathbf{F}^{r-1}$, and by Proposition 2.2(a), if q' and q'' in $\mathbf{F}^{r-1}$ are such that $|q'_l| = |q''_l|$, $l = 1, \dots, r-1$, the Φ 's corresponding by (4.1) and (4.2) to $\chi_{q'}$ and $\chi_{q''}$ are identical. If $\pi \approx \rho_p$, then by Proposition 3.2, $\rho_p(F)$ and A_u commute. Now for $\psi(z) = \psi_1(Z_1) \dots \psi_{r-1}(Z_{r-1})$ with $z = (Z_1, \dots, Z_{r-1})$, $Z_l = (z_{s(l-1)+1}, \dots, z_{sl})$, $l = 1, \dots, r-1$, we have

$$(A_u \psi)(z) = \psi_1(u_1 Z_1) \dots \psi_{r-1}(u_{r-1} Z_{r-1}).$$

Thus putting $\psi_l(Z_l) = Z_l^{n_l}$ with $n_l = (n_{1l}, \dots, n_{sl}) \in \mathbb{N}^s$, we note that A_u preserves the finite-dimensional subspaces of $\mathcal{H}_{p'}$, namely the spaces $\mathcal{H}^n = \bigotimes_{l=1}^{r-1} \mathcal{H}^{|n_l|}$, where $n = (|n_1|, \dots, |n_{r-1}|) \in \mathbb{N}^{r-1}$. $\mathcal{H}^{|n_l|}$ is the space of homogeneous polynomials in $z_{s(l-1)}, \dots, z_{sl}$ of degree $|n_l|$. Moreover, $\mathcal{H}_{p'} = \bigoplus_n \mathcal{H}^n$ —an orthogonal direct sum over $n \in \mathbb{N}^{r-1}$. We also note that A_u restricted to $\mathcal{H}^n$ is equal to $\bigotimes_{l=1}^{r-1} T^{|n_l|}(u_l^{-1}, \dots, u_{sl}^{-1})$ with T^k , $k = |n_l|$, being the representation of $U(s)$ on $\mathcal{H}^k$ given by $(T_u^k \psi)(Z) = \psi(u^{-1}Z)$. Since T^k is irreducible (cf., e.g., [13, pp. 204–209]), the representations $T^n = \bigotimes_l T^{|n_l|}$ of $U(s) \times \dots \times U(s)$, $r-1$ copies of $U(s)$, act irreducibly on $\mathcal{H}^n$, and $T^n \approx T^m$ iff $n = m$. Hence, by Schur's Lemma, every intertwining operator S for $\bigoplus_n T^n$ on $\mathcal{H}_{p'}$ is of the form $S = \bigoplus_n c_n(S) \text{Id}_{\mathcal{H}^n}$. In particular, each $\rho_p(F)$ with $F \in \mathcal{Q}$ is such. It follows from (4.1) that $\Phi(F)$ is equal to one of the constants $c_n(\rho_p(F))$, $n \in \mathbb{N}^{r-1}$. Conversely, for every fixed n , the mapping $F \mapsto c_n(\rho_p(F))$ defines a multiplicative linear functional on $\mathcal{Q}$. Now we shall derive explicit formulas for the constants c_n above. Since, e.g.,

$$c_n(\rho_p(F)) = (\rho_p(F) \psi_{n'}^{p'}, \psi_{n'}^{p'})_{\mathcal{H}_{p'}}$$

with $n' = (n_1, 0, \dots, 0; n_2, 0, \dots, 0; \dots; n_{r-1}, 0, \dots, 0) \in (\mathbb{N}^s)^{r-1}$, we calculate the integral, see Proposition 2.1(b),

$$\int_{\mathbb{C}^{s(r-1)}} [\rho_p(F) \psi_{n'}^{p'}](z) \bar{\psi}_{n'}^{p'}(z) \exp(-|p'| |z|^2) dz d\bar{z}, \quad (4.3)$$

which in expanded form is equal to (with $k = s(r-1)$)

$$\begin{aligned} & (2\pi)^{-k} (n!)^{-1} |p'|^{|n|+k} \\ & \times \int_{\mathbb{C}^k} \int_{\mathbb{F}^{r-1} \times \mathbb{F}_0} F(q(z_0), p) \exp(\sqrt{-1} \operatorname{Re}(\bar{p}'p) + |p'|(\bar{z}_0 \cdot z - \tfrac{1}{2}|z_0|^2)) \\ & \times (z - z_0)^{n'} \bar{z}^{n'} \exp(-|p'| |z|^2) dq(z_0) dp dz d\bar{z}. \end{aligned} \quad (4.4)$$

The integral

$$\int_{\mathbb{C}^{s(r-1)}} (z - z_0)^{n'} \bar{z}^{n'} \exp(-|p'| |z|^2) \exp(|p'| \bar{z}_0 \cdot z) dz d\bar{z} \quad (4.5)$$

is equal to

$$\begin{aligned} & (2\pi/|p'|)^{(s-1)(r-1)} \prod_{l=1}^{r-1} 2\pi n_l! |p'|^{-n_l-1} \sum_{j=0}^{n_l} (-|p'| |z_{s(l-1)+1}^0|^2)^j \binom{n_l}{j} (j!)^{-1} \\ & = (2\pi)^{s(r-1)} |p'|^{-|n|-s(r-1)} n! \prod_{l=1}^{r-1} L_{n_l}(|p'| |z_{s(l-1)+1}^0|^2) \end{aligned} \quad (4.5a)$$

with L_{n_l} being the Laguerre polynomial. (4.5a) is obtained (see [5]) by substituting the binomial formula for $(z - z_0)^n$, developing $\exp(-|p'| \bar{z}_0 \cdot z)$ in a power series and integrating this series term by term using the orthogonality relations for the

functions z^n in $\mathcal{H}_p$. Substituting (4.5a) in (4.4) we obtain that (4.3) is equal to

$$\begin{aligned} & \int_{\mathbf{F}^{r-1} \times \mathbf{F}_0} F(q(z_0), p) \exp(\sqrt{-1} \operatorname{Re}(\bar{p}'p) - \tfrac{1}{2}|p'| |z_0|^2) \\ & \quad \times \prod_{l=1}^{r-1} L_{n_l}(|p'| |z_{s(l-1)+1}^0|^2) dq(z_0) dp \\ & = \int_0^\infty dt_1 \dots \int_0^\infty dt_{r-1} \left(\int_{\mathbf{F}_0} f(t_1, \dots, t_{r-1}, p) \exp(\sqrt{-1} \operatorname{Re}(\bar{p}'p)) dp \right) \\ & \quad \times \exp(-\tfrac{1}{2}|p'| (t_1^2 + \dots + t_{r-1}^2)) \prod_{l=1}^{r-1} t_l^{\sigma-1} g_l, \end{aligned}$$

with f as in (3.0) and g_l given by

$$\begin{aligned} g_l &= \int_{S(\sigma-1)} L_{n_l}(|p'| |z_1|^2) dS(q(Z)), \\ Z &= (z_1, \dots, z_s) \in \mathbf{C}^s, \quad |Z| = t_l, \end{aligned}$$

$S(\sigma-1)$ being the unit sphere in $\mathbf{F}$. Since here $q(Z) = \Omega(PZ)/t_l$, with P as in Proposition 2.1(b) and $\Omega \in O(\sigma, \mathbf{R})$, in order to compute g_l one has to calculate the integrals

$$\int_{S(\sigma-1)} |z_1(q)|^{2j} dS(q), \quad j = 0, \dots, n_l, \quad (4.6)$$

where $|z_1(q)|^2 = t_l^2((q^1)^2 + (q^i)^2)$ with $q^1, q^i, \dots$ denoting the coordinates of q in the (standard) basis $\{1, i, \dots\}$ of $\mathbf{F}$ over $\mathbf{R}$. Now (4.6) is equal to

$$\begin{aligned} & \int_0^{\pi/2} \cos^{2j} \theta \cos \theta \sin^{\sigma-3} \theta d\theta \cdot 2\pi \cdot 2\pi^{s-1} [(s-2)!]^{-1} t_l^{2j} \\ & = 2\pi^s t_l^{2j} j! / (j+s-1)!. \end{aligned}$$

We summarize the results of this section in the following:

PROPOSITION 4. *The multiplicative linear functionals on $\mathcal{Q}$ fall into two classes:*

(a) *the functionals corresponding to $(r-1)$ -tuples $(t_1, \dots, t_{r-1})$ of nonnegative real numbers and given by*

$$F \mapsto \hat{F}(t_1, \dots, t_{r-1}) = \int_{\mathbf{F}^{r-1} \times \mathbf{F}_0} F(q, p) \exp(\sqrt{-1} \operatorname{Re}(\bar{q}' \cdot q)) dq dp$$

with $q' \in \mathbf{F}^{r-1}$ arbitrary provided $(|q'_1|, \dots, |q'_{r-1}|) = (t_1, \dots, t_{r-1})$.

(b) *the functionals corresponding to pairs $(p', n) \in \mathbf{F}_0 \setminus \{0\} \times \mathbf{N}^{r-1}$ and given by*

$$\begin{aligned} F \mapsto \hat{F}(p', n) &= (2\pi^s)^{r-1} \int_{\mathbf{R}_+^{r-1}} \exp\left(-\frac{|p'|}{2}(t_1^2 + \dots + t_{r-1}^2)\right) \prod_{l=1}^{r-1} L_{n_l}^{(s-1)}(|p'| t_l^2) t_l^{\sigma-1} \\ & \quad \times \left(\int_{\mathbf{F}_0} f(t_1, \dots, t_{r-1}, p) \exp(\sqrt{-1} \operatorname{Re}(\bar{p}'p)) dp \right) dt_1 \dots dt_{r-1}, \end{aligned}$$

where $f(|q_1|, \dots, |q_{r-1}|, p) = F(q_1, \dots, q_{r-1}, p)$ and

$$\begin{aligned} L_k^{(m)}(x) &= \sum_{j=0}^k \frac{(-x)^j}{(j+m)!} \binom{k}{j} \\ &= [(k+m)!]^{-1} x^{-m} e^x (d^k/dx^k)(x^{k+m} e^{-x}). \end{aligned} \quad (4.7)$$

5. Nonvanishing of the Gel'fand transform of P_a .

LEMMA 1. For $1 \leq m < 2k + \frac{1}{2}$, $k > \frac{3}{2}$ and $Q > 0$, the following formula holds:

$$\begin{aligned} \int_{\mathbf{R}^m} \frac{\exp(\sqrt{-1} x_0 \cdot x) dx}{(Q^2 + 4|x|^2)^k} &= 2^{-m} \pi^{m/2} \frac{\Gamma(k - m/2)}{\Gamma(k)} Q^{m-2k}, \quad \text{for } x_0 = 0, \\ &= 2^{m+1-4k} \pi^{(m+1)/2} \frac{r^{2k-m} e^{-(r/2)Q}}{\Gamma(k) \Gamma(k - (m-1)/2)} \\ &\quad \times \int_0^\infty e^{-(r/2)Qt} ((t+1)^2 - 1)^{k-(m+1)/2} dt, \\ &\quad \text{for } r = |x_0| \neq 0. \end{aligned} \quad (5.1)$$

PROOF. For $x_0 = 0$, the integral is equal to the "area" of the unit sphere in $\mathbf{R}^m$ ($= 2$ when $m = 1$) times $\int_{\mathbf{R}^+} r^{m-1} (Q^2 + 4r^2)^{-k} dr$ and we substitute $r = r'Q/2$. For $x_0 \neq 0$, the function $[4((\frac{1}{2}Q)^2 + x^2)]^{-k}$ is radial on $\mathbf{R}^m$, hence its Fourier transform (5.1) is equal to, see, e.g., [9, p. 155],

$$4^{-k} (2\pi)^{m/2} r^{-(m-2)/2} \int_0^\infty \left(\left(\frac{1}{2} Q \right)^2 + t^2 \right)^{-k} J_{(m-2)/2}(rt) t^{m/2} dt, \quad m \geq 1, \quad k > 3/2.$$

Combining now the Sonine formula [12, p. 434, (2)],

$$\int_0^\infty \frac{x^{\nu+1} J_\nu(ax) dx}{(x^2 + k^2)^{\mu+1}} = \frac{a^\mu k^{\nu-\mu}}{2^\mu \Gamma(\mu+1)} K_{\mu-\nu}(ak),$$

valid when $-1 < \operatorname{Re}(\nu) < 2 \operatorname{Re}(\mu) + \frac{3}{2}$, with the following expression for the function K [12, p. 172, (4)],

$$K_\nu(z) = \frac{\Gamma(\frac{1}{2}) (\frac{1}{2} z)^\nu}{\Gamma(\nu + \frac{1}{2})} \int_1^\infty e^{-zt} (t^2 - 1)^{\nu-1/2} dt,$$

valid for $\operatorname{Re}(\nu + \frac{1}{2}) > 0$, $|\arg z| < \pi/2$, we obtain (5.1).

LEMMA 2. For $\varepsilon > 0$, $m > 0$, $x \in \mathbf{R}^n$,

$$\begin{aligned} (\varepsilon + |x|^2)^{-m} &= (4\pi)^{-n/2} \Gamma(m)^{-1} \int_{\mathbf{R}^n} \exp(-\sqrt{-1} x \cdot y) \\ &\quad \times \left(\int_0^\infty t^{m-1-m/2} e^{-\varepsilon t} e^{-|y|^2/4t} dt \right) dy, \end{aligned}$$

i.e. $(\varepsilon + |x|^2)^{-m}$ is a Fourier transform of a positive function in $L^1(\mathbf{R}^n)$.

PROOF. Combine

$$(\varepsilon + |x|^2)^{-m} = \Gamma(m)^{-1} \int_0^\infty t^{m-1} e^{-(\varepsilon + |x|^2)t} dt, \quad m > 0,$$

with

$$\exp(-|x|^2 t) = (4\pi t)^{-n/2} \int_{\mathbf{R}^n} \exp(-|y|^2/4t) \exp(-\sqrt{-1} x \cdot y) dy,$$

and note that the obtained double integral is absolutely convergent.

PROPOSITION 5. For every $a \in A$, $\hat{P}_a$ does not vanish on $\mathfrak{M}(\mathcal{Q})$ —the maximal ideal space of $\mathcal{Q}$.

PROOF. (a) For the points $(t_1, \dots, t_{r-1}) \in \mathfrak{M}(\mathcal{Q})$, integrating over $\mathbf{F}_0$ in the formula of Proposition 4(a), according to Lemma 1 (with $x_0 = 0$) we get

$$\hat{P}_a(t_1, \dots, t_{r-1}) = c \int_{\mathbf{R}^{\sigma(r-1)}} \frac{\exp(\sqrt{-1} \operatorname{Re}(\bar{q}' \cdot q))}{(\varepsilon + |q|^2)^{d+1-\sigma}} dq,$$

with

$$c = 2^{1-\sigma} \pi^{s-1/2} \varepsilon^{d/2} c_{r,\mathbf{F}} \Gamma(d/2 - s + 1/2) \Gamma(d/2)^{-1},$$

and this is positive by Lemma 2 and the Fourier inversion formula.

(b) For the points $(p', n) \in \mathfrak{M}(\mathcal{Q})$ with $p' \in \mathbf{F}_0 \setminus \{0\}$, $n = (n_1, \dots, n_{r-1}) \in \mathbf{N}^{r-1}$, we use the formula from Proposition 4(b) for $\hat{P}_a(p', n)$. Applying Lemma 1 to the integral over $\mathbf{F}_0$ there (with $Q = \varepsilon + t_1^2 + \dots + t_{r-1}^2$, $m = \sigma - 1$, $x_0 = p'$, $k = d/2$), then interchanging the order of integration from $dt dt_1 \dots dt_{r-1}$ to $dt_1 \dots dt_{r-1} dt$, making the change of variables $(|p'|t_1^2, \dots, |p'|t_{r-1}^2) = (x_1, \dots, x_{r-1})$, and finally applying (4.7), we obtain

$$\hat{P}_a(p', n) = c \int_0^\infty e^{-|p'|et/2} ((t+1)^2 - 1)^{d/2-s} \prod_{l=1}^{r-1} \mathcal{J}_l(t) dt,$$

with

$$c = 2^{s(r-1)} (\varepsilon |p'|)^{d/2} \Gamma(d/2)^{-1} \exp(-|p'| \varepsilon / 2)$$

and

$$\mathcal{J}_l(t) = ((n_l + s - 1)!)^{-1} \int_0^\infty e^{-tx/2} (d^{n_l}/dx^{n_l})(x^{n_l+s-1} e^{-x}) dx.$$

Integrating by parts get

$$\begin{aligned} \mathcal{J}_l(t) &= ((n_l + s - 1)!)^{-1} \int_0^\infty e^{-tx/2} x^{n_l+s-1} e^{-x} dx \cdot (t/2)^{n_l} \\ &= (t/2)^{n_l} (t/2 + 1)^{-(n_l+s)}. \end{aligned}$$

Thus $\hat{P}_a(p', n)$ is positive.

REMARK. For M being the real hyperbolic space, i.e. for $\mathbf{F} = \mathbf{R}$, we have $N^- = \mathbf{R}^d$, $\mathfrak{M}(L^1(N^-)) = \hat{N}^-$, $P_a(X_{-a}) = c_{d,\mathbf{R}} \varepsilon^{d/2} (\varepsilon + |q|^2)^{-d}$ and, by Lemma 2, $\hat{P}_a > 0$ on $\hat{N}^-$.

6. Theorem on ideals in $L^1(N^-)$. Since the algebras $\mathcal{Q}$ we consider here have the same qualitative properties as the one considered in [5], similar facts can be proved about them. In particular, the following statement about ideals in $L^1(N^-)$ is a consequence of the Wiener property of $\mathcal{Q}$ and existence of the approximate identity for $L^1(N^-)$ in $\mathcal{Q}$ (the dilations δ_t , $t > 0$, on N^- used in the construction of the approximate identity are given by $\delta_t(q, p) = (t^{-1/2}q, t^{-1}p)$).

PROPOSITION 6. *If $\mathcal{G}$ is a proper closed right ideal in $L^1(N^-)$, then there is a Φ in $\mathfrak{M}(\mathcal{Q})$ such that $\hat{F}(\Phi) = 0$ for every $F \in \mathcal{G} \cap \mathcal{Q}$.*

7. Proof of the Theorem [5]. The Theorem follows now from Proposition 6, for if we put

$$\mathcal{G} = \left\{ f \in L^1(N^-) : \lim_{N^- \ni n \rightarrow \infty} \varphi * f(n) = c_0 \int_{N^-} f(n) \, dn \right\},$$

with $\varphi \in L^\infty(N^-)$ being the boundary value of the bounded harmonic function u on M , then $P_{a_0} \in \mathcal{G} \cap \mathcal{Q}$ and $\hat{P}_{a_0} \neq 0$ on $\mathfrak{M}(\mathcal{Q})$, so $\mathcal{G} = L^1(N^-)$. Hence $P_a \in \mathcal{G}$ for every a in A .

ADDED IN PROOF. Meanwhile Korányi [14] described the Gel'fand space, as well as the related Plancherel formula, for the commutative algebra $\mathcal{Q}$ of *biradial* functions in $L^1(N^-)$, i.e. the functions F such that

$$F(q, p) = f(|q|, |p|), \quad (q, p) \in N^-,$$

for some f on $\mathbf{R}_+ \times \mathbf{R}_+$, cf. §3. His approach uses neither the classification of symmetric spaces nor the representations of nilpotent groups.

REFERENCES

1. P. Bernat et al., *Représentations des groupes de Lie résolubles*, Dunod, Paris, 1972.
2. D. Geller, *Fourier analysis on the Heisenberg group*, Proc. Nat. Acad. Sci. U.S.A. **74** (1977), 1328–1331.
3. S. Helgason, *Differential geometry and symmetric spaces*, Academic Press, New York and London, 1962.
4. ———, *Application of the Radon transform to representations of semisimple Lie groups*, Proc. Nat. Acad. Sci. U.S.A. **63** (1969), 643–647.
5. A. Hulanicki and F. Ricci, *A Tauberian theorem and tangential convergence for bounded harmonic functions on balls in $\mathbf{C}^n$* , Invent. Math. **62** (1980), 325–331.
6. A. Kaplan and R. Putz, *Boundary behavior of harmonic forms on a rank one symmetric space*, Trans. Amer. Math. Soc. **231** (1977), 369–384.
7. A. Korányi, *Boundary behavior of Poisson integrals on symmetric spaces*, Trans. Amer. Math. Soc. **140** (1969), 393–409.
8. H. Leptin, *Harmonische Analyse auf gewissen nilpotenten Lieschen Gruppen*, Studia Math. **48** (1973), 201–205.
9. E. M. Stein and G. Weiss, *Introduction to Fourier analysis on Euclidean spaces*, Princeton Univ. Press, Princeton, N. J., 1971.
10. R. Takahashi, *Quelques résultats sur l'analyse harmonique dans l'espace symétrique non compact de rang un du type exceptionnel*, Analyse Harmonique sur les Groupes de Lie. II, Lecture Notes in Math., vol. 739, Springer-Verlag, Berlin and New York, 1979.

11. G. Warner, *Harmonic analysis on semi-simple Lie groups*. I, Springer-Verlag, Berlin and New York, 1972.
12. G. N. Watson, *A treatise on the theory of Bessel functions*, 2nd ed., Cambridge Univ. Press, Cambridge, 1966.
13. D. P. Zhelobenko, *Kompaktnye gruppy Li i ikh predstavleniya*, Nauka, Moskva, 1970.
14. A. Korányi, *Some applications of Gelfand pairs in classical analysis*, Harmonic Analysis and Group Representations (C.I.M.E., Cortona, 1980), Liguori, Naples (to appear).

INSTITUTE OF MATHEMATICS, UNIVERSITY OF WROCŁAW, PL. GRUNWALDZKI 2/4, 50-384 WROCŁAW, POLAND